

Weak nonleptonic hyperon decays in relativistic χ PT

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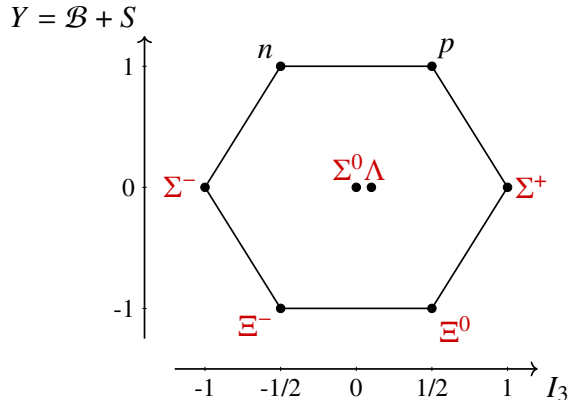
 GSI



Subject of study

- $\Delta S = 1$ transitions:

- $\Sigma^+ \rightarrow n\pi^+$
- $\Sigma^- \rightarrow n\pi^-$
- $\Lambda \rightarrow p\pi^-$
- $\Xi^- \rightarrow \Lambda\pi^-$



$$\sqrt{2}A(\Sigma^+ \rightarrow p\pi^0) - A(\Sigma^+ \rightarrow n\pi^+) + A(\Sigma^- \rightarrow n\pi^-) = 0$$

$$A(\Lambda \rightarrow p\pi^-) + \sqrt{2}A(\Lambda \rightarrow n\pi^0) = 0$$

$$A(\Xi^- \rightarrow \Lambda\pi^-) + \sqrt{2}A(\Xi^0 \rightarrow \Lambda\pi^0) = 0$$

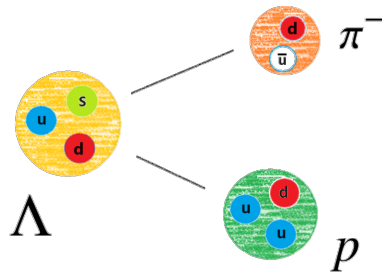
Nonleptonic decays

Transition amplitude

$$\mathcal{M}(B_i \rightarrow B_f \pi) = G_F m_{\pi^+}^2 \bar{u}_f (A^{(S)} + A^{(P)} \gamma_5) u_i$$

From angular momentum conservation:

$$J(1/2) = S_b(1/2) + S_\pi(0) + L_{b\pi}$$
$$S(L_{b\pi} = 0), \quad P(L_{b\pi} = 1)$$



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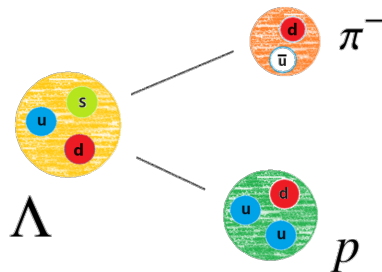
$$J(1/2) = S_b(1/2) + S_\pi(0) + L_{b\pi}$$

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Dimensionless l -wave amplitudes

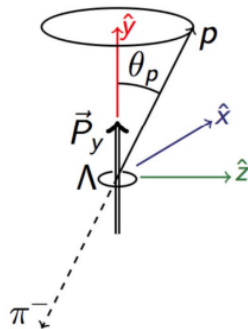
$$\text{parity-violating: } A^{(S)} \equiv S$$

$$\text{parity-conserving: } A^{(P)} \equiv \frac{|\vec{\mathbf{p}}_f|}{E_f + m_f} P$$



Nonleptonic decay parameters

From partial waves to observables:



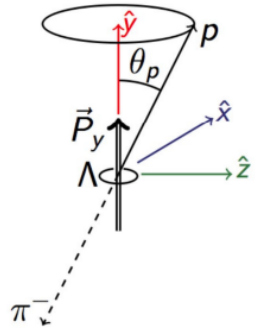
$\Lambda \rightarrow p \pi^-$ decay

Nonleptonic decay parameters

From partial waves to observables:

- Angular distribution $\frac{d\Gamma}{d\Omega} \propto 1 + \alpha \mathbf{P}_\Lambda \cdot \hat{\mathbf{n}}$

$$\alpha = \frac{2\Re(S^*P)}{|S|^2 + |P|^2}$$



$\Lambda \rightarrow p \pi^-$ decay

Nonleptonic decay parameters

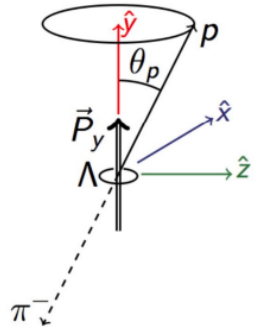
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- Spin $\mathbf{s}_\Lambda \rightarrow \mathbf{s}_p$ rotation

$$\beta = \frac{2\Im(S^*P)}{|S|^2 + |P|^2} = \sqrt{1 - \alpha^2} \sin \phi$$



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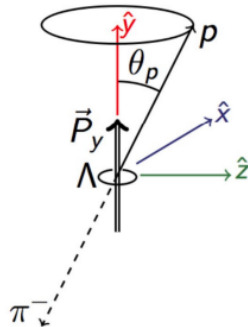
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- Decay width

$$\Gamma = \text{kin}(|\vec{\mathbf{p}}_f|, E_f, m_f)(|S|^2 + |P|^2)$$



$\Lambda \rightarrow p \pi^-$ decay

Motivation: new data landscape



Polarization and entanglement in baryon-antibaryon pair production in electron-positron annihilation

The BESIII Collaboration*

[Nature Phys. 15 (2019) 631]

Motivation: new data landscape



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Probing CP symmetry and weak phases with entangled double-strange baryons

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Precise Measurements of Decay Parameters and CP Asymmetry with Entangled $\Lambda-\bar{\Lambda}$ Pairs

M. Ablikim *et al.* (BESIII Collaboration)

Phys. Rev. Lett. **129**, 131801 – Published 22 September 2022

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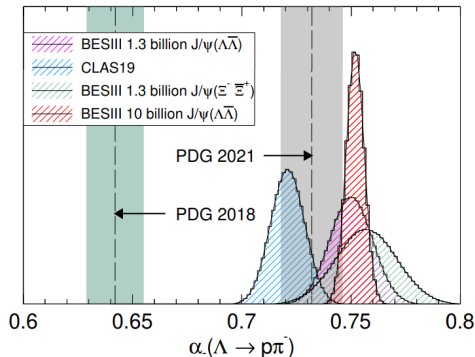
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Extrapolate data points

Next step

Use α , ϕ , Γ data to extrapolate updated experimental L -wave amplitude values.

L -wave amplitude extraction: assuming CP conservation, $\Delta I = 1/2$ rule

$$L = \sum_j L_j \exp(i\delta_j^L), \quad j \in \{2\Delta I, 2I\}, \quad L_j \in \mathbb{R}$$

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and final-interaction phase shifts [NS et al., PRD105, 116022 (2022)]

	$ \mathbf{q} $ [MeV/c]	δ_1^S [°]	δ_3^S [°]	δ_1^P [°]	δ_3^P [°]
$\Lambda \rightarrow N\pi$	103	6.52(9)	-4.60(7)	-0.79(8)	-0.75(4)
$\Sigma \rightarrow N\pi$	190	9.98(23)	-10.70(13)	-0.04(33)	-3.27(15)

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Relative sign between amplitudes fixed by Lee-Sugawara relation

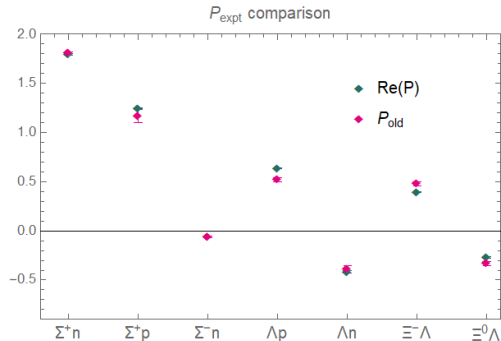
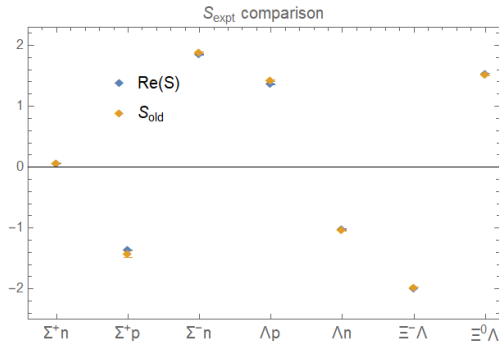
$$\frac{3}{\sqrt{6}} A^{(S)}(\Sigma^- \rightarrow n\pi^-) + A^{(S)}(\Lambda \rightarrow p\pi^-) + 2A^{(S)}(\Xi^- \rightarrow \Lambda\pi^-) = 0$$

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Use α , ϕ , Γ data to extrapolate updated experimental L -wave amplitude values.

New reference values extracted from current data compared to [E. Jenkins, NPB 375 (1992) 561-581]:



Previous values extracted on the assumption of real-valued amplitudes:
updated to complex-valued l -waves on most recent data.

Framework

- At low-energy regime, α_s too large for a perturbative description of hadron interactions from QCD.
- χ PT: EFT with hadrons as DF parametrizes meson-baryon interactions
[S. Weinberg, Physica A 96 (1979) 1-2, 327-340], [J. Gasser & H. Leutwyler, Annals Phys. 158 (1984) 142].

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Weak nonleptonic hyperon decays previously studied in

Nuclear Physics B261 (1985) 185-198
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ON THE VALIDITY OF CHIRAL PERTURBATION THEORY FOR WEAK HYPERON DECAYS*

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California Institute of Technology, Pasadena, CA 91125, USA

Received 23 January 1985
(Revised 14 June 1985)

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Non-leptonic hyperon decays in chiral perturbation theory

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Hyperon non-leptonic decays in chiral perturbation theory

Elizabeth Jenkins

*Department of Physics 0319, University of California, San Diego, 9500 Gilman Drive, La Jolla,
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CHIRAL PERTURBATION THEORY HYPERON DECAYS*

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Starting point

Computing 1-loop corrections using *Heavy-Baryon* χ PT (nonrelativistic approach):

Ref.	$O(\text{trunc})$	10	h_π	B^*
Bijnens, Sonoda & Wise [NPB 261 (1985) 185-198]	$M_K^2 \log M_K$	✗	✗	✗
Jenkins [NPB 375 (1992) 561-581]	$\parallel$	✓	✓	✗
Abd El-Hady & Tandean [PRD 61, 114014 (2000)]	$\parallel$	✓	✓ ¹	✗
Borasoy & Holstein [EPJC 6 (1999) 85-107]	$M_K^2 (a + b \log M_K)$	✗	✓	✗
Borasoy & Holstein [PRD 59, 094025 (1999)]	M_K^2	✗	✓	✓

¹Contradicting previous results.

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General conclusions

- LO chiral corrections to S -waves are in good agreement with experiment. P -waves are not well-described.
- A tree-level fit suggests **resonance exchange** is relevant.
- Results from simultaneous fitting are presumed unsatisfactory: **yet unexplored**.

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Relativistic χ PT and EOMS

- Dimensional regularization + modified minimal subtraction do not work for baryons;
- inclusion of baryon masses breaks power-counting for the loop diagrams.

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Previous approach: HB χ PT Jenkins & Manohar [PLB 255 (1991) 558-562]

- Expansion of the Lagrangian in powers of $1/m_B$ in the chiral limit.
- Power-counting is manifest, but Lorentz invariance is not.
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Extended On-Mass-Shell ren. scheme Gegelia & Japaridze [PRD 60 (1999) 114038]

- Subtract the power-counting-violating terms, i.e. choosing appropriate renormalization conditions.
- Manifest Lorentz invariance: the standard approach for B χ PT nowadays.

Procedure

Compute 1-loop corrections from relativistic LO Lagrangian

$$\mathcal{L}_{\phi B}^{\text{s}} + \mathcal{L}_{\phi B}^{\text{w}}$$

$$\mathcal{L}_{\phi B}^{\text{w}} = h_D \text{tr} \bar{B} \{ \xi^\dagger h \xi, B \} + h_F \text{tr} \bar{B} [\xi^\dagger h \xi, B] + h_C \text{tr} \bar{T}^\mu (\xi^\dagger h \xi) T_\mu$$

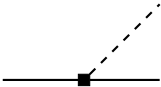
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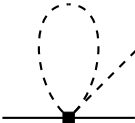
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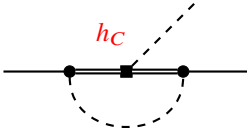
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E.g. S-wave contributions:


 $= O(M_K^0)$


 $= O(M_K^2 \log M_K^2)$

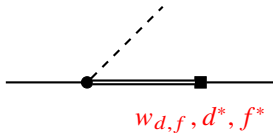



 $= O(M_K^2 (a + b \log M_K^2))$

Resonance saturation

Inclusion of lower-lying $\frac{1}{2}^{\pm}$ resonances [Borasoy & Holstein, PRD 59, 094025 (1999)]¹

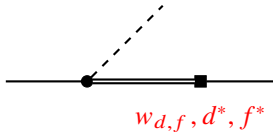
$$\mathcal{L}_{\text{res}}^{\text{w}} \propto d^* \left[\text{tr}(\bar{R}^+ \{ \xi^\dagger h \xi, B \}) + \text{tr}(\bar{B} \{ \xi^\dagger h \xi, R^+ \}) \right] + f^* \left[\text{tr}(\bar{R}^+ [\xi^\dagger h \xi, B]) + \text{tr}(\bar{B} [\xi^\dagger h \xi, R^+]) \right] \\ + i w_d \left[\text{tr}(\bar{R}^- \{ \xi^\dagger h \xi, B \}) - \text{tr}(\bar{B} \{ \xi^\dagger h \xi, R^- \}) \right] + i w_f \left[\text{tr}(\bar{R}^- [\xi^\dagger h \xi, B]) - \text{tr}(\bar{B} [\xi^\dagger h \xi, R^-]) \right]$$



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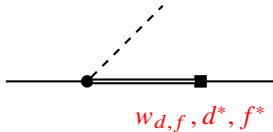
The resulting amplitudes:

$$S_{\text{theory}}, P_{\text{theory}} = l.c. (h_{D,F,C}, w_{d,f}, d^*, f^*)$$

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Update for Roper multiplet

LEC	SL new fit	Ref. [1]
D^*	0.72 ± 0.11	0.60 ± 0.41
F^*	0.36 ± 0.05	0.11 ± 0.41

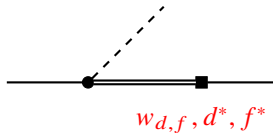
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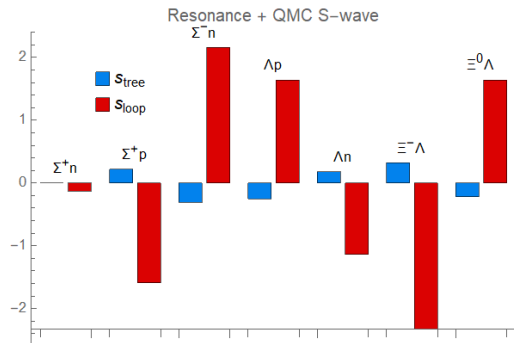
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Goal

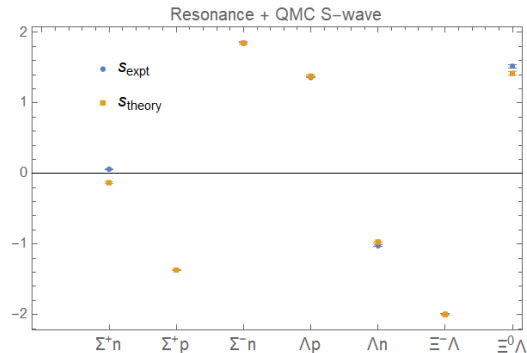
To fit L_{theory} to L_{expt} using least squares method to obtain LEC's values.

Preliminary results – S -waves

LECs extracted from fit to S -waves only:

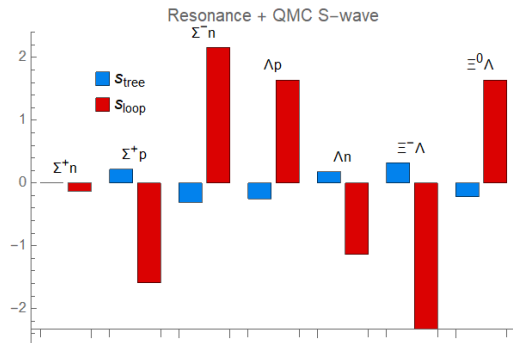


LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$
h_D	$-1/3 h_F$
h_F	-0.233 ± 0.007
h_C	1.99 ± 0.04
w_f	4.63 ± 0.07
w_d	-14.72 ± 0.11

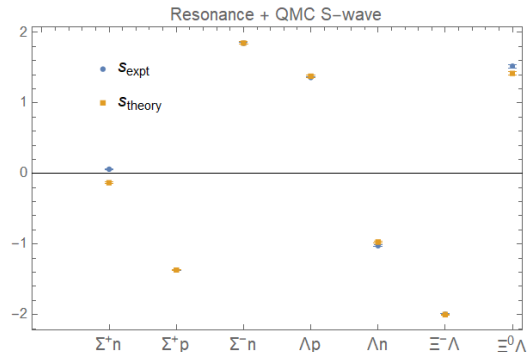


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w_d	-14.72 ± 0.11

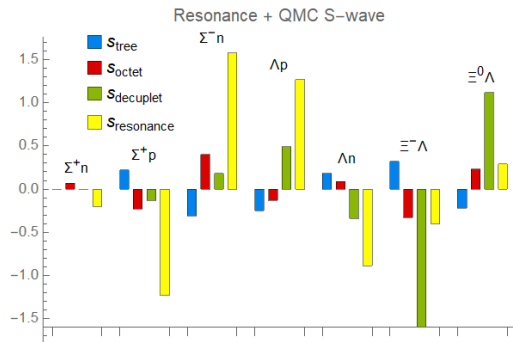


EOMS + resonances

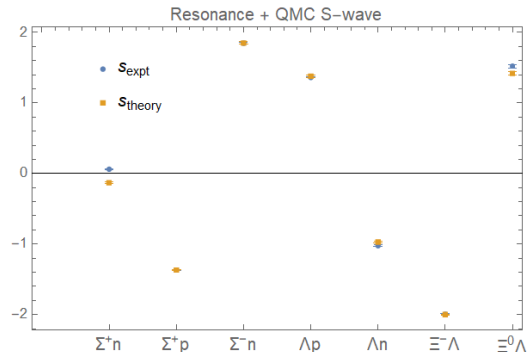
- Good agreement with experiment.
- Corrections to LO are large.

Preliminary results – S -waves

LECs extracted from fit to S -waves only:



LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$
h_D	$-1/3 h_F$
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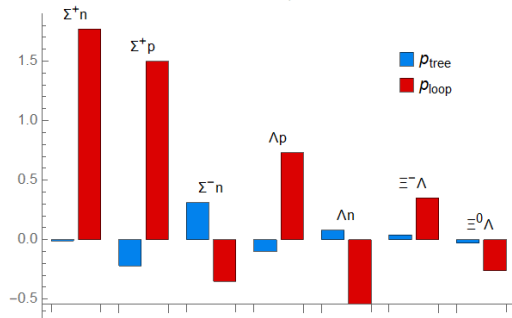
EOMS + resonances

- Good agreement with experiment.
- Corrections to LO are large.
- The importance of 10 and B^* is confirmed.

Preliminary results – P -waves

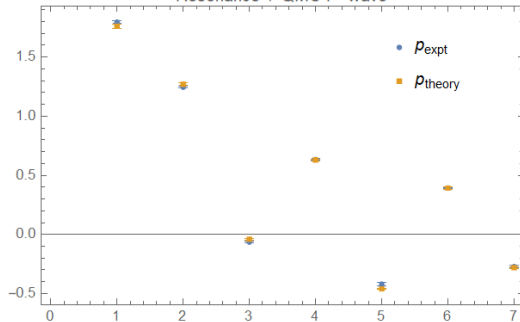
LECs extracted from fit to **P -waves only**:

Resonance + QMC P -wave



LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$
h_D	$-1/3 h_F$
h_F	0.221 ± 0.002
h_C	0.095 ± 0.003
d^*	1.61 ± 0.04
f^*	-3.955 ± 0.034

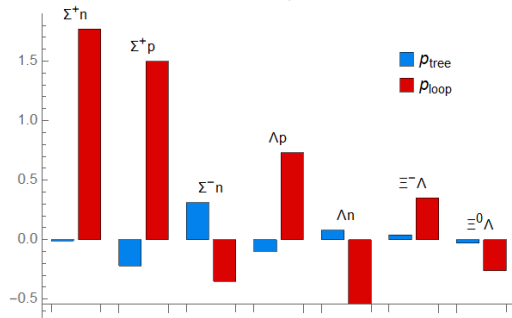
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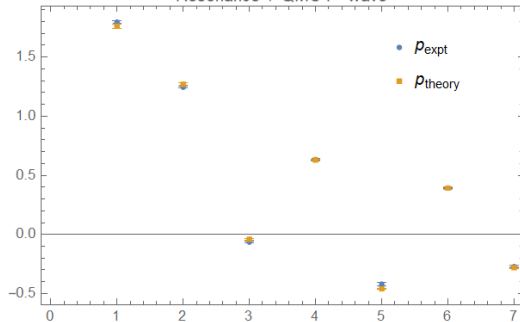
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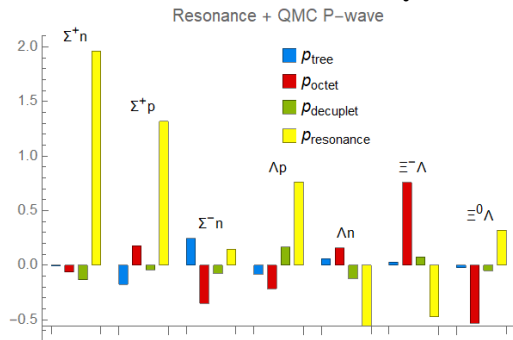


EOMS + resonances

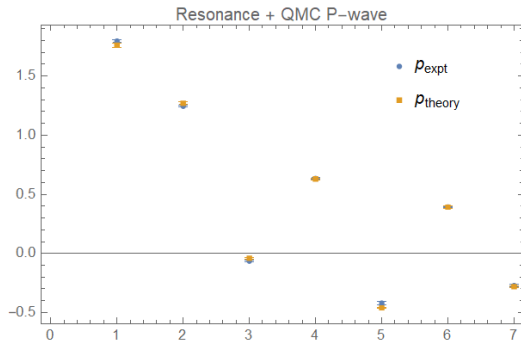
- Agreement with experiment better than in S -waves.

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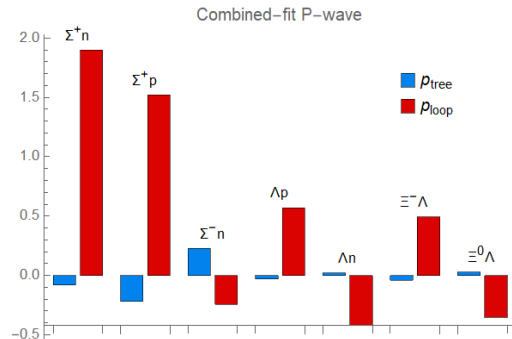
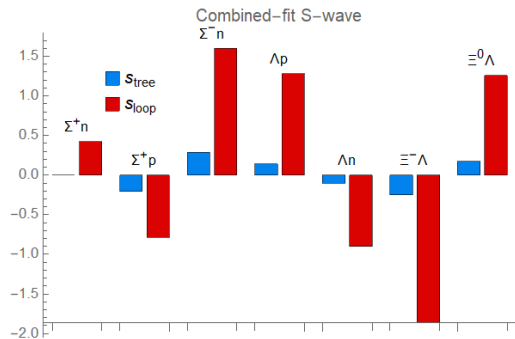


EOMS + resonances

- Agreement with experiment better than in S -waves.
- Similar resonance-to-tree-level relative size as in P -waves.

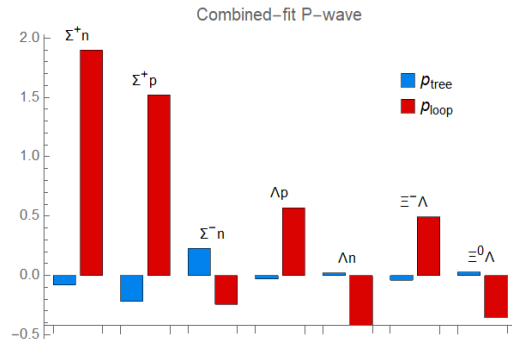
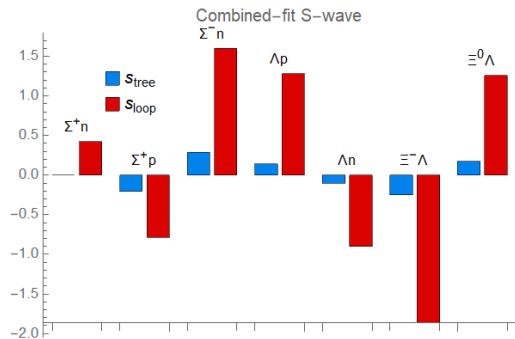
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LECs extracted from fit to **combined S - and P -waves**:



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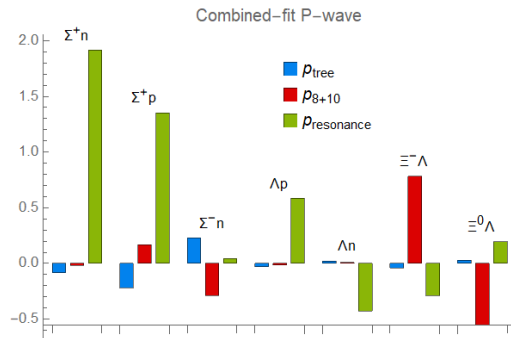
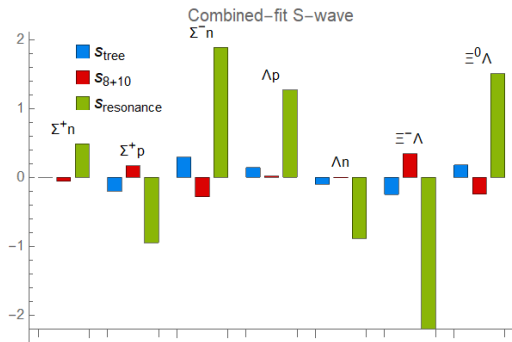
EOMS + resonances

- LEC's size somewhat consistent with other results.

LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$	LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$
h_D	-0.192 ± 0.004	w_d	-9.73 ± 0.06
h_F	0.186 ± 0.002	w_f	7.71 ± 0.02
h_C	-0.136 ± 0.008	d^*	2.72 ± 0.05
		f^*	-2.83 ± 0.04

Preliminary results – Combining S and P

LECs extracted from fit to **combined S - and P -waves**:



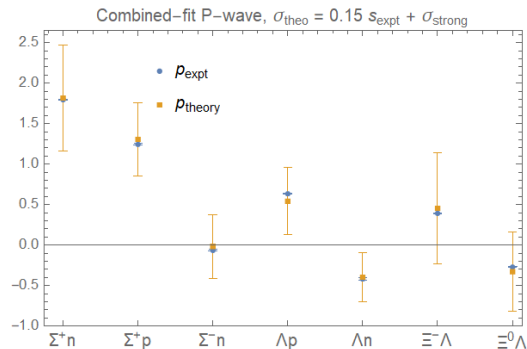
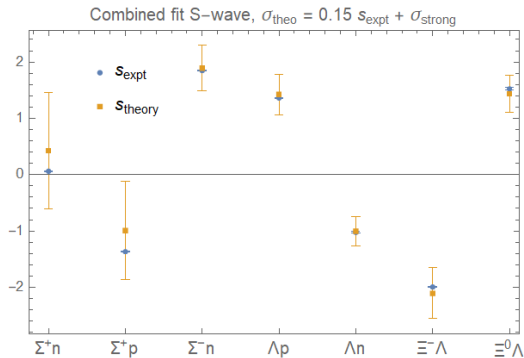
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- S agreement with experiment slightly worse, P remains good.

LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$	LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$
h_D	-0.13 ± 0.23	w_d	-7.3 ± 1.3
h_F	0.16 ± 0.09	w_f	7.7 ± 1.3
h_C	0.0 ± 0.4	d^*	1.1 ± 1.8
		f^*	-2.3 ± 1.5

Conclusions

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- Perform a comparison fit using 8 independent $\Delta I = 1/2$ amplitudes to circumvent isospin relations.
- Outlook: two-loops, $\Lambda(1405)$ as hadron molecular state, resonances included as DF...

Thank You for your
attention.

Any questions?

(this is the part where
you run)

someecards
user card



Relativistic chiral LO Lagrangian

Meson-baryon LO Lagrangian

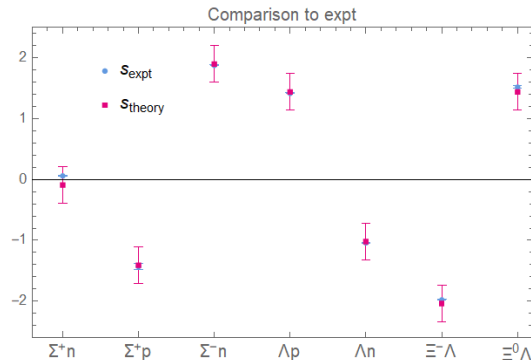
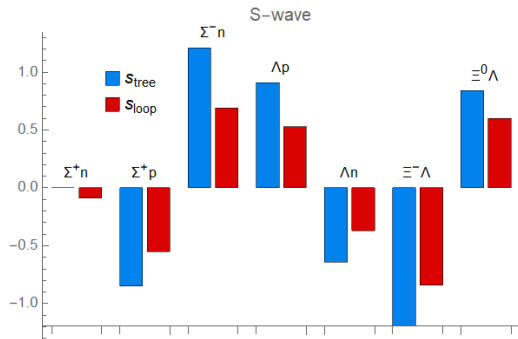
$$\begin{aligned}\mathcal{L}_{\phi B}^s = & i \operatorname{tr} \bar{B} \not{D} B - m_B \operatorname{tr} \bar{B} B + D \operatorname{tr} \bar{B} \gamma^\mu \gamma_5 \{A_\mu, B\} + F \operatorname{tr} \bar{B} \gamma^\mu \gamma_5 [A_\mu, B] - i \bar{T}^\mu \not{D} T_\mu \\ & + m_T \bar{T}^\mu T_\mu + C (\bar{T}^\mu A_\mu B + \bar{B} A_\mu T^\mu) + \mathcal{H} \bar{T}^\mu \gamma_\nu \gamma_5 A^\nu T_\mu + \frac{f^2}{4} \operatorname{tr} \partial_\mu \Sigma \partial^\mu \Sigma^\dagger\end{aligned}$$

Inclusion of $\frac{1}{2}^\pm$ resonances [B. Borasoy & B. Holstein, PRD 59, 094025 (1999)]

$$\begin{aligned}\mathcal{L}_{RB}^s = & 2s_d [\operatorname{tr}(\bar{R} \gamma_\mu \{A_\mu, B\}) - \operatorname{tr}(\bar{B} \gamma_\mu \{A_\mu, R\})] \\ & + 2s_f [\operatorname{tr}(\bar{R} \gamma_\mu [A_\mu, B]) - \operatorname{tr}(\bar{B} \gamma_\mu [A_\mu, R])] \\ \mathcal{L}_{B^*B}^s = & \frac{D^*}{2} [\operatorname{tr}(\bar{B}^* \gamma_\mu \gamma_5 \{A_\mu, B\}) + \operatorname{tr}(\bar{B} \gamma_\mu \gamma_5 \{A_\mu, B^*\})] \\ & + \frac{F^*}{2} [\operatorname{tr}(\bar{B}^* \gamma_\mu \gamma_5 [A_\mu, B]) + \operatorname{tr}(\bar{B} \gamma_\mu \gamma_5 [A_\mu, B^*])] \\ V^\mu = & \frac{1}{2} (\xi \partial^\mu \xi^\dagger + \xi^\dagger \partial^\mu \xi), \quad A^\mu = \frac{i}{2} (\xi \partial^\mu \xi^\dagger - \xi^\dagger \partial^\mu \xi) \\ \xi = & \exp \frac{i\pi}{f}, \quad \Sigma = \xi^2 = \exp \frac{2i\pi}{f}\end{aligned}$$

Previous work results – S -waves

LECs extracted from fit to S -waves only:



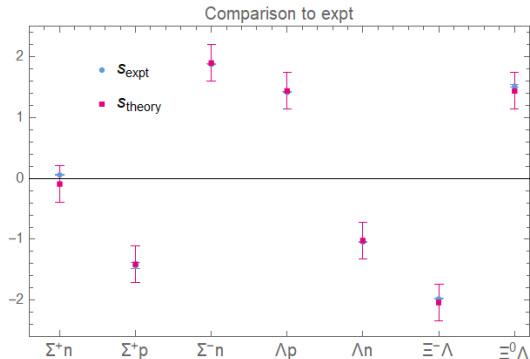
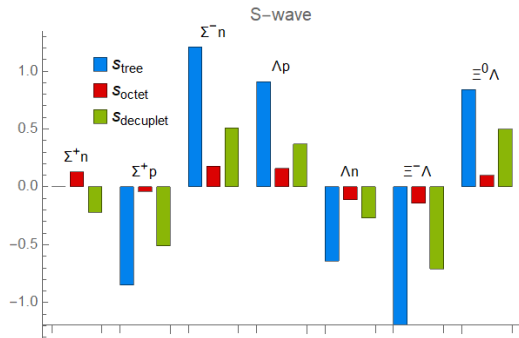
LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$
h_D	-0.35 ± 0.09
h_F	0.86 ± 0.05
h_C	-0.36 ± 0.65

[NPB 375 (1992) 561-581]

- Good agreement with experiment.

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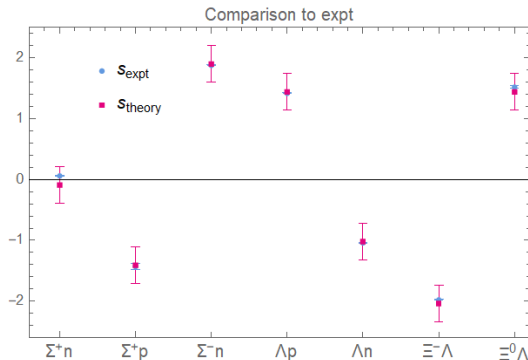
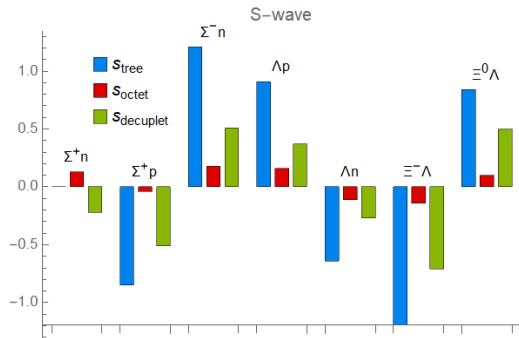
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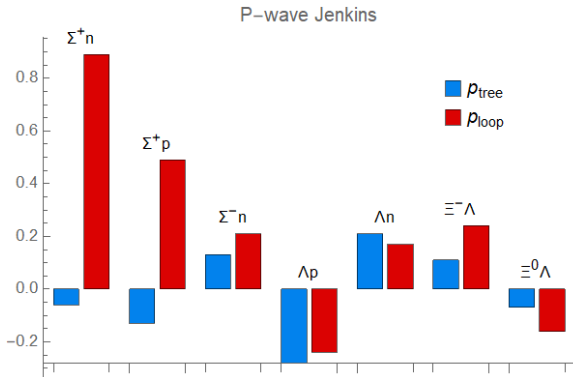
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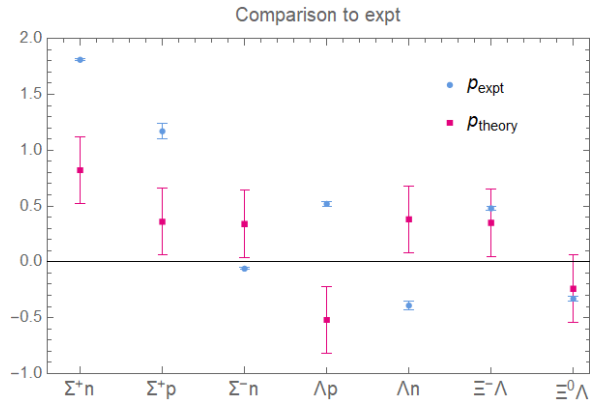
- Good agreement with experiment.
- Decuplet contribution dominates over octet.
- h_C not well determined by 1-loop fit.

Previous work results – P -waves

Using the LECs from S -wave fit:



LEC	$[G_F m_\pi^2 \sqrt{2} f_\pi]$
h_D	-0.35 ± 0.09
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- P -waves are poorly described.
- Simultaneous fitting seems unfeasible.