

Precyzyjne testy Modelu Standardowego i jego rozszerzeń

Janusz Gluza, Uniwersytet Śląski w Katowicach

XLIX Zjazd Fizyków i Fizyczek Polskich

9.09.2025, Katowice

'Precise measurements of known particles and interactions
are just as important as finding new particles'

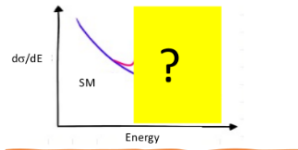
– Fabiola Gianotti



NARODOWE CENTRUM NAUKI



MS potwierdzone, dwa punkty widzenia



S. Dawson, BNL

Jan 14, 2025

This is sad scenario where there is no intermediate scale physics

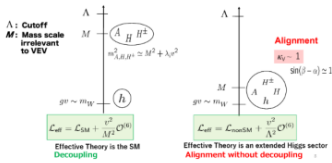
The Physics Landscape

We are in a fascinating situation: where to look and what will we find?

For the first time since Fermi theory, WE HAVE NO SCALE

The next facility must be versatile with as broad and powerful reach as possible, as there is **no precise target**

➔ **more Sensitivity, more Precision, more Energy**



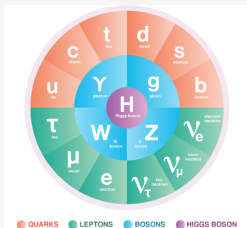
plot by Kanemura

A. Blondel, Epiphany 2021

Precyzyjne przewidywania SM stanowią podstawę do konfrontacji z każdym, szczególnie małym niestandardowym efektem, który poszukujemy!

PRECYZJA

Model Standardowy



"Whoever has only a hammer will see nothing but nails."

Serial dokumentalny BBC, [Precision: The Measure of All Things](#)

Stałe fizyczne, PDG2025 (Particle Data Group)

Quantity	Symbol, equation	Value	Uncertainty (ppb)
speed of light in vacuum	c	$299\,792\,458\text{ m s}^{-1}$	exact
Planck constant	h	$6.626\,070\,15 \times 10^{-34}\text{ J s}$ (or J/Hz) [§]	exact
Planck constant, reduced	$\hbar \equiv h/2\pi$	$1.054\,571\,817 \dots \times 10^{-34}\text{ J s}$ $= 6.582\,119\,569 \dots \times 10^{-22}\text{ MeV s}$	exact* exact*
electron charge magnitude	e	$1.602\,176\,634 \times 10^{-19}\text{ C}$	exact
conversion constant	$\hbar c$	$197.326\,980\,4 \dots \text{ MeV fm}$	exact*
conversion constant	$(\hbar c)^2$	$0.389\,379\,372\,1 \dots \text{ GeV}^2\text{ mbarn}$	exact*
electron mass	m_e	$0.510\,998\,950\,00(15)\text{ MeV}/c^2 = 9.109\,383\,7015(28) \times 10^{-31}\text{ kg}$	0.30
proton mass	m_p	$938.272\,088\,16(29)\text{ MeV}/c^2 = 1.672\,621\,923\,69(51) \times 10^{-27}\text{ kg}$ $= 1.007\,276\,466\,621(53)\text{ u} = 1836.152\,673\,43(11)\text{ } m_e$	0.31 0.053, 0.060
neutron mass	m_n	$939.565\,420\,52(54)\text{ MeV}/c^2 = 1.008\,664\,915\,95(49)\text{ u}$	0.57, 0.48
deuteron mass	m_d	$1875.612\,942\,57(57)\text{ MeV}/c^2$	0.30
unified atomic mass unit**	$u = (\text{mass } ^{12}\text{C atom})/12$	$931.494\,102\,42(28)\text{ MeV}/c^2 = 1.660\,539\,066\,60(50) \times 10^{-27}\text{ kg}$	0.30
fine-structure constant	$\alpha = e^2/4\pi\epsilon_0\hbar c$	$7.297\,352\,5693(11) \times 10^{-3} = 1/137.035\,999\,084(21)$ [†]	0.15
classical electron radius	$r_e = e^2/4\pi\epsilon_0 m_e c^2$	$2.817\,940\,3262(13) \times 10^{-15}\text{ m}$	0.45
$(e^- \text{ Compton wavelength})/2\pi$	$\lambda_e = \hbar/m_e c = r_e \alpha^{-1}$	$3.861\,592\,6796(12) \times 10^{-13}\text{ m}$	0.30
Bohr radius ($m_{\text{nucleus}} = \infty$)	$a_\infty = 4\pi\epsilon_0\hbar^2/m_e e^2 = r_e \alpha^{-2}$	$0.529\,177\,210\,903(80) \times 10^{-10}\text{ m}$	0.15
wavelength of 1 eV/c particle	$\hbar c/(1\text{ eV})$	$1.239\,841\,984 \dots \times 10^{-6}\text{ m}$	exact*
Rydberg energy	$\hbar c R_\infty = m_e e^4/2(4\pi\epsilon_0)^2\hbar^2 = m_e c^2 \alpha^2/2$	$13.605\,693\,122\,994(26)\text{ eV}$	1.9×10^{-3}
Thomson cross section	$\sigma_T = 8\pi r_e^2/3$	$0.665\,245\,873\,21(60)\text{ barn}$	0.91

^aUpdated values reflecting the “2022 CODATA recommended values”, released online in May 2024, will be incorporated in future editions of the *Review*.

[§]CODATA recommends that the unit be J/Hz to stress that in $h = E/\nu$ the frequency ν is in cycles/sec (Hz), not radians/sec.

*These are calculated from exact values and are exact to the number of places given (*i.e.* no rounding).

**The molar mass of ^{12}C is $11.999\,999\,9958(36)\text{ g}$.

[†]CODATA recommended value at $Q^2 = 0$. At $Q^2 \approx m_W^2$ the value is $\sim 1/128$. A world average of the latest data yields $\alpha^{-1} = 137.035999178(8)$. ^{‡‡}

[‡]Absolute laboratory measurements of G_N have been made only on scales of about 1 cm to 1 m.

Ku nowej definicji kilograma

Janusz Gluza, Agnieszka Grzanka, Agnieszka Pleban

Instytut Fizyki, Uniwersytet Śląski, Katowice

Towards a redefinition of the kilogram

Abstract: Recent attempts at a redefinition of the kilogram are shortly described.

1. Wstęp

Wszystkie podstawowe jednostki miar układu SI poza kilogramem zdefiniowane są w sposób uniwersalny, natomiast odzwierciedlenie masy wymaga bezpośredniego porównania danej próbki z oryginałem lub jedną z jego oficjalnych kopii. W artykule opisujemy próby zdefiniowania także jednostki masy w sposób fundamentalny, za pomocą precyzyjnych doświadczeń, opartych m.in. na efektach kwantowych. Fizycy pracujący nad tym zagadnieniem mają nadzieję, że ich eksperymenty będzie można w przyszłości powtórzyć, tak by odzwierciedlenie kilograma w precyzyjnym laboratorium fizycznym nie stanowiło problemu. Czy zmiana definicji kilograma nastąpi już w roku 2007, jak sugerują w swoich pracach niektórzy autorzy? Niekoniecznie, ale chyba nieuchronnie nadchodzi czas takiej zmiany.

2. Trochę historii i stan obecny

Król Ludwik XVI w 1790 r. zobowiązał francuskich uczonych do stworzenia spójnego systemu jednostek wag i miar – systemu metrycznego. Już 19 marca 1791 r. Francuska Akademia Nauk przyjęła definicję jednostki masy opartej na ściśle określonej objętości (1 dm³, czyli praktycznie jednego litra) destylowanej wody w temperaturze, w której ma ona największą gęstość, tj. ok. 4 °C. Za datę powstania międzynarodowego układu miar SI (Système International d'Unités) opartego na systemie dziesiętnym należy jednak uznać 22 czerwca 1799 r., kiedy to w Archives de la République w Paryżu zdeponowano dwa platynowe wzorce: metra oraz kilograma (tzw. Kilogramme des Archives). Ważnym krokiem w historii SI było utworzenie w 1875 r. Międzynarodowego Biura Miar i Wag (BIPM, Bureau International des Poids et Mesures) w Sèvres pod Paryżem.

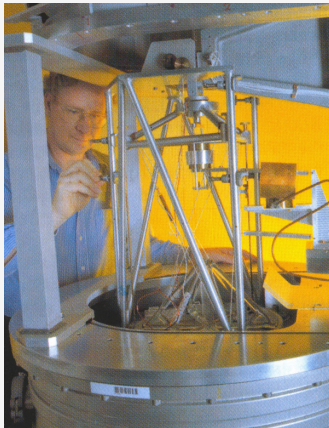
Kilogramme des Archives miał kształt walca i wdo-

wności na czynniki zewnętrzne. Dlatego przyjęto radę francuskiej sekcji międzynarodowej komisji metra CIM (Commission Internationale du Mètre), aby użyć wytworzonego przez francuskiego chemika Henriego Sainte-Claire Deville'a i jego ucznia Henriego Debray stopu platyny z irydem. Stop ten był zdecydowanie twardszy od czystszej platyny i miał dodatkowe pożądane właściwości (większą odporność na korozję, większą gęstość, dobre przewodnictwo ciepła i prądu, mniejszą wrażliwość na pole magnetyczne). Masową produkcję stopu zajęła się firma Johnson Matthey (JM).

Pierwsze trzy prototypy kilograma zostały wykonane właśnie w fabryce JM i nadano im nazwy KI, KII, KIII. Prototypy te porównano w Obserwatorium Paryskim z Kilogramme des Archives i ostatecznie jako międzynarodowy wzorzec kilograma przyjęto prototyp KIII. Wzorzec ten został usankcjonowany uchwałą I Generalnej Konferencji Miar w 1889 r., a za jego symbol przyjęto literę *k* (tzw. Grand K). Jest to wałek o średnicy i wysokości po 39 milimetrów (rys. 1), wykonany w 90% z platyny; resztą stanowi iryd.



Największa rewolucja w pomiarach od czasu rewolucji francuskiej:
nowa definicja jednostki masy



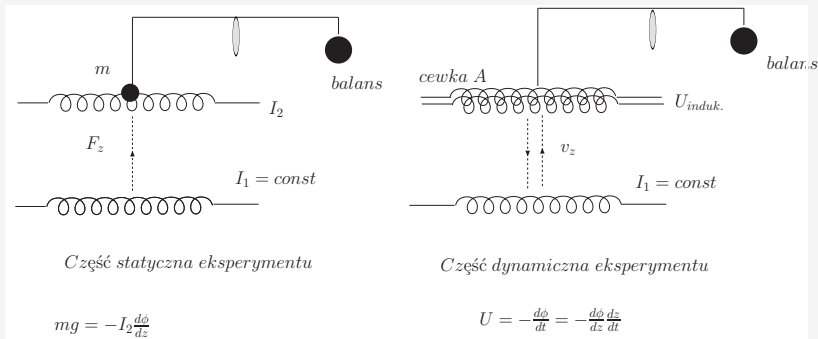
Definicje jednostki masy

- ▶ Nowa definicja, 2019:
"Kilogram, oznaczenie kg, jest to jednostka masy w układzie SI. Jest ona zdefiniowana poprzez przyjęcie ustalonej wartości liczbowej stałej Plancka h , wynoszącej $6.62607015 \cdot 10^{-34}$, wyrażonej w jednostce $J \cdot s$, która jest równa $kg \cdot m^2 \cdot s^{-1}$, przy czym metr i sekunda zdefiniowane są za pomocą c i f_{Cs} ."
- ▶ Czyli mamy związek:

$$kg \leftrightarrow h, m(c), s(f_{Cs})$$

.

Balans Watta: idea (Postępy fizyki, 2007)



Balans równoważy cewkę A

$$mgv = U_{\text{induk.}} I_2$$

U, I - kwantowe efekty (prąd tunelowy - stała Josphsona $\propto e/h$,
oporność - stała von Klitzinga $\propto h/e$).

How We're Redefining the kg

Filmik z NIST. <https://www.youtube.com/watch?v=Oo0jm1PPRu0>

specjalny budynek



Układ SI (od 1790) definiuje 7 podstawowych jednostek fizycznych

Wielkość	Jednostka	względna dokładność
Długość	m	1×10^{-12}
Masa	kg	1×10^{-8}
Czas	s	3×10^{-15}
Prąd	A	4×10^{-8}
Temperatura	K	3×10^{-7}
Światłość	kandela	1×10^{-4}
Miara substancji	mol	8×10^{-8}

Dokładność: błąd sekundy raz na 20 mln lat - ważne np. w GPS, komórki, internet.

<https://www.nature.com/articles/s41586-021-03253-4> - "Frequency ratio measurements at 18-digit accuracy using an optical clock network" - Nature, 2021

Stała struktury subtelnej - α_{QED} - nowy wynik (2020)

REPORT

Measurement of the fine-structure constant as a test of the Standard Model

Richard H. Parker^{1,*}, Chengshai Ye^{1,2,3}, Weicheng Zheng¹, Brian Estey⁴, Holger Müller^{1,3,5}

* See all authors and affiliations

Science 13 Apr 2020
Vol. 368, Issue 6485, pp. 101-105
DOI: 10.1126/science.aba7706

Article | Published: 02 December 2020

Determination of the fine-structure constant with an accuracy of 81 parts per trillion

Léo Morel, Zhibin Yao, Pierre Cladé & Salda Guellati-Khélifa

Nature 588, 61-65(2020) | Cite this article

6367 Accesses | 1 Citations | 300 Altmetric | Metrics

$$\alpha^{-1}(Cs) = 137.035\,999\,046(27)$$
$$\alpha^{-1}(Rb) = 137.035\,999\,206(11)$$
$$\alpha^{-1}(a_e) = 137.035\,999\,139(31)$$



Uwaga: (i) nowy wynik już odchylenie od SM w kierunku jak dla $(g-2)_\mu$, (ii) duża rozbieżność z Cs ($\sim 5.4\sigma$).

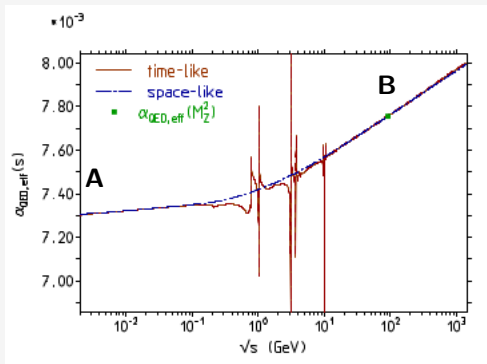
Guellati-Khélifa poprawiała eksperyment 22 lata¹

<https://www.nature.com/articles/s41586-020-2964-7> [02 December 2020]

$$R_\infty = \frac{\alpha^2 m_e c}{4\pi\hbar} \longrightarrow \alpha^2 = \frac{R_\infty}{c} \times \frac{h}{m_e} \longrightarrow \frac{h}{m_e} = \frac{u}{m_e} \frac{M_X}{u} \frac{h}{M_X}$$

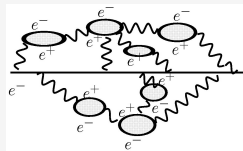
¹ Jaki najdłuższy eksperyment w historii?

Fizyka wysokich energii zmienia obraz, $\alpha_{QED}(s)$, polaryzacja próżni



A: $\alpha_{QED}(0) \simeq 1/137$

B: $\alpha_{QED}(M_Z^2) \simeq 1/128$



F. Jegerlehner, <http://dx.doi.org/10.23731/CYRM-2020-003.9>

Efektywna $\alpha(s)$ w funkcji fotonowej polaryzacji próżni $\Delta\alpha(s)$

$$\alpha(s) = \frac{\alpha}{1 - \Delta\alpha(s)} \Delta\alpha(s) = \Delta\alpha_{\text{lep}}(s) + \Delta\alpha_{\text{had}} + \Delta\alpha_{\text{top}}(s).$$

Fizyka cząstek elementarnych

gravitational constant [†]	G_N	$6.674\,30(15) \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$ $= 6.708\,83(15) \times 10^{-39} \hbar c \text{ (GeV}/c^2\text{)}^{-2}$	2.2×10^4 2.2×10^4
standard gravitational accel.	g_N	$9.806\,65 \text{ m s}^{-2}$	exact
Avogadro constant	N_A	$6.022\,140\,76 \times 10^{23} \text{ mol}^{-1}$	exact
Boltzmann constant	k	$1.380\,649 \times 10^{-23} \text{ J K}^{-1}$ $= 8.617\,333\,262 \dots \times 10^{-5} \text{ eV K}^{-1}$	exact exact*
molar volume, ideal gas at STP	$N_A k \text{ (273.15 K)/(101 325 Pa)}$	$22.413\,969\,54 \dots \times 10^{-3} \text{ m}^3 \text{ mol}^{-1}$	exact*
Wien displacement law constant	$b = \lambda_{\text{max}} T$	$2.897\,771\,955 \dots \times 10^{-3} \text{ m K}$	exact*
Stefan-Boltzmann constant	$\sigma = \pi^2 k^4 / 60 \hbar^3 c^2$	$5.670\,374\,419 \dots \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$	exact*
Fermi coupling constant ^{††}	$G_F / (\hbar c)^3$	$1.166\,378\,8(6) \times 10^{-5} \text{ GeV}^{-2}$	510
weak-mixing angle ^{††}	$\sin^2 \theta(M_Z) \text{ (}\overline{\text{MS}}\text{)}$	$0.231\,29(4)^{\dagger\dagger}$	1.7×10^5
$W^\pm$ boson mass	m_W	$80.3692(133) \text{ GeV}/c^2$ [¶]	1.7×10^5
Z^0 boson mass	m_Z	$91.1880(20) \text{ GeV}/c^2$	2.2×10^4
strong coupling constant	$\alpha_s(m_Z)$	$0.1180(9)$	7.6×10^6
$\pi = 3.141\,592\,653\,589\,793\,238 \dots$ $e = 2.718\,281\,828\,459\,045\,235 \dots$ $\gamma = 0.577\,215\,664\,901\,532\,860 \dots$			
1 in $\equiv 0.0254 \text{ m}$	1 G $\equiv 10^{-4} \text{ T}$	1 eV $= 1.602\,176\,634 \times 10^{-19} \text{ J (exact)}$	kT at 300 K $= [38.681\,727\,0718 \dots]^{-1} \text{ eV (exact*)}$
1 Å $\equiv 0.1 \text{ nm}$	1 dyne $\equiv 10^{-5} \text{ N}$	(1 kg) $c^2 = 5.609\,588\,603 \dots \times 10^{35} \text{ eV (exact*)}$	0 °C $\equiv 273.15 \text{ K}$
1 barn $\equiv 10^{-28} \text{ m}^2$	1 erg $\equiv 10^{-7} \text{ J}$	1 C $= 2.997\,924\,58 \times 10^9 \text{ esu}$	1 atmosphere $\equiv 760 \text{ Torr} \equiv 101\,325 \text{ Pa}$

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[‡]Absolute laboratory measurements of G_N have been made only on scales of about 1 cm to 1 m.

- ▶ (i) odkrycie mionu, J/Ψ
- (ii) $(g-2)_e$, $(g-2)_\mu$
- (iii) V-A, parzystość;

EKSPERYMENT $\rightarrow$ TEORIA

- ▶ (i) $\tau^\pm$ (tau lepton);
- (ii) Tevatron - odkrycie kwarku top;
- (iii) H^0 (skalarny boson (Higgs-Englert...))

TEORIA $\rightarrow$ EKSPERYMENT

ważne SM poprawki kwantowe! LEP, SLAC, LHC
M. Veltman (1977) ρ -parametr $\sim m_t^2, \ln(m_H^2)$;
 $\rightarrow$ Acta Physica Polonica B

- ▶ Neutrino (masy, mieszania, CP faza(y);
Super-K, Hyper-K, T2K, NOvA, Antares, KM3NeT, Juno, DUNE,
SNO+, Daya Bay, Double Chooz, RENO, ...

BOOK
Special volume of Acta Physica
Polonica B commemorating
Martinus Veltman

by Jachew, Stanisław; Jachew, Marek;
Praszkiewicz, Michał

Published by Acta Physica Polonica 2021

This memorial volume of Acta Physica Polonica B is a special tribute to the Nobel Laureate Martinus Veltman. Professor Veltman was a member of the APFB International Editorial Council, lecturer at the Cracow School of Theoretical Physics in 1977 and 1994, and the author of two of the most cited articles ever published in our journal: «Second Threshold in Weak Interactions» [Acta Phys. Pol. B 8, 475 (1977)] and «The Infrared-Ultraviolet Connection» [Acta Phys. Pol. B 12, 437 (1981)].

Przyszłe zderzacze; TEORIA $\leftrightarrow$ EKSPERYMENT

Aside: factor-of-2 improvements can matter!

Search for $K_L \rightarrow \pi\pi$

ANNALS OF PHYSICS: 8, 156-181 (1958)

Long-lived Neutral K Mesons*

M. BARDON, K. LANDE, AND L. M. LEDERMAN

*Columbia University, New York, New York, and Brookhaven
National Laboratories, Upton, New York*

AND

WILLIAM CHINOWSKY

Brookhaven National Laboratories, Upton, New York

set an upper limit $<0.6\%$ on the reactions

$< 0.6\%$

$$K_S^0 \rightarrow \begin{cases} \mu^+ + e^- \\ e^+ + e^- \\ \mu^+ + \mu^- \end{cases}$$

and on $K_S^0 \rightarrow \pi^+ + \pi^-$.

VOLUME 6, NUMBER 10

PHYSICAL REVIEW LETTERS

MAY 15, 1961

DECAY PROPERTIES OF K_S^0 MESONS*

D. NAGAI, E. O. OIKARINEN, N. I. PETROV, A. M. RUSANOVA, AND V. A. RUSAKOV
Joint Institute of Nuclear Research, Moscow, U.S.S.R.
(Received April 20, 1961)

Combining our data with those obtained in reference 7, we set an upper limit of 0.3% for the relative probability of the decay $K_S^0 \rightarrow \pi^- + \pi^+$. Our

$< 0.3\%$

"At that stage the search was terminated by administration of the Lab."

[Okun, hep-ph/0112031]

VOLUME 13, NUMBER 4

PHYSICAL REVIEW LETTERS

27 JULY 1964

EVIDENCE FOR THE 2π DECAY OF THE K_S^0 MESON*†

J. H. CHRISTENSEN, J. W. CRONIN,[‡] V. L. FITCH,[‡] AND R. TURLAY[§]
Princeton University, Princeton, New Jersey
(Received 10 July 1964)

$= 0.2 \pm 0.04\%$

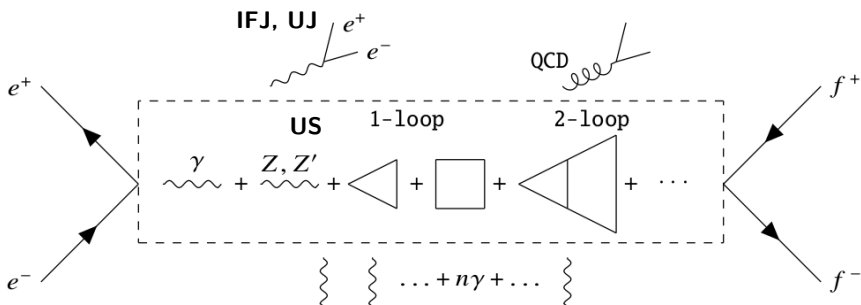
We would conclude therefore that K_S^0 decays to two pions with a branching ratio $R = (K_S^0 \rightarrow \pi^+ + \pi^-) / (K_S^0 \rightarrow \text{all charged modes}) = (2.0 \pm 0.4) \times 10^{-3}$ where the error is the standard deviation. As empha-

particles.us.edu.pl/czastki.us.edu.pl

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Theory and Phenomenology of Particle Physics, University of Szécsen research group

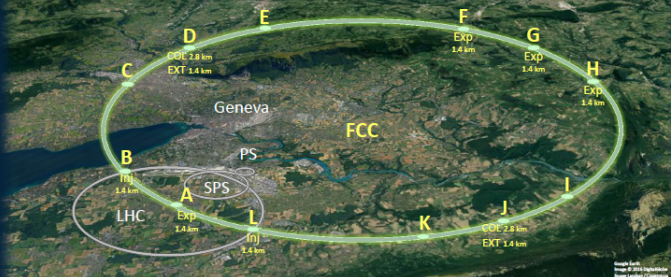
Generatory MC, teoria (rezonans bozonu Z), perturbacyjne poprawki kwantowe



Współpraca Kraków - Katowice, 2024-2029, grant NCN

The Future Circular Collider (FCC) study is an international collaboration aimed at designing the particle accelerator that will replace the LHC once it has completed its operational lifetime. The FCC will expand the current energy and luminosity frontiers in order to help answer the most fundamental questions in science: What is dark matter? Are there extra dimensions in the universe? Are there other forces in nature?

The FCC collaboration, hosted by CERN, is open to universities, research institutes and high-tech companies. A conceptual design will be delivered before the end of 2018, in time for the next update of the European Strategy for Particle Physics.



FCC-hh – A discovery machine

The 100 TeV proton-proton collider (FCC-hh) will have an energy seven times higher than the LHC. Such a collider will give access to the smallest scales and the most energetic phenomena in nature.

New fundamental forces and particles can be discovered, extending the reach for searching dark matter particles, supersymmetric partners of quarks and gluons, and possible substructure inside quarks. Billions of Higgs bosons and trillions of top quarks will be produced, creating new opportunities for the study of rare decays, flavor physics, and the mechanism of electroweak symmetry breaking.

The FCC-hh collider provides also the opportunity to push the exploration of the collective structure of matter at the most extreme density and temperature conditions to new frontiers through the study of heavy-ion collisions.

FCC-ee – A machine for precision

The second scenario of the FCC design study (FCC-ee) is a high-luminosity, high-precision electron-positron collider with center-of-mass collision energies between 90 and 350 GeV. Located in the same 100 km long tunnel as the FCC-hh it is considered a potential intermediate step towards the realization of the hadron facility, and complementary to it.

Clean experimental conditions give electron-positron colliders the capability to measure known particles with the highest precision.

FCC-ee would measure the properties of the Z, W, Higgs and top particles with unequal accuracy, offering the potential for discovering dark matter or heavy neutrinos. The FCC-ee could enable profound investigations of electroweak symmetry breaking and open a broad indirect search for new physics over several orders of magnitude in energy.

FCC-he – New opportunities

With the huge energy provided by the 50 TeV proton beam and the potential availability of an electron beam with energies of the order of 60 GeV, new horizons open up for the physics of deep inelastic electron-proton scattering.

The FCC-he collider would be both a high-precision Higgs factory and a powerful microscope to discover new particles. It would be the most accurate tool for studying quark-gluon interactions, possible substructure of matter and unprecedented measurements of strong and electroweak interaction phenomena. The hadron-electron collider is a unique complement to the exploration of nature at high energies within the FCC complex.

The FCC study explores three different scenarios: a hadron-hadron collider (FCC-hh), an electron-positron collider (FCC-ee), and a hadron-lepton (FCC-he) collider. The hadron-hadron collider defines the overall infrastructure for the FCC. With a target center-of-mass energy of 100 TeV, and 16-Tesla bending magnets, such a machine will have a circumference of 100 km.

Main parameters and geometrical aspects

	LHC	FCC
Circumference [km]	26.7	100
Dipole field [T]	8.33	16
Strategic sections	8 × 303 m	8 × 3000 m (2 × 4000 m)
Number of IPs	2+2	2+2
Injection energy [TeV]	0.45	3.3

FCC-hh compared with LHC and High-Luminosity LHC

	LHC	HL-LHC	FCC-hh baseline	FCC-hh ultimate
Energy at center of mass [TeV]	14	14	100	300
Bunch spacing [ns]	25	25	25	25
Number of bunches	2808	2808	10600	10600
Transverse emittance [nm]	0.25	0.25	0.44	0.44
Beam current [A]	0.584	1.12	0.5	0.5
Peak luminosity [$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$]	1.0	3.0	3.0	<30.0

FCC-ee compared with the Large Electron-Positron collider (LEP2)

The main center-of-mass operating points with strong physics interest for FCC-ee are 91 GeV (Z pole), 160 GeV (W pair production threshold), 240 GeV (Higgs resonance) and 350 GeV (t $\bar{t}$ threshold).

	LEP2	FCC-ee			
		Z	W	H	t
Energy at center of mass [GeV]	208	91	160	240	350
Bunch spacing [ns]	247/494	7.5	2.5	50	400
Number of bunches	4	30180	91500	5260	780
Emittance (horizontal) [nm]	22	0.2	0.09	0.26	0.61
Emittance (vertical) [nm]	250	1	1	1.2	2
Beam current [mA]	3.04	1450	152	30	6
Peak luminosity [per 2 IPs] [$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$]	0.012	207	90	19.1	5.1

Contacts and further information

FCC - FCC Office
fcc.office@cern.ch

EuroCircol - Prof. Carsten P. Welsch
carsten.welsch@cockcroft.ac.uk



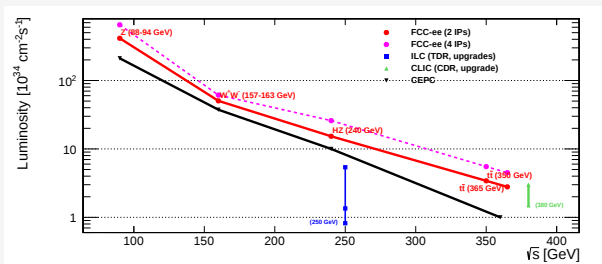
<http://fcc.web.cern.ch>

<http://www.eurocircol.eu>



The project has received funding from the European Union's Horizon research and innovation programme under grant agreement 101019754. The information contained herein is confidential and intended solely for the individual named. It is not to be distributed outside the named individual's organisation. The project is not responsible for any use that may be made of the information contained herein.

Motywacja do precyzyjnych badań: fabryki elektroślabych Z, W, H, t i zapachów



<https://arxiv.org/abs/2203.06520> (z podsumowania aktywności US Snowmass)

Phase	Run duration (years)	Center-of-mass Energies (GeV)	Integrated Luminosity (ab^{-1})	Event Statistics
FCC-ee-Z	4	88-94	150	$5 \cdot 10^{12}$ Z decays
FCC-ee-W	2	157-163	10	10^8 WW events
FCC-ee-H	3	240	5	10^6 ZH events 25k $WW \rightarrow H$
FCC-ee-tt	5	340-365	$0.2 \div 1.5$	10^6 $t\bar{t}$ even ts 200k ZH 50k $WW \rightarrow H$

FCC-hh, optymalne dla 84 TeV [pdf](#) (referat ML Mangano)

Higgs Factories

- The Higgs boson has a size/wavelength. What's inside?



Precision measurements are different ways of probing the "compositeness of the Higgs".



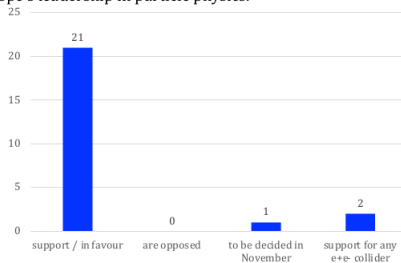
$$\lambda_h \approx 10^{-17} \text{ m}$$

$$\lambda_{10 \text{ TeV}} \approx 10^{-19} \text{ m}$$

Referat "National Inputs" Calina Alexa, <https://agenda.infn.it/event/44943/>

a) Which is the preferred next major/flagship collider project for CERN?

- **Broad consensus** among CERN Member States in support of the Future Circular Collider (FCC) as a key long-term project to maintain Europe's leadership in particle physics.



Support for FCC	Belgium, Czech Republic, Denmark, Estonia, Finland, France, Germany, Greece, Hungary, Israel, Italy, Norway, Poland, Portugal, Romania, Serbia, Slovak Republic, Spain, Sweden, Switzerland, United Kingdom
Opposed	None
To be finalized in November	Netherlands
Support for any e ⁺ e ⁻ collider	Austria, Bulgaria

Horizon Europe 2028-2034

16 July 2025, Commission President Ursula von der Leyen presented the Commission's proposal for the next Multiannual Financial Framework (MFF) 2028–34 and the European Competitiveness Fund (ECF), [link - pdf](#)

Based on orientation of the steering mechanism for the next MFF including the Competitiveness Coordination Tool, the Horizon Europe programme and the European Competitiveness Fund could finance a coherent sequence between research and innovation, demonstration, development and deployment, focusing efforts and funding, from the EU and national, public and private sector of 'moonshots' projects with a strong scientific component, boosting EU-wide value creation and strategic autonomy (see examples below) .

Possible 'moonshots':

- investing in the European Organization for Nuclear Research's (CERN) Future Circular Collider, alongside other CERN's participating countries. The objective is to maintain Europe's leadership in particle physics research. The funding (up to 20% of the overall cost) could come from Horizon Europe.
- Developing Smart and Clean Aviation and European leadership in the next generation CO₂-free aircraft and automated air traffic management: It would require a partnership with industry, together with a strong scientific and engineering capacity, supported by Horizon Europe, but also a robust industrial deployment component from the Competitiveness Fund.

Moonshots

Objective: ambitious technology driven projects that boost the EU's strategic autonomy through research, development and deployment.

▶ the future circular collider ▶ clean aviation ▶
quantum computing ▶ next-generation AI ▶ data
sovereignty ▶ automated transport and mobility
▶ regenerative therapies ▶ fusion energy ▶ space
economy ▶ zero water pollution ▶ ocean observation

Jakiej precyzji oczekujemy na poziomie Tera-Z w FCC-ee?

→ 10^{12} rozpadów bozonu Z w rezonansie e^+e^- ?

⇒ 1-2 rzędy wielkości poprawa w wyznaczeniu mierzonych obserwabli!

Aby dojść do eksperymentalnej precyzji, musimy poprawić obliczenia!

Expected precision in 2040

Conclusion of the 2018 Workshop

J. Gluza

"We anticipate that, at the beginning of the FCC-ee campaign of precision measurements, the theory will be precise enough not to limit their physics interpretation. This statement is however conditional to sufficiently strong support by the physics community and the funding agencies, including strong training programmes".

Numerical evaluation with three-loops calculations:

arXiv:1901.02648

	$\delta\Gamma_Z$ [MeV]	δR_l [10^{-4}]	δR_b [10^{-5}]	$\delta \sin_{eff}^{2,l} \theta$ [10^{-6}]
Present EWPO theoretical uncertainties				
EXP-2018	2.3	250	66	160
TH-2018	0.4	60	10	45
EWPO theoretical uncertainties when FCC-ee will start				
EXP-FCC-ee	0.4 0.025	10	2 ÷ 6	6 3
TH-FCC-ee	0.07	7	3	7

0.5 → 0.4
Pięć lat!

- 500 person-years needed over 20 years – **Recognized as strategic priority.**

President Clinton's talk at MIT's 1998 Commencement



President William Jefferson Clinton—1998 MIT Commencement

Just yesterday in Japan, physicists announced a discovery that tiny neutrinos have mass. Now, that may not mean much to most Americans, but **it may change our most fundamental theories -- from the nature of the smallest subatomic particles to how the universe itself works, and indeed how it expands.**

.....

The larger issue is that these kinds of findings have **implications that are not limited to the laboratory.** They affect the whole of society -- not only our economy, but our very view of life, our understanding of our relations with others, and **our place in time.**

Slajd: Kajita, Neutrino'2022, <https://www.youtube.com/watch?v=9LheUWrXUHU>

Niestandardowe efekty neutrin możemy też bardzo ładnie badać poprzez precyzyjne pomiary i obliczenia w FCC-ee.

N_ν

QED deconvolution

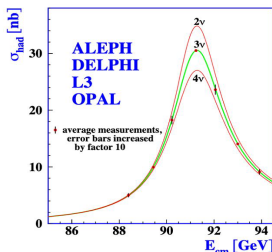


Reminder: Measuring N_ν at LEP

- The Z lineshape determination**

Phys. Rep. 427 (2006) 257

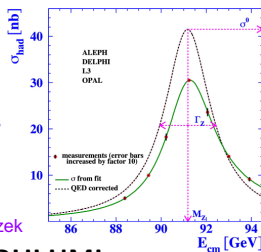
- Measure hadronic and leptonic cross sections (σ_{had} and $\sigma_{\ell\ell}$) as a function of E_{cm} ($\sqrt{s}$)



Fit to a Breit-Wigner shape
folded with QED ISR effects

Get m_Z , Γ_Z and σ^0
Get also $R_t^0 = \sigma_{\text{had}}^0 / \sigma_{\ell\ell}^0$

Staszek



KKMC, BHLUMI

- The peak cross section σ^0 is very sensitive to N_ν**

- The smaller the peak cross section, the larger the number of light neutrino active species



Reminder: Measuring N_ν at LEP

- What was done in practice to extract N_ν :

- Total Z decay width : $\Gamma_Z = \Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{\text{had}} + N_\nu \Gamma_{\nu\nu}$. $R_\ell^0 = \frac{\Gamma_{\text{had}}}{\Gamma_{\ell\ell}}$
- Divide by $\Gamma_{\ell\ell}$: $\Gamma_Z / \Gamma_{\ell\ell} = 3 + \delta_\tau + R_\ell^0 + N_\nu (\Gamma_{\nu\nu} / \Gamma_{\ell\ell})$
 - δ_τ is a small phase-space correction due to the finite τ mass
 - $(\Gamma_{\nu\nu} / \Gamma_{\ell\ell})$ rather immune to SM parameters (m_{top} , m_H , ...): taken from SM
 - $\Gamma_Z / \Gamma_{\ell\ell}$ taken from Breit-Wigner peak expression $\sigma_{\text{had}}^0 = \frac{12\pi}{m_Z^2} \frac{\Gamma_{ee} \Gamma_{\text{had}}}{\Gamma_Z^2}$
- Solve for N_ν :

$$N_\nu \left(\frac{\Gamma_{\nu\nu}}{\Gamma_{\ell\ell}} \right)_{\text{SM}} = \left(\frac{12\pi}{m_Z^2} \frac{R_\ell^0}{\sigma_{\text{had}}^0} \right)^{\frac{1}{2}} - R_\ell^0 - 3 - \delta_\tau \approx -2.263 \cdot 10^{-3}$$

SM prediction:

= 1.99125 ± 0.00083 in 2005

= 1.99060 ± 0.00021 in 2019

Dubovik, Freitas, Gluza, Riemann, Usovitch

Phys. Lett. B 783 (2018) 86

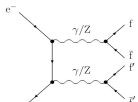
Measured **BHLUMI**

Phys. Rep. 427 (2006) 257

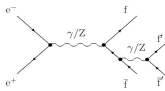
Light fermion pair production

Four-fermion final state (with at least one e^+e^- pair) may pass the event selection

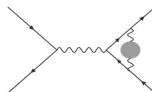
Conversion



Annihilation



Virtual correction at the same order
(interference with tree-level graph)



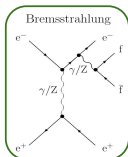
J. Gluza et al, arXiv:0807.4691
Phys. Rev. D 78 (2008) 085019

Specific four fermion
MC generators

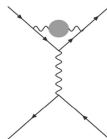
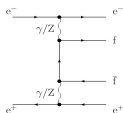
Positive correction

Negative correction

Dominant



Multiperipheral



Delicate cancelation w/o cuts, but **negative correction** when selection cuts are included
(smaller momentum and larger acoplanarity angle for the e^+e^- pair in the four-fermion final state)



Our result made it way to the Hall of Fame in 2023

Z HADRONIC POLE CROSS SECTION $41.4802 \pm 0.0325 \text{ nb}$

OUR EVALUATION is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 2006). Corrections as discussed in VOUTSINAS 2020 and JANOT 2020 are also included. This quantity is defined as

$$\sigma_h^0 = \frac{12\pi}{s} \frac{\Gamma(e^+e^-) \Gamma(\text{hadrons})}{\Gamma_Z^2}$$

It is one of the parameters used in the Z lineshape fit.

VALUE [nb]	EVTS	DOCUMENT ID	TECN
41.4802 ± 0.0325		OUR EVALUATION	
41.4802 ± 0.0325		¹ JANOT 2020	
		² VOUTSINAS 2020	
		LEP-SLC 2006	
41.500 ± 0.037			
41.541 ± 0.037		³ ABBENDI 2001A	OPAL
41.501 ± 0.055	4.10M	ABREU 2000F	DUPH
41.578 ± 0.069	3.70M	ACCARIERI 2000C	L3
41.535 ± 0.055	3.54M	BARATE 2000C	ALEP
41.559 ± 0.058	4.07M	ABRAMS 1989B	MRK2
42 ± 4	450		

¹ JANOT 2020 applies a correction to LEP-SLC 2006 using an updated Bhabha cross section calculation account for correlated luminosity bias as presented in VOUTSINAS 2020.

² VOUTSINAS 2020 applies a correction to LEP-SLC 2006 to account for correlated luminosity bias.

See <https://pdglive.lbl.gov/>

- [Number of neutrino types](#)
- [Z properties](#) (under Gauge and Higgs bosons)

Number of Light ν Types 2.996 ± 0.007
(already in the PDG in 2020)

VALUE	DOCUMENT ID	TECN
2.9963 ± 0.0074	¹ JANOT 2020	
	² VOUTSINAS 2020	
	³ LEP-SLC 2006	RVUE
2.9918 ± 0.0081		
2.9840 ± 0.0082		

• • We do not use the following data for averages, fits, limits, etc. • •

$E_{\text{CM}}^{\text{Z}} = 88 - 94 \text{ GeV}$

Z WIDTH $2.4955 \pm 0.0023 \text{ GeV}$

OUR EVALUATION is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 2006). Corrections as discussed in VOUTSINAS 2020 and JANOT 2020 are also included.

VALUE [GeV]	EVTS	DOCUMENT ID	TECN	COMMENT
2.4955 ± 0.0023		OUR EVALUATION		
2.4955 ± 0.0023		¹ JANOT 2020		
		² VOUTSINAS 2020		
		LEP-SLC 2006		
2.4955 ± 0.0023				
2.4952 ± 0.0023				

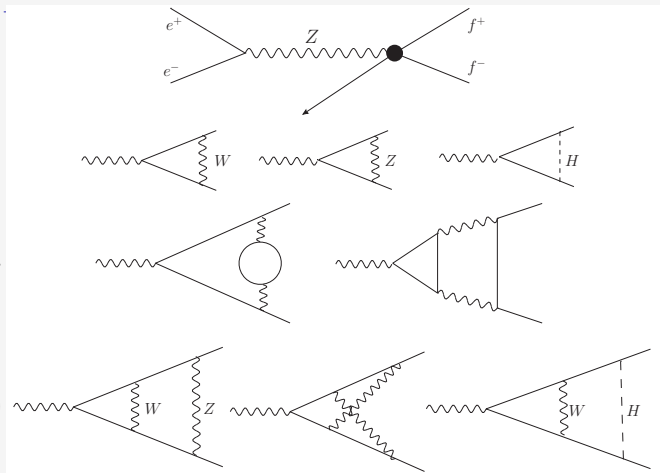
• • We do not use the following data for averages, fits, limits, etc. • •

$E_{\text{CM}}^{\text{Z}} = 88 - 94 \text{ GeV}$

$N_\nu = 2.9963 \pm 0.0074$ - "wirtualny" ślad "nowej fizyki"!
Świetny przykład połączenia teorii (nasze obliczenia NNLO plus BHLUMI) i eksperymentalna analiza danych.

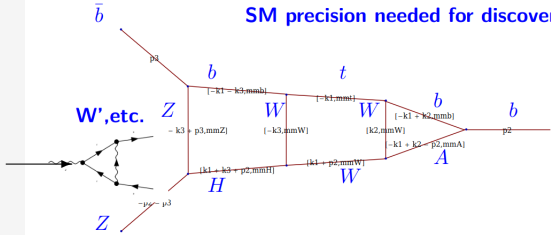
Flagowy przykład precyzyjnej fizyki w fizyce wysokich energii e^+e^- :

wierzchołek $Zf\bar{f}$ i poprawki elektrosłabe



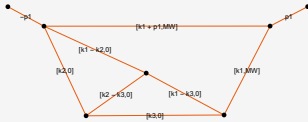
I. Dubovyk, A. Freitas, JG, T. Riemann, J. Usovitsch,
<https://doi.org/10.1016/j.physletb.2016.09.012>
<https://doi.org/10.1016/j.physletb.2018.06.037>
[https://doi.org/10.1007/JHEP08\(2019\)113](https://doi.org/10.1007/JHEP08(2019)113)

SM precision needed for discovery studies (indirect effects)

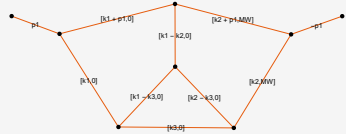


Four scales :

$$\left\{ \frac{M_H^2}{M_Z^2}, \frac{M_W^2}{M_Z^2}, \frac{m_t^2}{M_Z^2}, \frac{s+i\epsilon}{M_Z^2} \right\}$$



1-dim



4-dim

$$-18.779406962 - 6.390785027i \quad -22.5213 + 4.74442i \pm (0.001 + 0.001i)$$

$$I = -\frac{1}{(-s)^{1+3\epsilon}} \int_{-i\infty}^{+i\infty} \prod_{i=1}^4 dz_i \left(-\frac{M_W^2}{s} \right)^{z_3} \frac{\Gamma(-\epsilon - z_1) \Gamma(-z_1) \Gamma(1 + 2\epsilon + z_1)}{\Gamma(1 - 2\epsilon) \Gamma(1 - 3\epsilon - z_1)} \\ \times \frac{\Gamma(-2\epsilon - z_{12}) \Gamma(1 - \epsilon + z_2) \Gamma(1 + z_{12}) \Gamma(1 + \epsilon + z_{12}) \Gamma(1 + 3\epsilon + z_3) \Gamma(1 - \epsilon - z_4)}{\Gamma(1 - z_2) \Gamma(2 + \epsilon + z_{12})} \dots$$

Reprezentacje Mellin-Barnesa w HEP - metoda

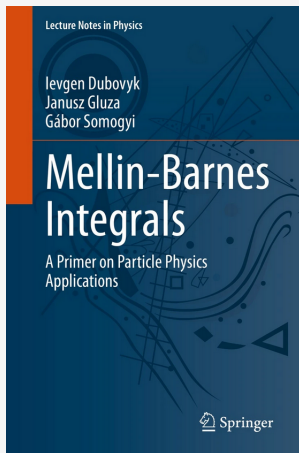
- "Om definita integraler", R. H. Mellin, Acta Soc. Sci. Fenn. 20(7), 1 (1895),
"The theory of the gamma function", E. W. Barnes Messenger Math. 29(2), 64 (1900).

$$\begin{aligned} \text{matematyka} &\longrightarrow \frac{1}{(A+B)^\lambda} = \frac{1}{\Gamma(\lambda)} \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \Gamma(\lambda+z) \Gamma(-z) \frac{B^z}{A^{\lambda+z}} \\ \text{fizyka} &\longrightarrow \frac{1}{(p^2 - m^2)^a} = \frac{1}{\Gamma(a)} \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \Gamma(a+z) \Gamma(-z) \frac{(m^2)^z}{(p^2)^{a+z}} \end{aligned}$$

Rekurencja $\implies$ wielowymiarowe zespolone całki.

$$\begin{aligned} \frac{1}{(A_1 + \dots + A_n)^\lambda} &= \frac{1}{\Gamma(\lambda)} \frac{1}{(2\pi i)^{n-1}} \int_{-i\infty}^{i\infty} dz_1 \dots dz_{n-1} \\ &\times \prod_{i=1}^{n-1} A_i^{z_i} A_n^{-\lambda - z_1 - \dots - z_{n-1}} \prod_{i=1}^{n-1} \Gamma(-z_i) \Gamma(\lambda + z_1 + \dots + z_{n-1}) \end{aligned}$$

Analityczne i numeryczne obliczanie całek Mellin-Barnesa



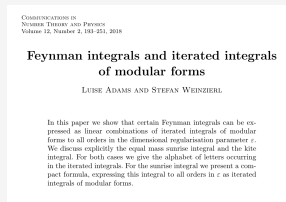
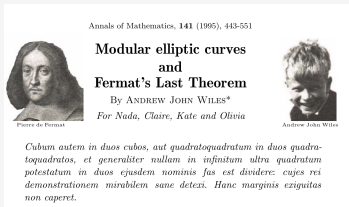
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<https://arxiv.org/abs/2211.13733>

<https://doi.org/10.1007/978-3-031-14272-7>

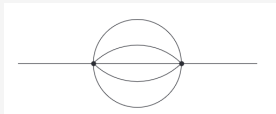
$$T = 2\pi\sqrt{\frac{l}{g}} \times {}_2F_1\left[\frac{1}{2}, \frac{1}{2}; 1; \sin^2 \theta\right]$$

Rozwiązania analityczne dla masywnych całek wielopetlowych, opisujących procesy/rozpady rozpraszania, wykraczają poza funkcje eliptyczne – jak daleko?

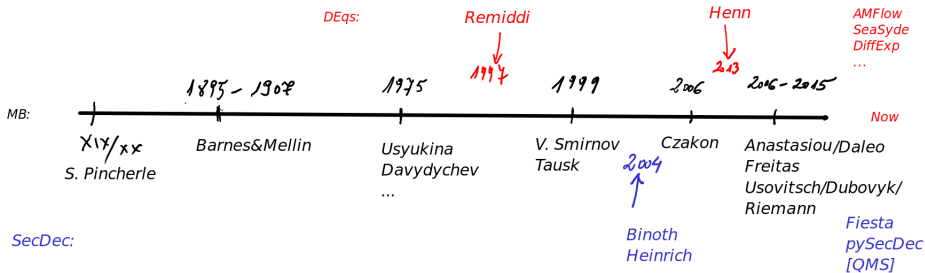


"Epsilon-factorized form" of banana 4-loop integrals → fast evaluations to nearly arbitrary precision.

'Taming Calabi-Yau Feynman Integrals: The Four-Loop Equal-Mass Banana Integral', S. Pögel, X. Wang, S. Weinzierl,

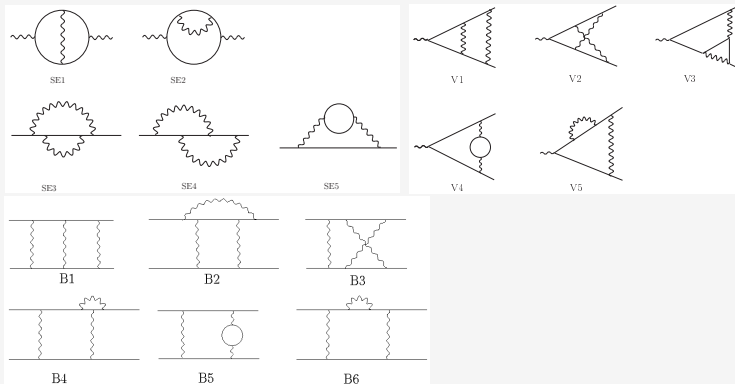


Trzy wiodące metody obliczeń (analityczne oraz numeryczne)



Analityczne rozwiązania, przykład rozpraszanie Bhabhy $e^+e^- \rightarrow e^+e^-$

$$G(X) = \frac{1}{(i\pi^{d/2})^L} \int \frac{d^d k_1 \dots d^d k_L}{(q_1^2 - m_1^2)^{\nu_1} \dots (q_j^2 - m_j^2)^{\nu_j} (q_N^2 - m_N^2)^{\nu_N}} X$$



Rozpraszanie Bhabhy $e^+e^- \rightarrow e^+e^-$

PRL **100**, 131602 (2008)

PHYSICAL REVIEW LETTERS

week ending
4 APRIL 2008

energy kernel is trivial, $K_{SE}(z) = 1/(q^2 - z)$. Whereas the two-loop vertex kernel K_v can be found in Eq. (5) of [15], the cross-section corrections due to the double boxes of Fig. 1(c) depend on three such kernels $K_a(z)$, $a = A, B, C$ [12,16]. Notice that, unlike the vertex kernel, the box kernels are infrared divergent, but, analogously to the one-loop box, they have no singularity in the electron mass. The net cross-section contribution from the eight double box diagrams is still infrared divergent. These boxes become, as well as the reducible diagrams with

one-loop vertices and boxes, infrared finite after adding real soft photon emission. The anatomy of that is nicely detailed in Section 2.2 of [7]. In order to construct an infrared-finite quantity, we combine: (i) Born diagrams interfering with the two-loop box diagrams [Fig. 1(c)] and with reducible vertices [Fig. 1(b)], (ii) Born diagrams with a one-loop vacuum polarization function interfering with single one-loop boxes and vertices, and finally (iii) real single-photon emission with a one-loop vacuum polarization [18]. The resulting cross section becomes

$$\frac{d\hat{\sigma}}{d\Omega} = c \int_{4M_e^2}^{\infty} dz \frac{R_{\text{had}}(z)}{z} \frac{1}{t-z} F_1(z) + c \int_{4M_e^2}^{\infty} \frac{dz}{z(s-z)} \left\{ R_{\text{had}}(z) \left[F_2(z) + F_3(z) \ln \left| 1 - \frac{z}{s} \right| \right] - R_h(s) \left[F_2(s) + F_3(s) \ln \left| 1 - \frac{z}{s} \right| \right] \right\} \\ + c \frac{R_h(s)}{s} \left\{ F_2(s) \ln \left(\frac{s}{4M_e^2} - 1 \right) - 6\zeta_2 F_a(s) + F_3(s) \left[2\zeta_2 + \frac{1}{2} \ln^2 \left(\frac{s}{4M_e^2} - 1 \right) + \text{Li}_2 \left(1 - \frac{s}{4M_e^2} \right) \right] \right\}, \quad (5)$$

with $c = \alpha^4/(\pi^2 s)$ and $R_h(s) = \theta(s - 4M_e^2) R_{\text{had}}(s)$. Further,

$$F_1(z) = \frac{1}{3} \left\{ 9\bar{c}(s, t) \ln \left(\frac{s}{m_e^2} \right) + \left[-z^2 \left(\frac{1}{s} + \frac{2}{t} + 2\frac{s}{t^2} \right) + z \left(4 + 4\frac{s}{t} + 2\frac{t}{s} \right) + \frac{1}{2} \frac{t^2}{s} + 6\frac{s^2}{t} + 5s + 4t \right] \ln \left(-\frac{t}{s} \right) + s \left(-\frac{z}{t} + \frac{3}{2} \right) \right. \\ \times \ln \left(1 + \frac{t}{s} \right) + \left[\frac{1}{2} \frac{z^2}{s} + 2z \left(1 + \frac{s}{t} \right) - \frac{11}{4} s - 2t \right] \ln^2 \left(-\frac{t}{s} \right) - \left[\frac{1}{2} \frac{z^2}{t} - z \left(1 + \frac{s}{t} \right) + \frac{t^2}{s} + 2\frac{s^2}{t} + \frac{9}{2} s + \frac{15}{4} t \right] \\ \times \ln^2 \left(1 + \frac{t}{s} \right) + \left[\frac{z^2}{t} - 2z \left(1 + \frac{s}{t} \right) + 2\frac{s^2}{t} + 5s + \frac{5}{2} t \right] \ln \left(-\frac{t}{s} \right) \ln \left(1 + \frac{t}{s} \right) - 4 \left[\frac{t^2}{s} + 2\frac{s^2}{t} + 3(s+t) \right] \\ \times \left[1 + \text{Li}_2 \left(-\frac{t}{s} \right) \right] - \left[\frac{t^2}{s} + 2\frac{s^2}{t} + 3(s+t) \right] \ln \left(\frac{z}{s} \right) \ln \left(1 + \frac{t}{s} \right) - \left[2\frac{z^2}{t} - 4z \left(1 + \frac{s}{t} \right) - 4\frac{t^2}{s} - 2\frac{s^2}{t} + s - \frac{11}{2} t \right] \zeta_2 \\ + \left[z^2 \left(\frac{1}{s} + 2\frac{s}{t^2} + \frac{2}{t} \right) - z \left(\frac{t}{s} + 2\frac{s}{t} + 2 \right) \right] \ln \left(\frac{z}{s} \right) - \left[z^2 \left(\frac{1}{s} + \frac{1}{t} \right) + 2z \left(1 + \frac{s}{t} \right) + s + 2\frac{s^2}{t} \right] \ln \left(\frac{z}{s} \right) \ln \left(1 + \frac{z}{s} \right) \\ + \left[\frac{z^2}{s} + 4z \left(1 + \frac{s}{t} \right) - \frac{t^2}{s} - 4(s+t) \right] \ln \left(\frac{z}{s} \right) \ln \left(1 - \frac{z}{t} \right) - \left[z^2 \left(\frac{1}{s} + 2\frac{s}{t^2} + \frac{2}{t} \right) - 2z \left(\frac{t}{s} + \frac{s}{t} + 2 \right) + \frac{t^2}{s} + 2(s+t) \right] \\ \times \ln \left(1 - \frac{z}{t} \right) + \left[\frac{z^2}{t} - 2z \left(1 + \frac{s}{t} \right) + 2\frac{t^2}{s} + 8s + 4\frac{s^2}{t} + 7t \right] \ln \left(1 - \frac{z}{t} \right) \ln \left(1 + \frac{t}{s} \right) + \left[\frac{z^2}{s} + 4z \left(1 + \frac{s}{t} \right) - \frac{t^2}{s} - 4(s+t) \right] \\ \times \text{Li}_2 \left(\frac{z}{t} \right) - \left[z^2 \left(\frac{1}{s} + \frac{1}{t} \right) + 2z \left(1 + \frac{s}{t} \right) + s + 2\frac{s^2}{t} \right] \text{Li}_2 \left(-\frac{z}{s} \right) - \left[\frac{z^2}{t} - 2z \left(1 + \frac{s}{t} \right) + \frac{t^2}{s} + 5s + 2\frac{s^2}{t} + 4t \right] \\ \times \text{Li}_2 \left(1 + \frac{z}{u} \right) \left. \right\} + 4\bar{c}(s, t) \ln \left(\frac{2\omega}{\sqrt{s}} \right) \left[\ln \left(\frac{s}{m_e^2} \right) + \ln \left(-\frac{t}{s} \right) - \ln \left(1 + \frac{t}{s} \right) - 1 \right]. \quad (6)$$

Harmoniczne polilogarytmy

$$H(0; x) = \ln x ,$$

$$H(1; x) = \int_0^x \frac{dx'}{1-x'} = -\ln(1-x) ,$$

$$H(-1; x) = \int_0^x \frac{dx'}{1+x'} = \ln(1+x)$$

$$H(\vec{0}_w; x) = \frac{1}{w!} \ln^w x$$

$$H(\vec{m}_w; x) = \int_0^x dx' f(a; x') (\vec{m}_{w-1}; x')$$

Proces Bhabhy, prosty alfabet i litery (tylko 3):

$$f(0, x) = \frac{1}{x}, \quad f(1, x) = \frac{1}{1-x}, \quad f(-1, x) = \frac{1}{1+x}$$

$$H(0; x) = \ln x, H(0, 0; x) = \int \frac{H[0, x]}{x} dx = \frac{1}{2!} \ln^2 x, \dots$$

$$H(1; x) = \int_0^x \frac{dx'}{1-x'} = -\ln(1-x),$$

$$H(-1; x) = \int_0^x \frac{dx'}{1+x'} = \ln(1+x)$$

$$H(0, 1; x) = \int_0^x \frac{dx'}{x'} H(1; x') = - \int_0^x \frac{dx'}{x'} \ln(1-x'),$$

$$H(0, 1; x) = {}_2(x),$$

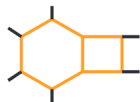
$$H(1, 0; x) = \int_0^x \frac{dx'}{1-x'} H(0; x') = \int_0^x \frac{dx'}{1-x'} \ln x'$$

$$\begin{aligned}
 H(1, 0; x) &= \int_0^x \frac{dx'}{1-x'} H(0; x') = \int_0^x \frac{dx'}{1-x'} \ln x' , \\
 H(1, 0; x) &= -\ln x \ln(1-x) + {}_2(x) \\
 H(1, 0; x) &= H(0, x)H(1, x) - H(0, 1; x)
 \end{aligned}$$

$$\begin{aligned}
 H(\vec{a}, x)H(\vec{b}, x) &= \sum_{\vec{c}=\vec{a}+\vec{b}} H(\vec{c}, x) \\
 H(m_1, \dots, m_q, x) &= H(m_1, x)H(m_2, \dots, m_q, x) \\
 &\quad - H(m_1, m_2, x)H(m_3, \dots, m_q, x) \\
 &\quad + \dots + (-1)^{q+1} H(m_q, \dots, m_1, x)
 \end{aligned}$$

Two-Loop Integrals – Hexabox Topology

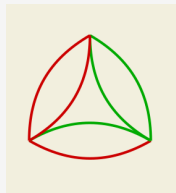
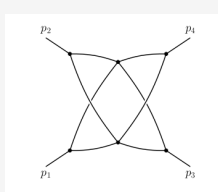
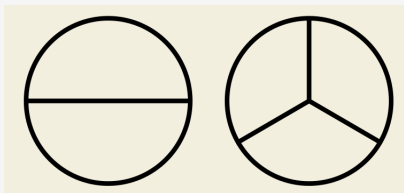
$$dI = \epsilon \mathbf{M} I, \quad \mathbf{M} = M_\alpha d \log W_\alpha$$



- **202 master integrals**, in 126 different sectors
 - 185 are five-point one-mass topologies. All this for 17 new integrals!
- **128 letters**, only 11 new ones (associated with six-point topologies)
 - 9 appear in the differential equation on the maximal cut, 2 appear in the off-shell penta-triangle
 - 2 'off-shell letters' are odd under $\sqrt{\Delta_5}$, 9 'on-shell' letters are even
- Closure of the alphabet under dihedral symmetry: **245 letters**
- From here we know how to proceed to get basis of functions, evaluate integrals, ...

Rodzaje diagramów

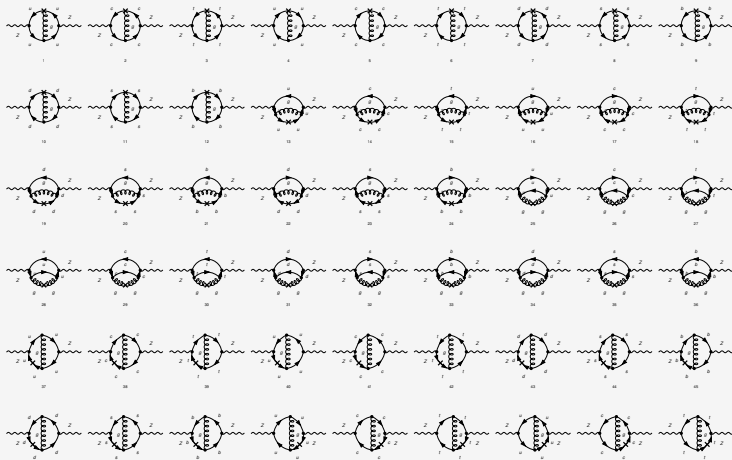
Mamy więc ciekawe topologicznie struktury diagramów Feynmana i odpowiadające im całki, np. sunrise (sunset), mercedes, crown, neckles, ladder, ...,



W przypadku analitycznych rozwiązań mamy wielo-literowe alfabety do znajdowania matematycznych funkcji im odpowiadających.

I numerycznie: Wielopętlowe obliczenia, rozbieżności, renormalizacja ...

$$\left(\text{YInt}(\text{dd}(k_1, 0), \text{dd}(k_2, 0), \text{dd}(k_1 + k_2, 0), \text{dd}(k_2 - p_1, 0), \text{dd}(k_1 + p_1, 0), \text{dn}_2(k_1, k_2), \text{dn}_2(k_1, k_2)) \rightarrow (-13.62737119439388771410934383 + 20.1668958438921923453391 \right. \\ (11.49351870466377347494786009 - 21.85421682411580147240037042 i) \text{eps}^2 + (1.010805534687752347048684410 + 8.050331174823845173560523670 i) \text{eps} + \\ \left. (1.281250000000000000000000000000 + 1.178097245096172464423491269 i) + \frac{0.187500000000000000000000000000 + 0. \times 10^{-29} i}{\text{eps}} \right),$$



Zamiast podsumowania

*"in this field, almost everything is
already discovered, and all that remains
is to fill a few unimportant holes"*



Philipp von Jolly
(1809-1884)

advice to the young Max Planck
not to go into physics, Munich 1878

Albert Michelson (1894):

"It seems probable that most of the grand underlying principles have been firmly established (...) the future truths of physical science are to be looked for in the sixth place of decimals"

Q: Dear Albert: What about special and general relativity, and quantum mechanics, particle physics, ...?

Dziękuję za uwagę.

Dodatkowe slajdy

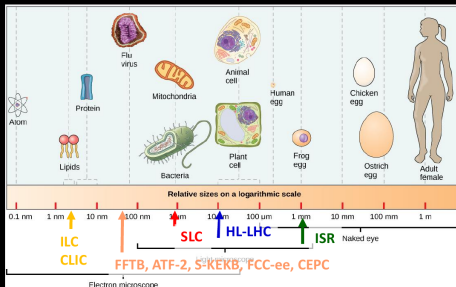
$$F = k \frac{qQ}{r^2} \equiv \frac{1}{4\pi\epsilon_0} \frac{qQ}{r^2} \rightarrow \alpha \equiv \frac{\phi e}{mc^2} = \frac{\frac{e^2}{4\pi\epsilon_0 L}}{mc^2} = \frac{e^2}{4\pi\epsilon_0 \hbar c}, \quad L = \frac{\hbar}{mc}$$

Stała struktury subtelnej α wynosi liczbowo około 1/137 [137.035999206(11)]

Czy 1/136 lub 1/138 robi różnicę?

Procentowe zmiany w wartości α implikują czerwone lub niebieskie gwiazdy
("Gravitation", Misner, Thorne, Wheeler)

vertical spot size challenge



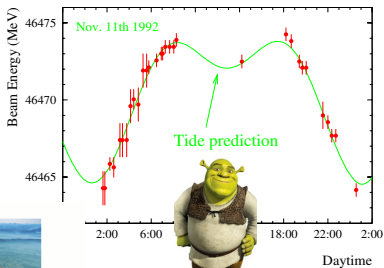
FCC-ee in the regime of FFTB, ATF-2, and especially SuperKEKB



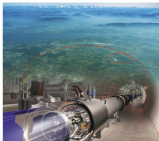
Moonrise over LEP



Fall of 1992 : The historic tide experiment !

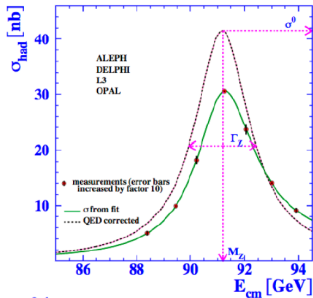


total strain is 4×10^{-8} ($\Delta C = 1$ mm)



Aby dojść do eksperymentalnej precyzji, musimy poprawić obliczenia!

□ Cross section : Z mass and width



♦ ~30% QED corrections (ISR)

CERN Yellow Reports:
Monographs

CERN-2019-003

Standard Model Theory for
the FCC-ee Tera-Z stage

A. Blondel
J. Gluza
S. Jadach
P. Janot
T. Riemann



$$\frac{\delta\Gamma_Z}{\Gamma_Z} = \frac{.025}{2495} \simeq 10^{-5}!$$

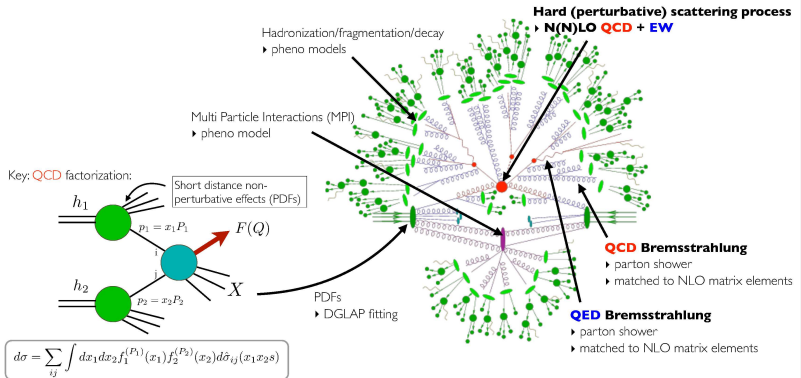
Observable	present value $\pm$ error	FCC-ee Stat.	FCC-ee Syst.	Comment and leading exp. error
m_Z (keV)	91186700 ± 2200	4	100	From Z line shape scan Beam energy calibration
Γ_Z (keV)	2495200 ± 2300	4	25	From Z line shape scan Beam energy calibration
$\sin^2 \theta_W^{\text{eff}} (\times 10^6)$	231480 ± 160	2	2.4	from $A_{\text{FB}}^{\mu\mu}$ at Z peak Beam energy calibration
$1/\alpha_{\text{QED}}(m_Z^2) (\times 10^3)$	128.2 ± 14	3	See II	from $A_{\text{FB}}^{\mu\mu}$ off peak QED&EW errors dominate
$R_\ell^Z (\times 10^3)$	20767 ± 25	0.1	0.2-1	ratio of hadrons to leptons acceptance for leptons
$\alpha_s(m_Z^2) (\times 10^4)$	1196 ± 30	0.1	0.4-1.6	from R_ℓ^Z above
$\sigma_{\text{had}}^0 (\times 10^3 \text{ nb})$	41541 ± 37	0.1	4	peak hadronic cross section luminosity measurement
$\sigma_Z (\times 10^3)$	2996 ± 7	0.005	1	Z peak cross sections Luminosity measurement
$R_b (\times 10^6)$	216290 ± 660	0.3	< 60	ratio of $b\bar{b}$ to hadrons stat. extrapol. from SLD
$A_{\text{FB},0}^b (\times 10^{-1})$	992 ± 16	0.02	1-3	b-quark asymmetry at Z pole from jet charge
$A_{\text{FB}}^{\tau,\tau} (\times 10^4)$	1498 ± 49	0.15	< 2	τ polarization asymmetry τ decay physics
τ lifetime (fs)	290.3 ± 0.5	0.001	0.04	radial alignment
τ mass (MeV)	1776.86 ± 0.12	0.004	0.04	momentum scale
τ leptonic ($\mu\nu_\mu\nu_\tau$) B.R. (%)	17.38 ± 0.04	0.0001	0.003	e/μ /hadron separation
m_W (MeV)	80350 ± 15	0.25	0.3	From WW threshold scan Beam energy calibration

PRECISION

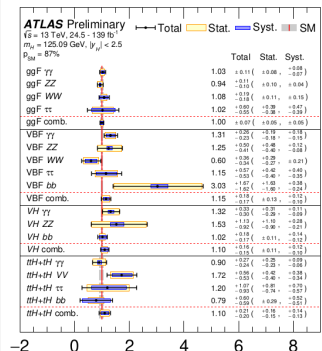
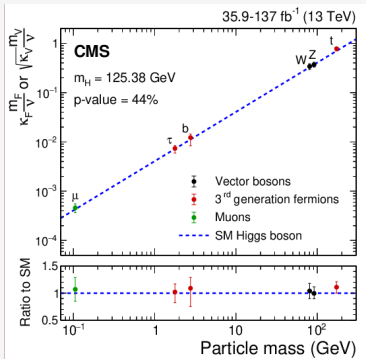
LHC and HL-LHC

HL-LHC aims at percent precision

Theoretical Predictions for the LHC



Od 2012 roku: stały wzrost precyzji, szeroka gama procesów

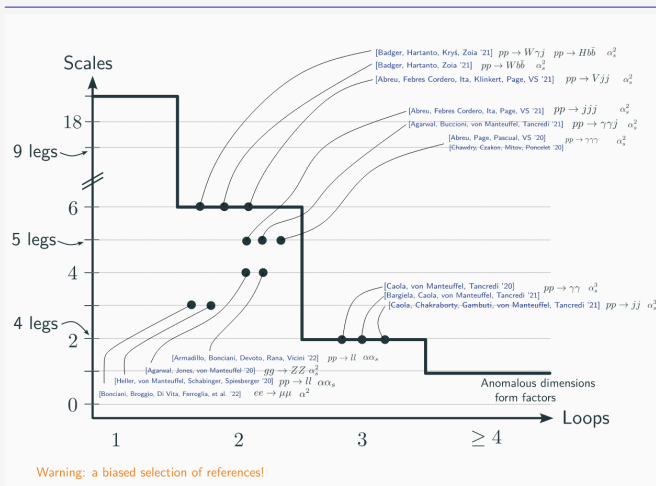


ATLAS 2020: <https://cds.cern.ch/record/2725733>

CMS 2020: <http://cds.cern.ch/record/2730058>

Workshop: Precision calculations for future e^+e^- colliders: targets and tools,

CERN 2022, talk by V. Sotnikov



Still not automated, systematic, even at 2nd order, though the progress in methods and tools impressive

just one example

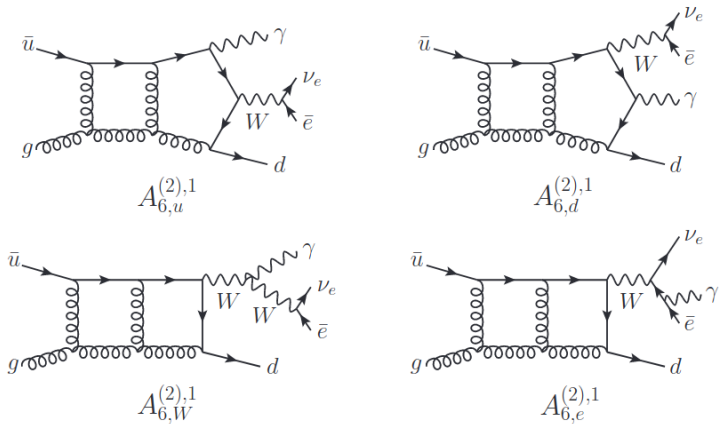
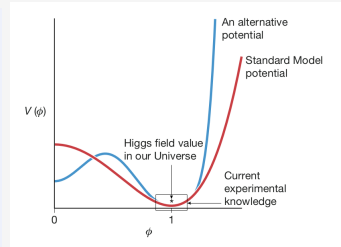
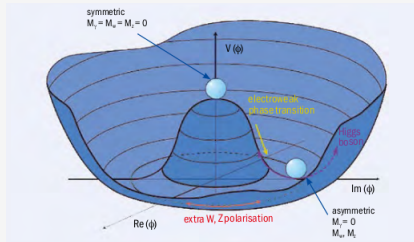


Figure 1. Sample two-loop Feynman diagrams for $W^+\gamma jj$ production.

Łamanie symetrii i kształt potencjału

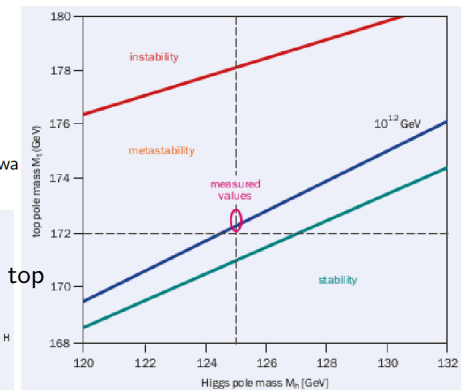
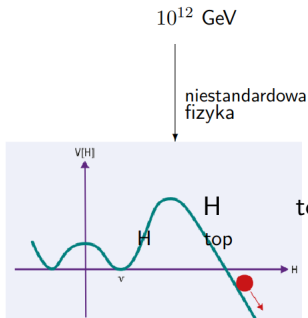


$$\Phi \equiv \Phi_{SM} = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$$

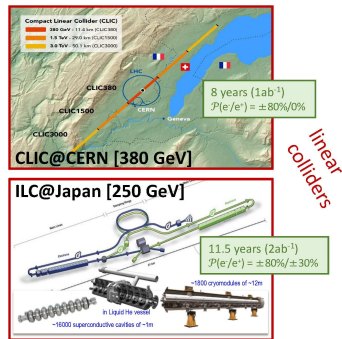
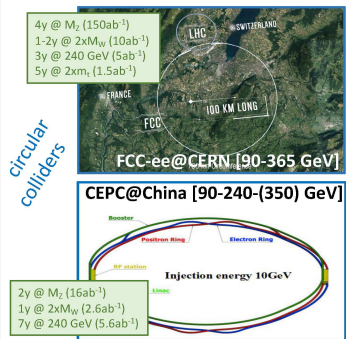
$$V = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 \Leftrightarrow y = ax + bx^2, y \equiv V, x \equiv \Phi^\dagger \Phi$$

$$V_{min} = v/\sqrt{2}, v = \sqrt{\mu^2/\lambda} \simeq 250 \text{ GeV}$$

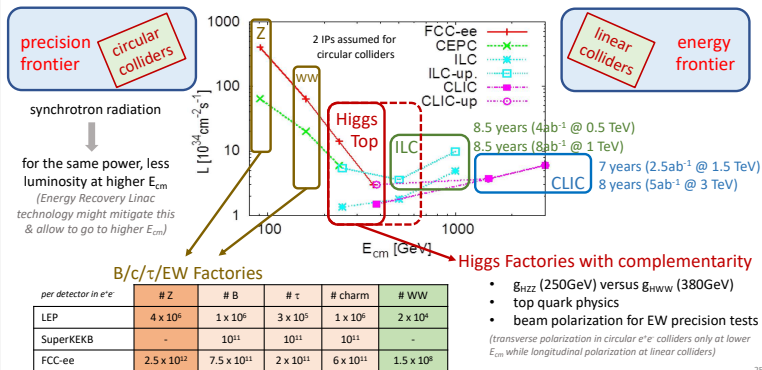
Łamanie symetrii i kształt potencjału, konsekwencje



e^+e^- Higgs Factories



e^+e^- Higgs Factories (incl. B/c/ τ /EW/top factories)



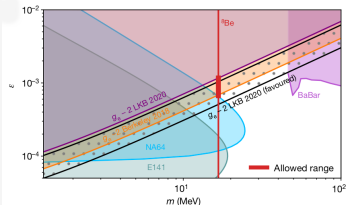
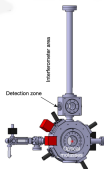
Czy pomiar α ma wpływ na niestandardowe modele (BSM)?

Table 1 | Error budget on α

Source	Correction ($\times 10^{-7}$)	Relative uncertainty ($\times 10^{-7}$)
Gravitational gradient	-0.6	0.1
Alignment of the beams	0.5	0.5
Combinational acceleration	1.2	1.2
Frequency of the laser	0.3	0.3
Wave-front curvature	0.8	0.3
Wave-front distortion	3.9	1.9
Gravitational phase	108.2	5.4
Residual Raman light shift	2.3	2.3
Index of refraction	0	<0.1
Internal interaction	0	<0.1
Light shift (two-photon interaction)	-15.0	2.3
Second-order Zeeman effect	0.1	0.1
Phase shift in Raman phase lock loop	-26.8	0.6
Global systematic effects	64.2	6.9
Statistical uncertainty	3.4	3.4
Relative mass of ^{87}Rb ($86.0001369126(8)$)	3.5	3.5
Relative mass of the electron ($5.48579909064(28) \times 10^{-4}$)	15	15
Rydberg constant ($10973731.568160(21) \text{ m}^{-1}$)	0.1	0.1
Total $\alpha = 137.035999084(21)$	6.9	6.9

For each systematic effect, more discussion can be found in Methods. *From ref. 1. †From <https://pdg.lbl.gov/2018/Constants/>.

Morel, L., Yao, Z., Cladé, P. et al. Determination of the fine-structure constant with an accuracy of 81 parts per trillion. *Nature* 586, 61–65 (2020).



Podstruktura: m^* , rozmiar $L = \hbar/(m^*c) \rightarrow$ dodatkowa modyfikacja rzędu $\delta a_e \simeq m_e/m^*$
stąd wykluczone:

$$m^* < 520 \text{ GeV} \equiv L > 2 \times 10^{-18} / 4 \times 10^{-19} \text{ m}$$

Eksperyment ma polepszyć dokładność δa_e o rząd w następnych latach, będzie na poziomie czułości $(g - 2)_\mu$.

Ograniczenia

U wielomian dla nieplanarnego 3-pętlowego pudełka (64 człony)

$$\begin{aligned} & x[1] x[2] x[4] + x[1] x[3] x[4] + x[2] x[3] x[4] + x[1] x[2] x[5] + \\ & x[1] x[3] x[5] + x[2] x[3] x[5] + x[1] x[4] x[5] + x[2] x[4] x[5] + \\ & x[2] x[4] x[6] + x[3] x[4] x[6] + x[2] x[5] x[6] + x[3] x[5] x[6] + \\ & x[4] x[5] x[6] + x[2] x[4] x[7] + x[3] x[4] x[7] + x[2] x[5] x[7] + \\ & x[3] x[5] x[7] + x[4] x[5] x[7] + x[1] x[2] x[8] + x[1] x[3] x[8] + \\ & x[2] x[3] x[8] + x[1] x[4] x[8] + x[2] x[4] x[8] + x[2] x[6] x[8] + \\ & x[3] x[6] x[8] + x[4] x[6] x[8] + x[2] x[7] x[8] + x[3] x[7] x[8] + \\ & x[4] x[7] x[8] + x[1] x[2] x[9] + x[1] x[3] x[9] + x[2] x[3] x[9] + \\ & x[2] x[4] x[9] + x[3] x[4] x[9] + x[1] x[5] x[9] + x[3] x[5] x[9] + \\ & x[4] x[5] x[9] + x[2] x[6] x[9] + x[3] x[6] x[9] + x[5] x[6] x[9] + \\ & x[2] x[7] x[9] + x[3] x[7] x[9] + x[5] x[7] x[9] + x[1] x[8] x[9] + \\ & x[3] x[8] x[9] + x[4] x[8] x[9] + x[6] x[8] x[9] + x[7] x[8] x[9] + \\ & x[1] x[2] x[10] + x[1] x[3] x[10] + x[2] x[3] x[10] + \\ & x[1] x[4] x[10] + x[2] x[4] x[10] + x[2] x[6] x[10] + \\ & x[3] x[6] x[10] + x[4] x[6] x[10] + x[2] x[7] x[10] + \\ & x[3] x[7] x[10] + x[4] x[7] x[10] + x[1] x[9] x[10] + \\ & x[3] x[9] x[10] + x[4] x[9] x[10] + x[6] x[9] x[10] + x[7] x[9] x[10] \end{aligned}$$

F-wielomian Symanzika: problem $\Gamma[0]$, twierdzenie Cheng-Wu

- ▶ 2-loop: $\delta(1 - v_1 - v_2)$, $U(\vec{v}) = v_3 + v_1 v_2$
brak dodatkowych całek MB from U
- ▶ 3-loop: $\delta(1 - v_1 - v_2 - v_3)$
 - ▶ "ladder" - 2 dodatkowe całki MB 64-dim $\rightarrow$ 2-dim (!)
 - ▶ "mercedes" - 4 dodatkowe całki MB

Aby otrzymać minimalny wymiar:

- ▶ 1-loop: $U(\vec{x}) \equiv 1$ gdziekolwiek możliwe
- ▶ 2- i 3-pętle: wyrażenie dla F nie rozwijane

$$F = F_0 + U \sum_{i=1}^N x_i m_i^2$$

Drugi człon można rozszerzyć (generując progi, bez potrzeby deformacji konturu) lub nie (mniejsza wymiarowość kosztem gorszej zintegrowanej zbieżności)

Metoda AMFlow, $\eta = \infty \longrightarrow \eta = 0^+$

analityczna kontynuacja (pływ dodatkowej masy)

2. Prace Jan 27 2022 papers by Zhi-Feng Liu, Yan-Qin Ma and Xiao Liu:

<https://inspirehep.net/literature/2020677>, <https://inspirehep.net/literature/2020676>,
<https://inspirehep.net/literature/2020880> and 1711.09572
<https://inspirehep.net/literature/1639025>.

$$\tilde{I}_{\vec{\nu}}(\eta) = \int \left(\prod_{i=1}^L \frac{d^D \ell_i}{i\pi^{D/2}} \right) \frac{\tilde{\mathcal{D}}_{K+1}^{-\nu_{K+1}} \dots \tilde{\mathcal{D}}_N^{-\nu_N}}{\tilde{\mathcal{D}}_1^{\nu_1} \dots \tilde{\mathcal{D}}_K^{\nu_K}}.$$

$$\tilde{\mathcal{D}}_1 = \ell_1^2 - m^2 + i\eta$$

$$I_{\vec{\nu}} = \lim_{\eta \rightarrow 0^+} \tilde{I}_{\vec{\nu}}(\eta)$$

$$i \frac{\partial}{\partial \eta} \vec{J}(\eta) = A(\eta) \vec{J}(\eta)$$

Key point: warunki brzegowe przy $\eta \rightarrow \infty$ są całkami pęcherzykowymi (bubbles) pojedynczej skali masy, rozwiązywanymi iteracyjnie .

MI's dużej dokładności z AMFlow, wyniki

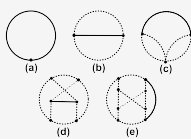


Diagram (a) is a circle with a single vertex at the bottom. Diagram (b) is a circle with a horizontal line across the middle. Diagram (c) is a circle with two vertices at the top and bottom, connected by a line. Diagram (d) is a circle with two vertices at the top and bottom, connected by two lines. Diagram (e) is a circle with two vertices at the top and bottom, connected by two lines, with an additional line connecting the two vertices.

$$I_{\vec{\nu}} = \int \left(\prod_{i=1}^L \frac{D\ell_i}{\pi^{D/2}} \right) \frac{\mathcal{D}_{K+1}^{-\nu_{K+1}} \dots \mathcal{D}_N^{-\nu_N}}{\mathcal{D}_1^{\nu_1} \dots \mathcal{D}_K^{\nu_K}}, \quad \mathcal{D}_1 = \ell_1^2 - m^2 + 0^+$$

$$\hat{I}_{\vec{\nu}'}(\ell_1^2) = \int \left(\prod_{i=2}^L \frac{D\ell_i}{\pi^{D/2}} \right) \frac{\mathcal{D}_{K+1}^{-\nu_{K+1}} \dots \mathcal{D}_N^{-\nu_N}}{\mathcal{D}_2^{\nu_2} \dots \mathcal{D}_K^{\nu_K}}, \quad I_{\vec{\nu}} = \{\Gamma[\dots]\} \hat{I}_{\vec{\nu}'}(-m^2)$$

$\uparrow$
L-loop

$\uparrow$
(L-1)-loop

$$\begin{aligned}
 I[(e)] = & -2.073855510286740\epsilon^{-2} - 7.812755312590133\epsilon^{-1} \\
 & - 17.25882864945875 + 717.6808845492140\epsilon \\
 & + 8190.876448160049\epsilon^2 + 78840.29598046500\epsilon^3 \\
 & + 566649.1116484678\epsilon^4 + 3901713.802716081\epsilon^5 \\
 & + 23702384.71086095\epsilon^6 + 14214293.68205112\epsilon^7,
 \end{aligned}$$

10 orders in ϵ , 16-digit precision.

Such an exact boundary point can be transported by DiffExp to any physical point.