

NNLO phase-space integrals for semi-inclusive deep-inelastic scattering

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(IMSc, India)

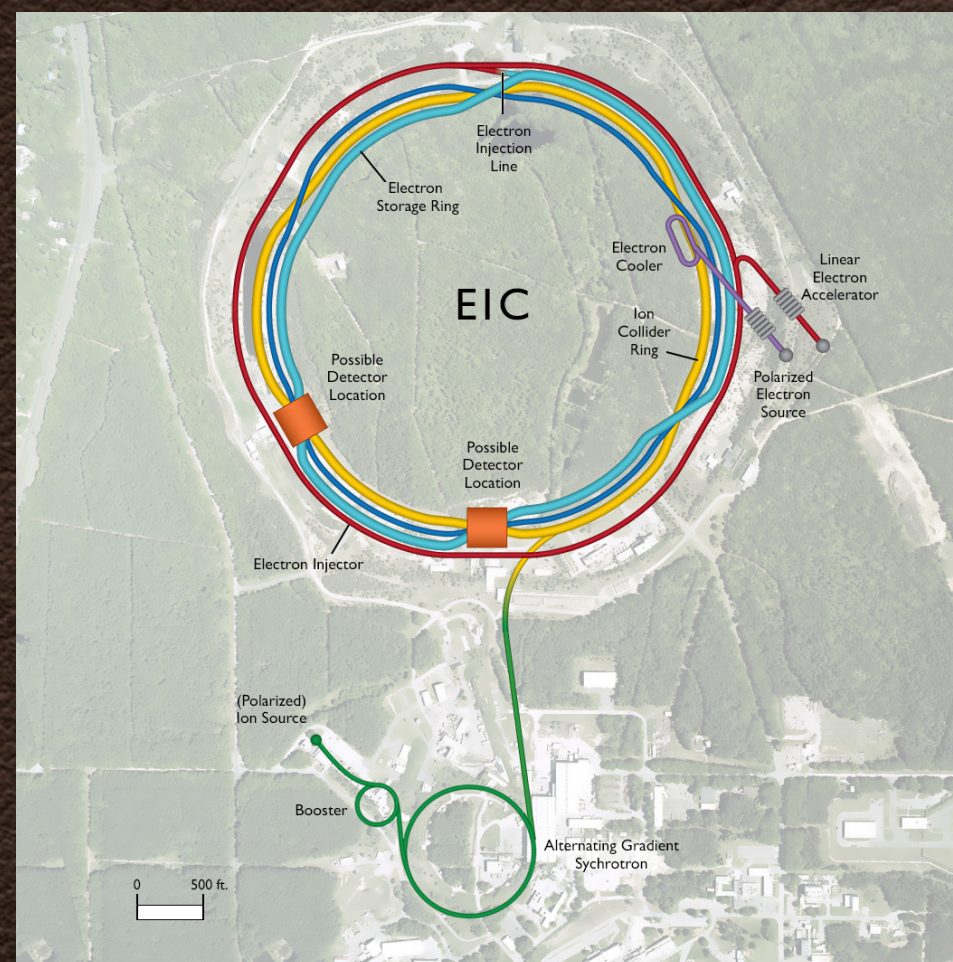
Roman Lee (Novosibirsk, IYF)

Narayan Rana (Bhubaneswar, NISER)

arXiv: 2412.16509v1



Deutsche
Forschungsgemeinschaft

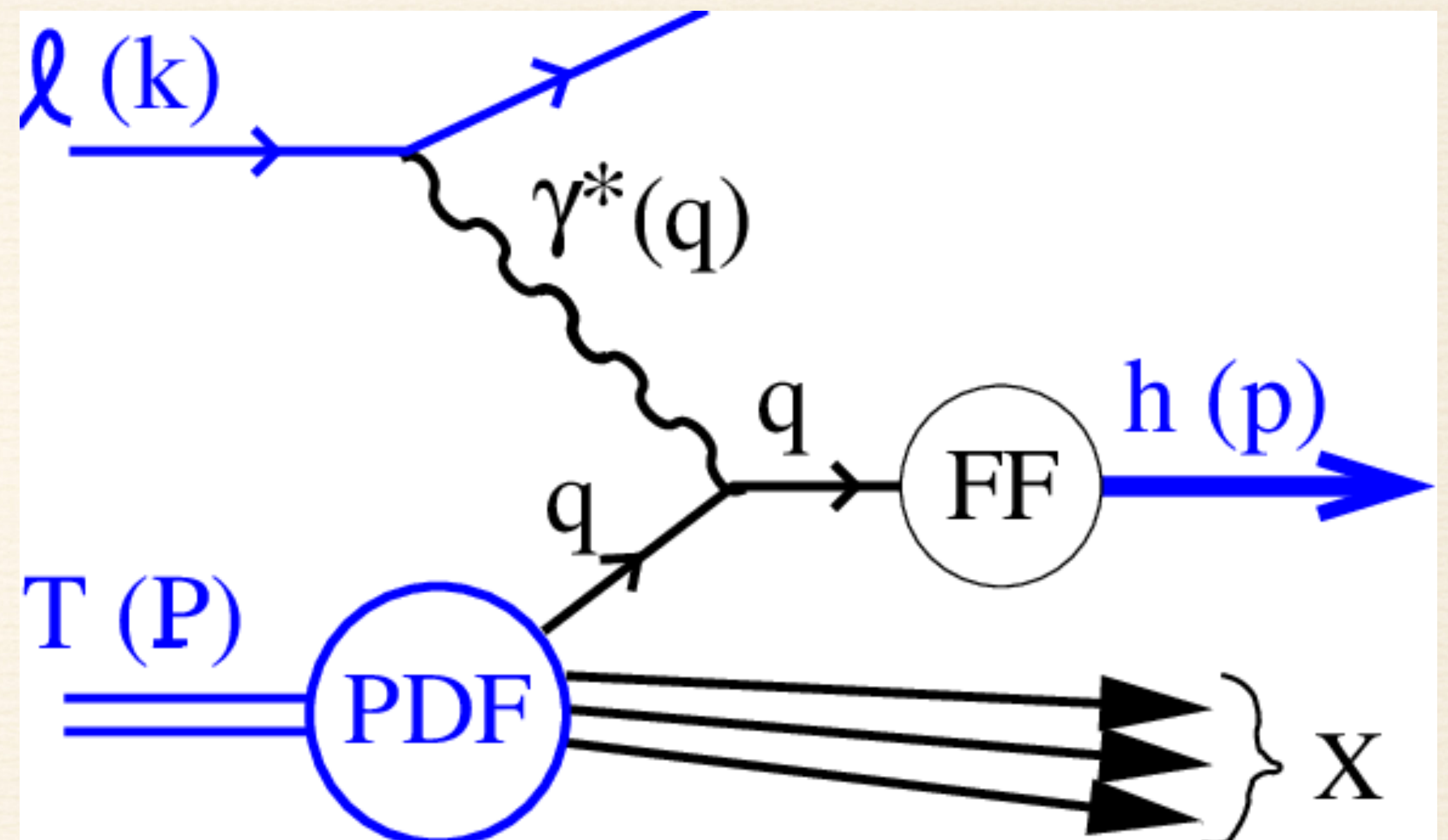


Matter To The Deepest
Katowice 2025

Semi-Inclusive-Deep-Inelastic-Scattering process



- ❖ *DIS + one final state hadron observed*
- ❖ *Important tool for extracting PDF's, TMD's*
- ❖ *3D tomography of the proton*
- ❖ *EIC*



State of the art

Recent NNLO calculation of SIDIS cross-section:

Ravindran, Sven-Olaf Moch, Roman Lee, Saurav Goyal, Vaibhav Pathak, Narayan Rana '24

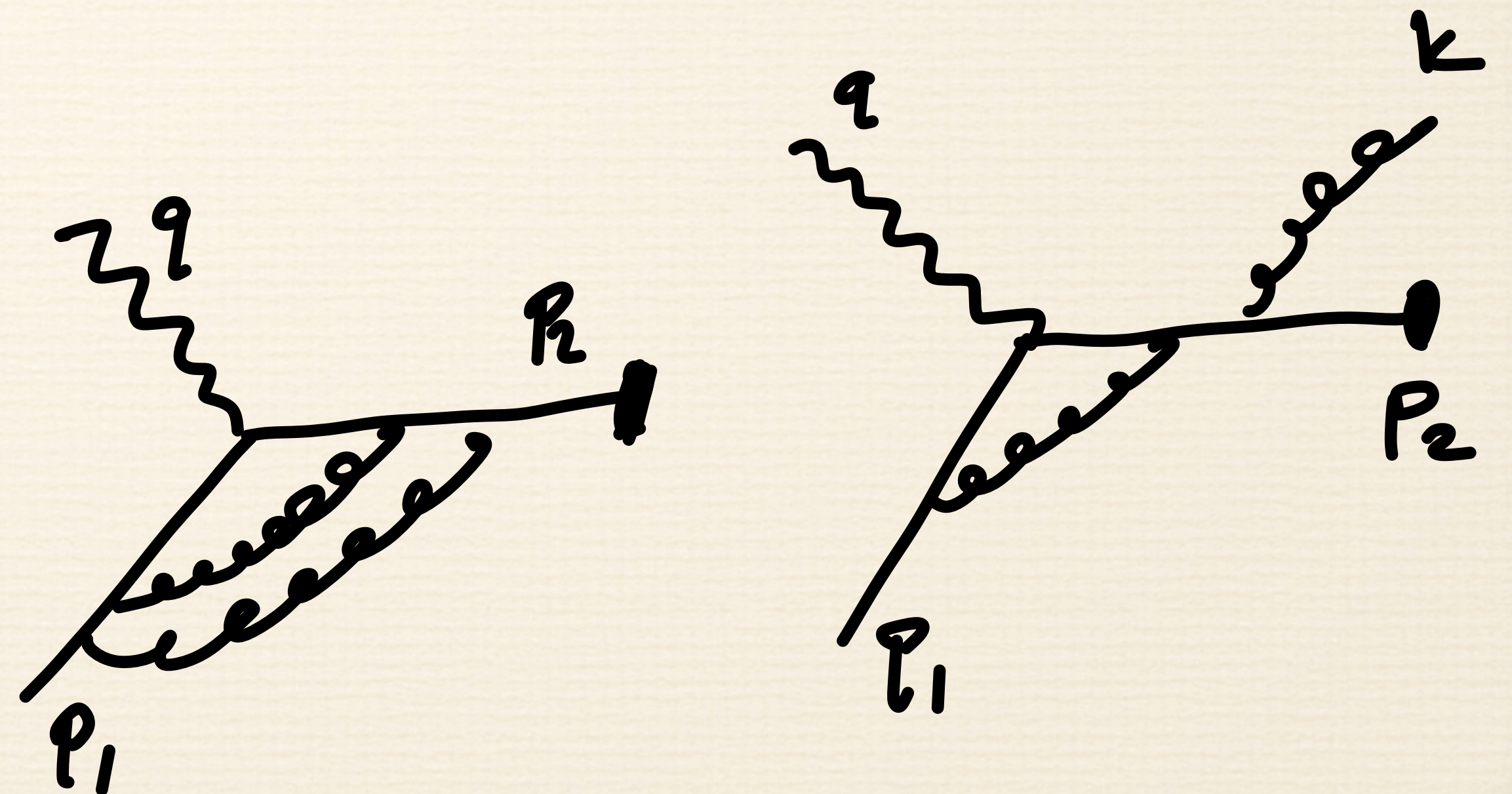
Gehrmann, Bonino, Stagnitto, Schönwald, Löchner '24

Is needed:

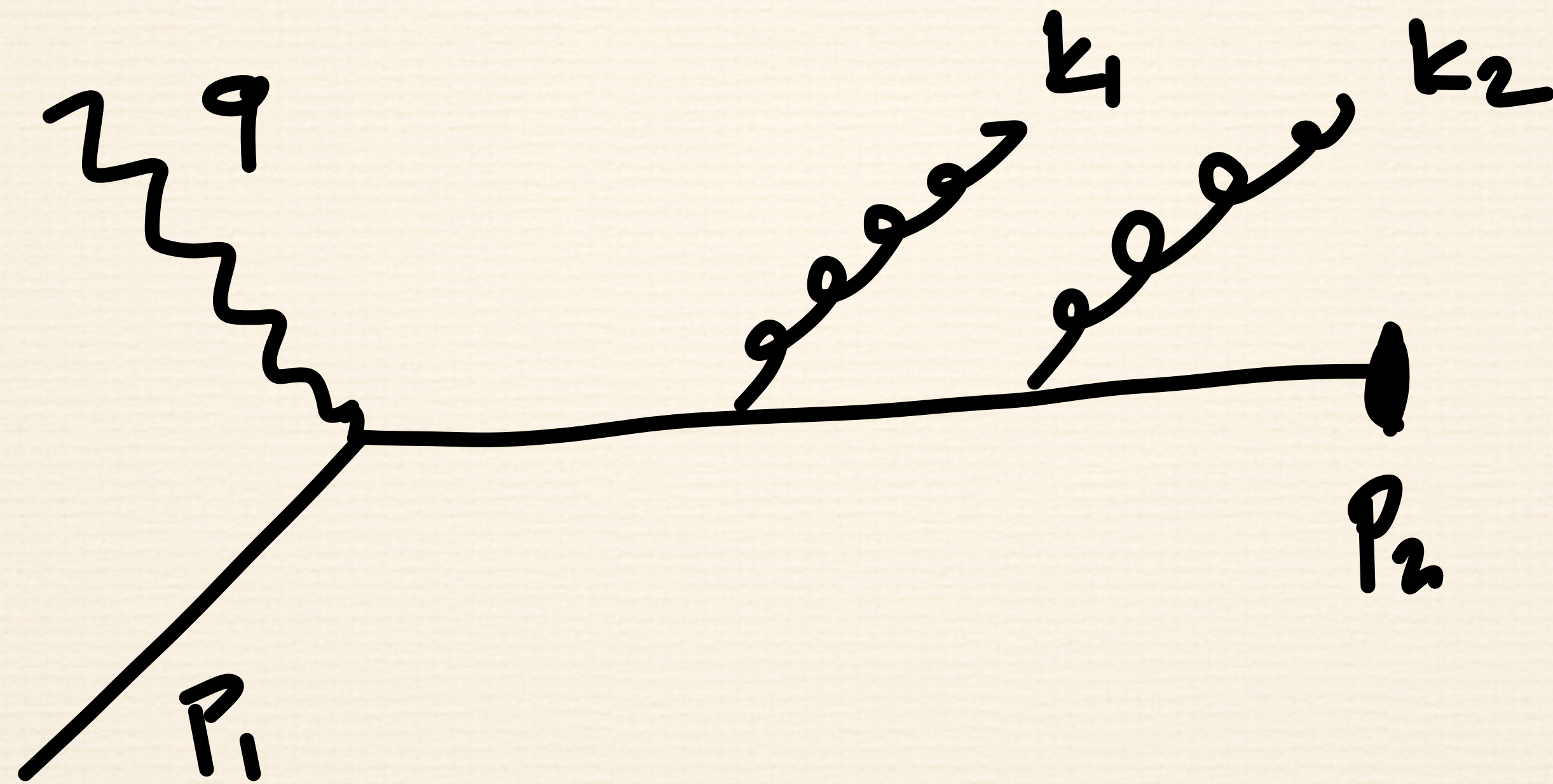
❖ 2 loop virtual $\rightarrow$ Long Known

❖ 1 loop + 1 real

❖ 2 real $\rightarrow$ Our Focus



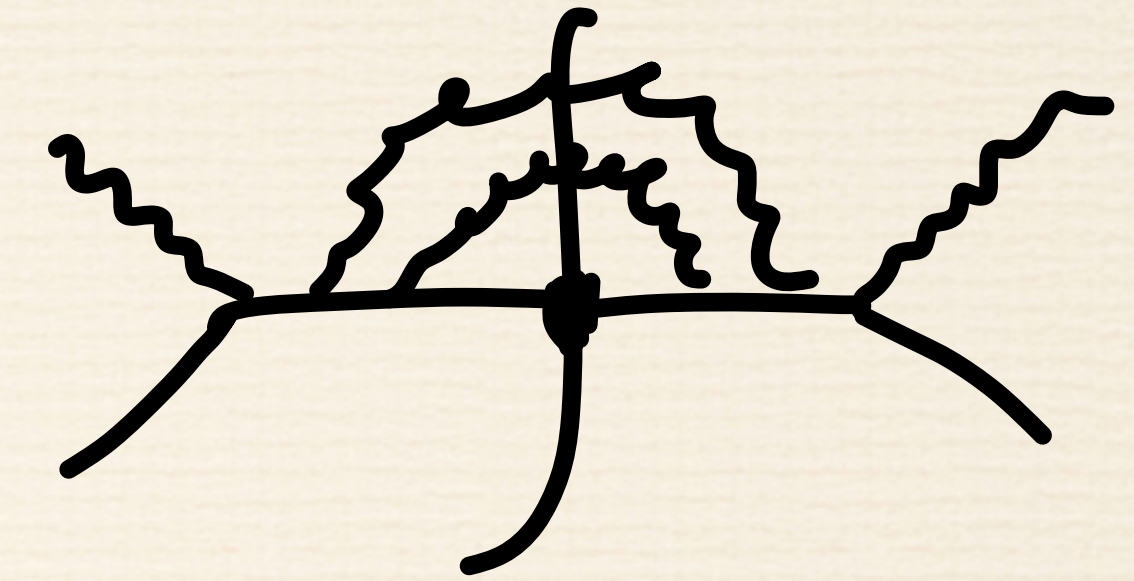
Goal: Calculate Phase space integrals



Reverse Unitarity

Anastasiou, Melnikov 2002

Real Emission Diagrams $\rightarrow$ Loop Diagrams



Reason: Apply IBP identities

Multi-loop Feynman Integrals in dim reg are convergent



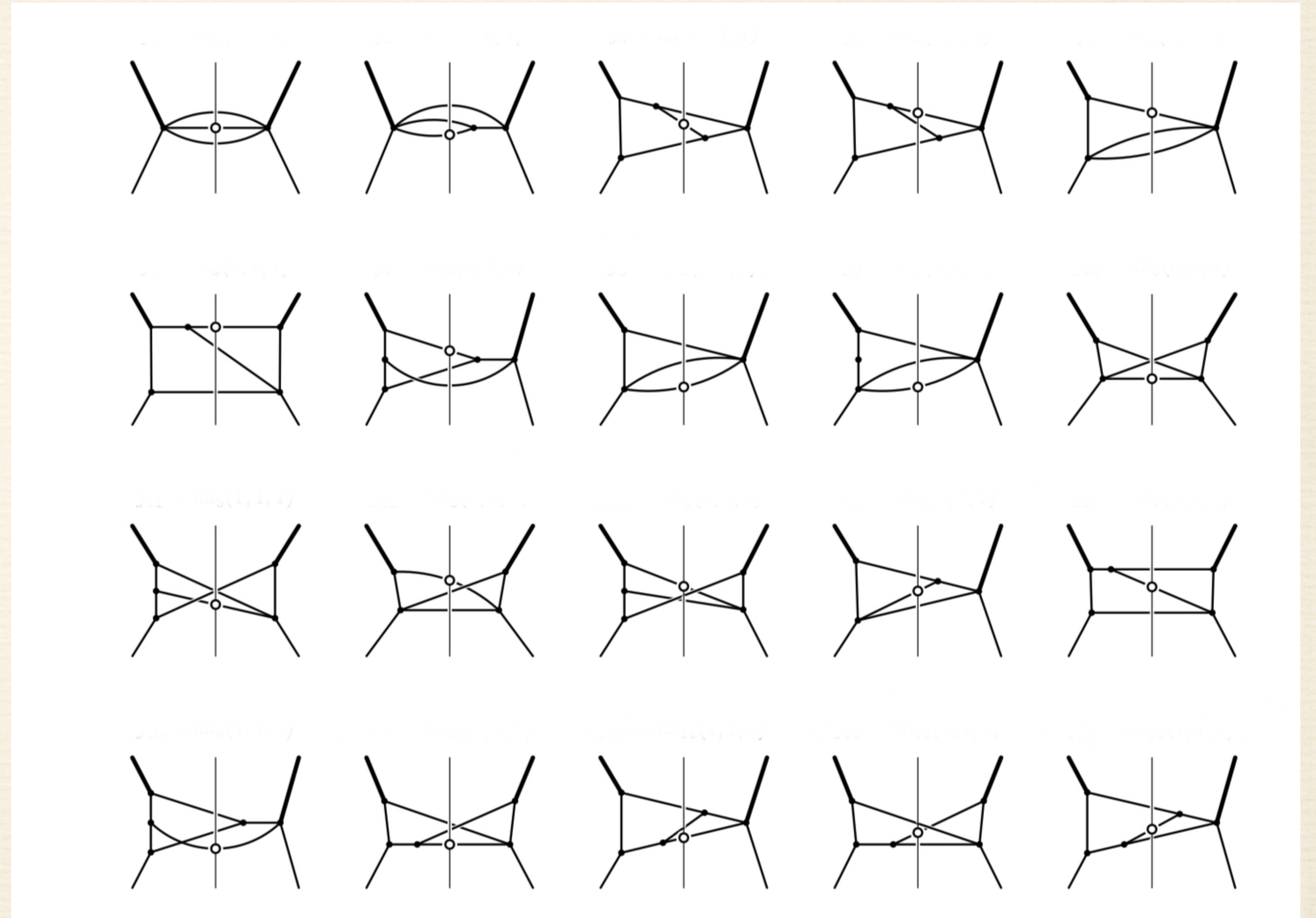
Boundary terms = 0 $\rightarrow$ Relation among integrals



Master Integrals

20 Master Integrals

$$\int dPS_3 \frac{1}{D_1^{n_1} D_2^{n_2} D_3^{n_3}} \delta(z' - \frac{p_1 \cdot p_2}{p_1 \cdot q})$$



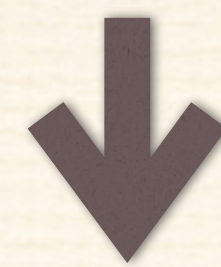
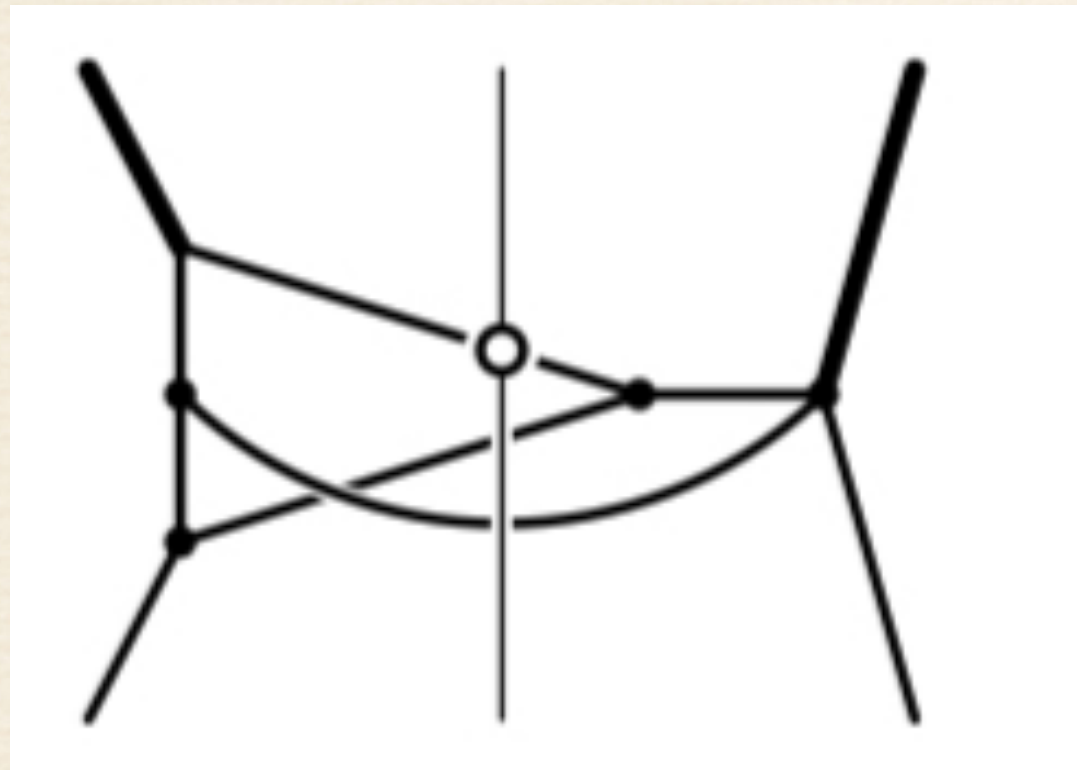
Methods

❖ DE Method (Canonical form, Iterative solution) Henn 2013

❖ Radial - Angular Decomposition (RAD) → Focus

Demonstration: Example

$$\int dPS_3 \frac{1}{(k_1 - p_1)^2 (k_2 - p_1 - q)^2 (k_1 + k_2 - p_1)^2} \delta\left(z' - \frac{p_1 \cdot p_2}{p_1 \cdot q}\right)$$

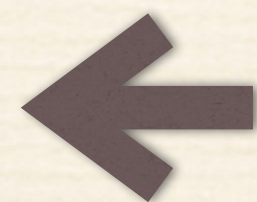


Spherical coordinates

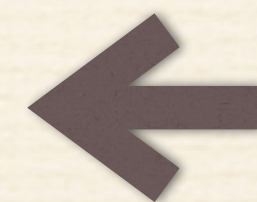
$$\int dPS'_{3,r} \frac{1}{r} \int dPS'_{3,\text{ang}} \frac{1}{(a + b \cos \theta) (A + B \cos \theta + C \sin \theta \cos \phi)}$$



Radial



MB



Angular



Angular Integrals

Somogyi '11

Lyubovitskij, Wunder, Zhevlakov '21

General two denominator case in Neerven parametrization:

$$\int d\Omega_{d-1} \frac{1}{(a + b\cos\theta_1)^j (A + B\cos\theta_1 + C\sin\theta_1\cos\theta_2)^l}$$

General two denominator case in Somogyi parametrization:

$$\frac{1}{a^j A^l} \int d\Omega_{d-1} \frac{1}{(p_1 \cdot q)^j (p_2 \cdot q)^l}$$

We categorize it

Massless: $a^2 = b^2, A^2 = B^2$

Single massive: $a^2 = b^2$ or $A^2 = B^2 + C^2$

Double massive: $a^2 \neq b^2, A^2 \neq B^2 + C^2$

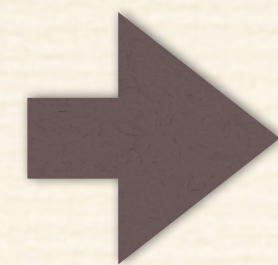
$$p_1 = (1, 0_{d-2}, -\frac{b}{a}) \quad , \quad p_2 = (1, 0_{d-3}, -\frac{C}{A}, -\frac{B}{A})$$

$$q = (1, \dots, \sin\theta_1 \cos\theta_2, \cos\theta_1)$$

Mellin-Barnes Technique

Taushif Ahmed, Syed Mehedi Hasan, **AR**, '24, arxiv: 2410.18886v1

Angular Integral
Three denominator massless



Mellin-Barnes representation

$$\int d\Omega_{d-1}(q) \frac{1}{(p_1 \cdot q)^{j_1} (p_2 \cdot q)^{j_2} (p_3 \cdot q)^{j_3}}$$

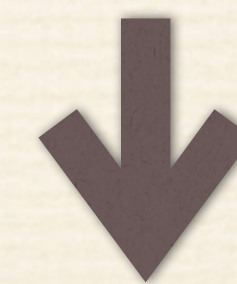
$$p_1^\mu = (1, 0, 0, 1)$$

$$p_2^\mu = (1, 0, \sin x_2^{(1)}, \cos x_2^{(1)})$$

$$p_3^\mu = (1, \sin x_3^{(2)} \sin x_3^{(1)}, \cos x_3^{(2)} \sin(x_3^{(1)}), \cos x_3^{(1)})$$

$$q^\mu = (1, \cos \theta_3 \sin \theta_2 \sin \theta_1, \cos \theta_2 \sin \theta_1, \cos \theta_1)$$

$$\begin{aligned} & 2^{2-j-k-l-2\epsilon} \pi^{1-\epsilon} \frac{1}{\Gamma(j)\Gamma(k)\Gamma(l)\Gamma(2-j-k-l-2\epsilon)} \\ & \times \int_{-i\infty}^{+i\infty} \frac{dz_{12} dz_{13} dz_{23}}{(2\pi i)^3} \Gamma(-z_{12}) \Gamma(-z_{13}) \Gamma(-z_{23}) \Gamma(j+z_{12}+z_{13}) \Gamma(k+z_{12}+z_{23}) \\ & \times \Gamma(l+z_{13}+z_{23}) \Gamma(1-j-k-l-\epsilon-z_{12}-z_{13}-z_{23}) (u_{12})^{z_{12}} (u_{13})^{z_{13}} (u_{23})^{z_{23}} \end{aligned}$$



Haug, Wunder '24

Expand in ϵ + solving the coefficients

Mellin-Barnes Technique

Four denominator case:

Taushif Ahmed, Syed Mehedi Hasan, **AR**, '25, [arxiv:2508.15952v1](#)

Haug, A.Smirnov, Wunder '25

Mellin-Barnes Technique

Gamma functions $\rightarrow$ Beta functions

$$B(a, b) = \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)}$$

Real integral representations:

$$B(a, b) = \int_0^1 dx x^{a-1} (1-x)^{b-1}, \quad \mathcal{R}(a), \mathcal{R}(b) > 0$$

$$B(a, b) = \int_0^\infty dx x^{a-1} (1+x)^{-a-b}, \quad \mathcal{R}(a), \mathcal{R}(b) > 0$$

Mellin-Barnes Technique

we get:

$$\int_0^1 \left(\prod_{k=1}^K dx_k \right) R_0(x, u) \int_{-i\infty}^{+i\infty} \prod_{l=1}^L \frac{dz_l}{2\pi i} \left[R_l(x, u) \right]^{z_l}$$

Mellin-Barnes Technique

Use the formula:

$$\int_{-i\infty+z_0}^{+i\infty+z_0} \frac{dz}{2\pi i} A^z = \delta(1-A), \quad A > 0$$

Then:

$$\int_0^1 \left(\prod_{k=1}^K dx_k \right) R_0(x, u) \prod_{l=1}^L \delta \left[1 - R_l(x, u) \right]$$

Mellin-Barnes Technique

Real Integrations over rational functions $\rightarrow$ Multiple Polylogarithms (MPLs):

$$G(a_1, \dots, a_n; z) = \int_0^z \frac{dt}{t - a_1} G(a_2, \dots, a_n; t)$$

Goal: Combining Radial + Angular

Result in terms of GPLs or MPLs



Angular part GPLs manipulation → Pole Subtraction Method



Singularity Resolution

Angular Part: Singularity Resolution

Why can we expand the angular part?

- ❖ Angular part → Collinear Singularities
- ❖ Parametric part → Soft Singularities

Conjecture:

Angular part soft singularity → All order Resummation + Factorization

Conjecture Verification → NNLO SIDIS

Example

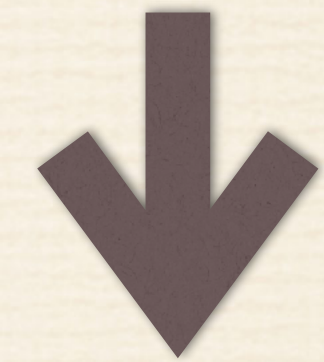
Angular Part :

$$\begin{aligned}\Omega'_z = & \frac{2\pi}{\varepsilon z} (z' + z - zz') + \frac{\pi}{z} (z' + z - zz') \ln \left(\frac{z}{z' + z - zz'} \right) \\ & + \frac{\pi\varepsilon}{4z} (z' + z - zz') \left\{ 2\text{Li}_2 \left(\frac{z}{zz' - z' - z} + 1 \right) + \ln^2 \left(\frac{z}{-zz' + z' + z} \right) \right\} \\ & + \frac{\varepsilon^2}{24z} \pi (z' + z - zz') \left\{ -6\text{Li}_3 \left(\frac{z}{-z'z + z + z'} \right) - 6\text{Li}_3 \left(\frac{z}{zz' - z' - z} + 1 \right) \right. \\ & + \ln^3 \left(\frac{z}{-zz' + z' + z} \right) - 3\ln \left(\frac{z}{zz' - z' - z} + 1 \right) \ln^2 \left(\frac{z}{-zz' + z' + z} \right) \\ & \left. + \pi^2 \ln \left(\frac{z}{-zz' + z' + z} \right) + 6\zeta(3) \right\} + \frac{\pi\varepsilon^3}{8z} (zz' - z' - z) \text{Li}_4 \left(\frac{zz' - z'}{z} \right) + O(\varepsilon^4)\end{aligned}$$

Extract



$$z^{-1+\frac{\varepsilon}{2}}$$



Push to Radial Part

Radial Part: Singularity Resolution

Consider $g(z)$ singular at $z = b$. Then separate the singular part $S(z)$ and the finite part $F(z)$ as:

$$g(z) = S(z)F(z)$$

Next we define:

$$\tilde{F} \equiv \lim_{z \rightarrow b} (F(z))$$

So we can separate the singularity as:

$$\begin{aligned} \int_a^b g(z) dz &= \int_a^b [S(z)F(z) - \tilde{F}S(z)] dz \\ &\quad + \int_a^b \tilde{F}S(z) dz \end{aligned}$$

The integrand $S(z)F(z) - \tilde{F}S(z)$ is non-singular and $\int_a^b \tilde{F}S(z) dz$ can be integrated easily.

Results

- ❖ Fully analytic form in GPLs
- ❖ One-fold numerical integration over GPLs

Verification through DE method $\rightarrow$ Perfect numerical agreement

Up to $\mathcal{E}^2$ $\rightarrow$ Useful for N^3LO SIDIS

Advantages of the RAD

- ❖ Angular Part → Universal Structure → Process Independent
- ❖ Transparent Singularity Structure
- ❖ It can handle Massive cases

Conclusion & Outlook

- ❖ Solution of NNLO 2 real SIDIS MI
- ❖ DE and RAD methods
- ❖ MB Technique for the Angular Part
- ❖ Transparent singularity structure in RAD
- ❖ RAD can handle Massive Cases

Thank you for listening!