

DHOST to unify Dark Matter and Dark Energy

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"Standard Model and Beyond" @Katowice, 21-23 Października 2022r.

Based upon

- Laudato E., Salzano V. & Umetsu K., “*Multicomponent DHOST analysis in galaxy clusters*”, MNRAS 511 (2022) 2, 1878-1892;
- Laudato E. & Salzano V., “*DHOST gravity in Ultra-diffuse galaxies - Part I: the case of NGC1052-DF2*”, EPJC 82 (2022) 7, 935;
- Laudato E. & Salzano V., “*DHOST gravity in Ultra-diffuse galaxies - Part II: NGC 1052-DF4 and Dragonfly 44*”, to be submitted.

Extended Theories of Gravity (ETGs)

- GR has theoretical and observational issues
 - Remind: “consensus” model (Λ CDM)
- GR tested directly only in the Solar System
 - Inference over 60 order of magnitudes in scale
- Dark Energy and Dark Matter
 - They are everywhere ($\sim 95\%$) in the Universe
- But... what are they?
 - DE: too many ideas \equiv no idea?
 - DM: no detection so far (sooner or later?)

Existence of DM - DE $\not\propto$ GR as THE gravity theory

ETGs: cons and pros

Cons (to make ETGs' opponents happy):

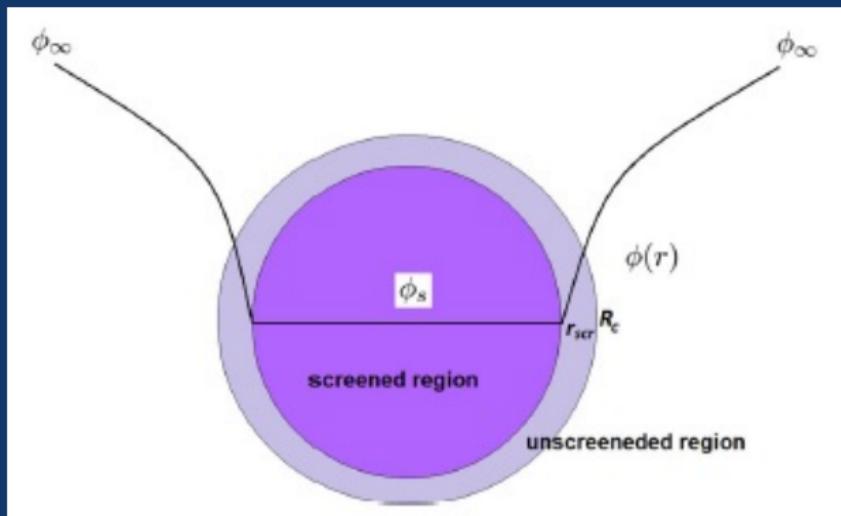
- too many possibilities!
- observations tune them almost undistinguishable from GR models
- how to discriminate among ETGs and GR? No smoking gun!
- statistically speaking, mostly disfavoured with respect to GR

Pros (to make me happy):

- the more ETGs are discarded, the more GR is confirmed (!)
- statistics is not everything
- solving “tensions” is more valuable (?)
- **DM and DE could be expression of totally NEW physics!**

ETGs: challenges

- ETGs are generally NOT proposed for DM, but mostly for DE
- once you modify GR on large scales, you use Screening Mechanisms to avoid influences also on smaller scales



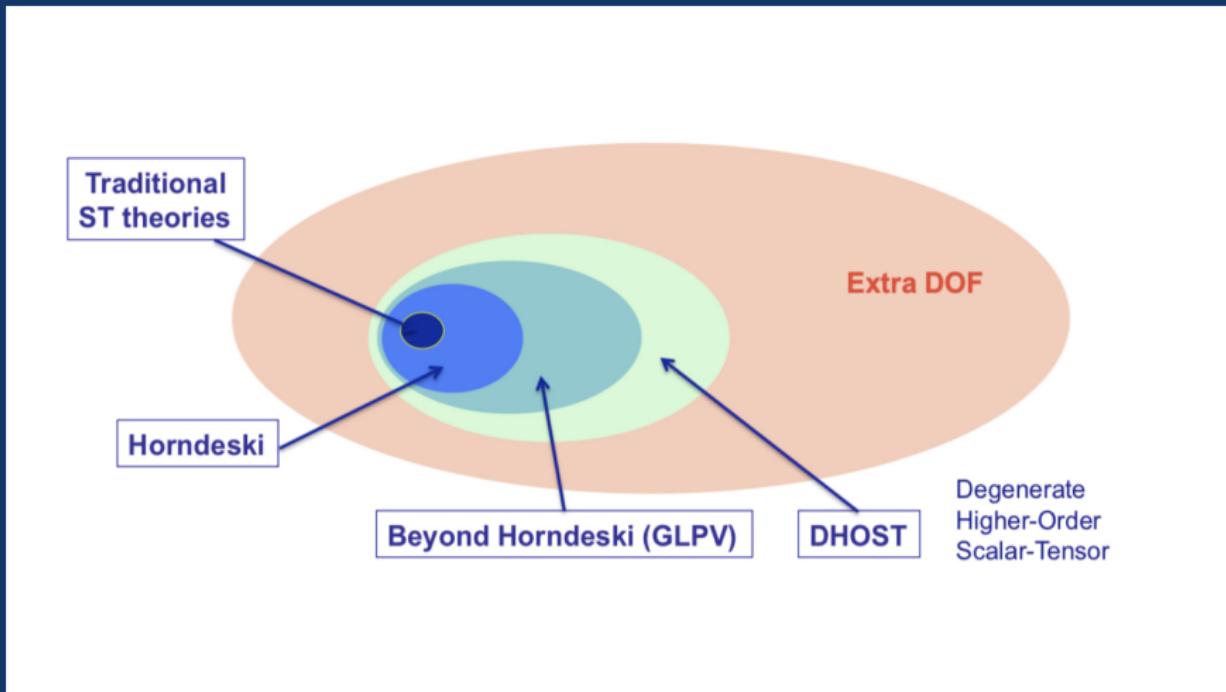
Key Idea: WHAT IF...

... Screening is partially broken, and ETGs effects might leak also into astrophysical scales?

- this is possible, but authors try to “heal” it
- when not, you might have a theory behaving as DE on cosmological scales and mimicking DM on smaller scales
- this would be a plus w.r.t. to GR, because you would unify the dark sector
- we will test this idea
 - Scenario 1: ETG as DE only, plus DM
 - Scenario 2: ETG as both DE and DM (no DM added!)

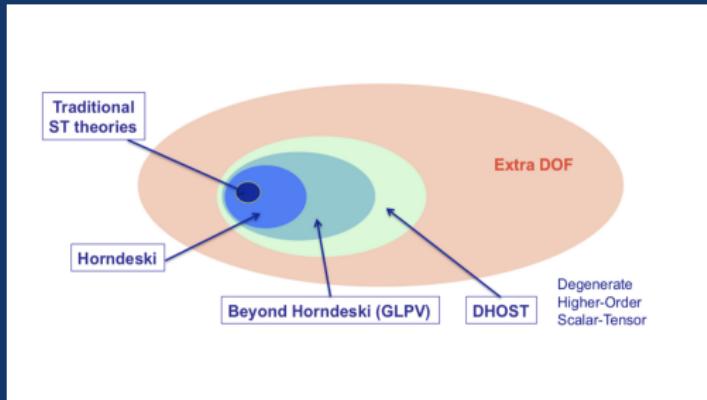
Our works: theoretical background

- DHOST: Degenerate Higher-Order Scalar-Tensor theories



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- The most general Lagrangian is (1711.06661):

$$\mathcal{L} = F_{(2)}(\phi, X)^{(4)} R + \sum_{i=1}^5 A_i(\phi, X) \mathcal{L}_{i,\phi}^{(2)}$$

Our works: theoretical background

- DHOST: Degenerate Higher-Order Scalar-Tensor theories
- partially broken (Vainshtein) screening (1712.04731)
- Gravitational and the metric potentials look like

$$\frac{d\Phi}{dr} = \frac{G_N M(r)}{r^2} + \Xi_1 M''(r)$$
$$\frac{d\Psi}{dr} = \frac{G_N M(r)}{r^2} + \Xi_2 \frac{M'(r)}{r} + \Xi_3 M''(r)$$

- With EFT parameters: (1712.04731)

$$\Xi_1 = -\frac{1}{2} \frac{(\alpha_H + c_T^2 \beta_1)^2}{c_T^2 (1 + \alpha_V - 4\beta_1) - \alpha_H - 1}$$
$$\Xi_2 = -\frac{\alpha_H (\alpha_H - \alpha_V + 2 (1 + c_T^2) \beta_1) + \beta_1 (c_T^2 - 1) (1 + c_T^2 \beta_1)}{c_T^2 (1 + \alpha_V - 4\beta_1) - \alpha_H - 1}$$

$$\Xi_3 = -\frac{(\alpha_H + c_T^2 \beta_1)}{c_T^2 (1 + \alpha_V - 4\beta_1) - \alpha_H - 1}, \quad \gamma_0 = \alpha_V - 3\beta_1, \quad G_N = \frac{G}{1 + \gamma_0}$$

Our works: theoretical background

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- With EFT parameters (1712.04731) + constraints from GWs (1710.05834):

$$\Xi_1 = -\frac{(\alpha_H + \beta_1)^2}{2(\alpha_H + 2\beta_1)}, \quad \Xi_2 = \alpha_H,$$
$$\Xi_3 = -\frac{\beta_1(\alpha_H + \beta_1)^2}{2(\alpha_H + 2\beta_1)}, \quad \gamma_0 = -\alpha_H - 3\beta_1, \quad G_N = \frac{G}{1 + \gamma_0}$$

Our works: data - clusters of galaxies

- **galaxy clusters**: sample of 16 objects from CLASH
- why galaxy clusters?
 - in between cosmological and astrophysical scales
 - more sensitive to large scale DE effects
 - DM and DE effects must coexist
 - multi-messenger approach:
 - X-ray astronomy
 - optical

Our works: data - clusters of galaxies

- **galaxy clusters**: sample of 16 objects from CLASH
- galaxy clusters: X-ray observations \Rightarrow hot gas
 - spherical symmetry + **hydrostatic equilibrium** \Rightarrow collisionless Boltzmann eq:

$$-\frac{d\Phi(r)}{dr} = \frac{kT_{gas}(r)}{\mu m_p r} \left[\frac{d \ln \rho_{gas}(r)}{d \ln r} + \frac{d \ln T_{gas}(r)}{d \ln r} \right]$$

$$M_{tot}(r) = \frac{r^2}{G_N} \frac{d\Phi(r)}{dr} = -\frac{kT_{gas}(r)}{\mu m_p G_N} r \left[\frac{d \ln \rho_{gas}(r)}{d \ln r} + \frac{d \ln T_{gas}(r)}{d \ln r} \right]$$

$$M_{tot}(r) = M_{gas}(r) + M_{gal}(r) + M_{DM}(r) \quad (GR)$$

$$M_{tot}(r) = M_{gas}(r) + M_{gal}(r) + corr. \, terms \quad (DHOST)$$

- matter densities are provided by observations (gas, galaxies) or models (NFW dark matter)
- **caveat**: nongravitational effects can **bias** mass estimation

Our works: data - clusters of galaxies

- **galaxy clusters**: sample of 16 objects from CLASH
- galaxy clusters: optical \Rightarrow gravitational lensing
 - convergence map

$$\kappa(R) = \frac{1}{c^2} \frac{D_I D_{ls}}{D_s} \int_{-\infty}^{+\infty} \nabla_r \left(\frac{\Phi(R, z) + \Psi(R, z)}{2} \right) dz$$

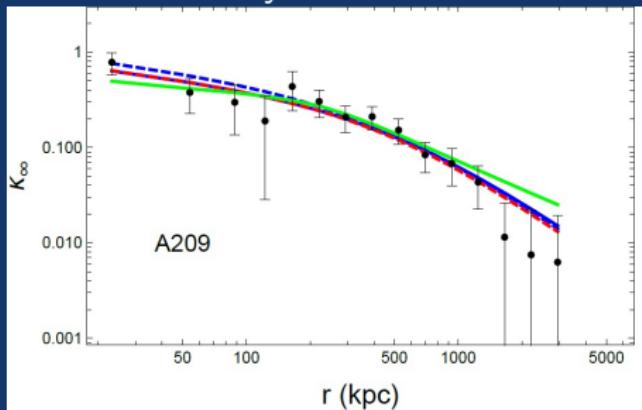
- **unbiased** mass estimation
- in GR: **tension** between X-ray and lensing mass
 - known thermal local disturbances contribute by 10 – 20%
 - tension is larger and with radial dependence
- how can we solve this tension?

Results: clusters of galaxies

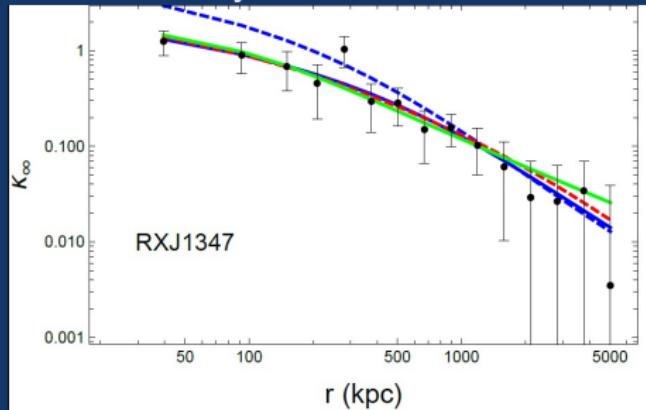
Scenario 1: DHOST as DE + DM

- DHOST as DE can solve tension!
- Statistically (Bayesian Evidence) favored w.r.t. GR!!

No X-ray vs GL tension



X-ray vs GL tension



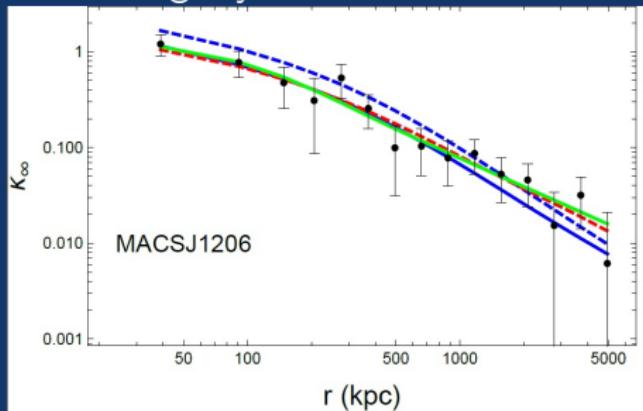
— GR:X+lens; — GR:lens; - - DHOST+DM; — DHOST-noDM

Results: clusters of galaxies

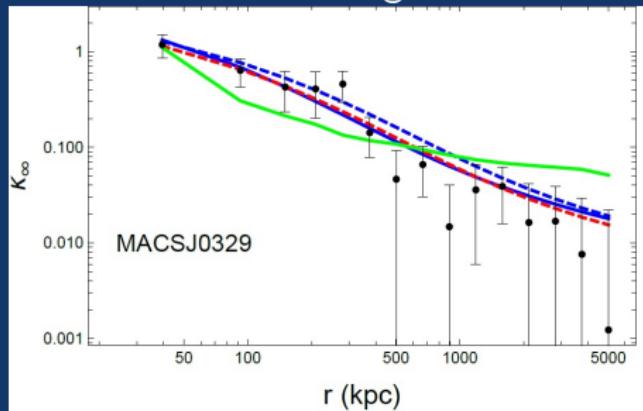
Scenario 2: DHOST as DE and DM

- Statistically disfavoured compared to GR in most of the cases
- Much more sensitive to mass components modeling!

Slightly better than GR



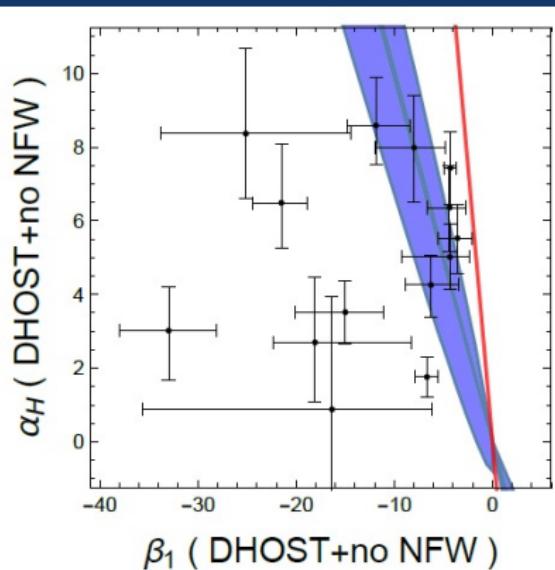
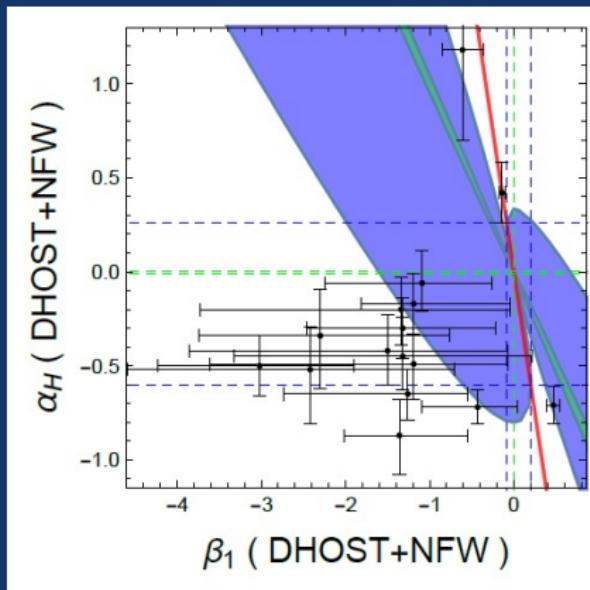
Mass modeling issues



— GR:X+lens; — GR:lens; - - DHOST+DM; — DHOST-noDM

Results: clusters of galaxies

- Blue lines: stellar stability + Hulse-Taylor pulsar (red line) (1712.04731)
- Green lines: helioseismology + Hulse-Taylor pulsar (red line) (1909.02552)



Our works: data - Ultra-diffuse galaxies

- **UDGs:** 3 from observations of the Dragonfly Telescope Array
 - NGC1052-DF2 (DF2) and NGC1052-DF4 (DF4): lack of DM ($\sim 1\%$)
 - Dragonfly44: DM dominated ($\sim 99\%$)
- mass modelling: stars + DM
- data: velocity dispersions \Rightarrow Jeans eq. (in GR)

$$\frac{d(\ell\sigma_r^2)}{dr} + \frac{\beta(r)}{r}(\ell\sigma_r^2) = \ell(r)\frac{d\Phi(r)}{dr}$$

$$\text{integrating} \quad \Rightarrow \quad \ell(r)\sigma_r^2 = \frac{1}{f(r)} \int_r^\infty ds f(s)\ell(s) \frac{M(s)}{s^2},$$

with $f(r) = d \log f(r) / d \log r = 2\beta(r)$. The observational quantity is

$$\sigma_{los}^2(R) = \frac{2}{I(R)} \left[\int_R^\infty dr r \frac{\ell(r)\sigma_r^2}{\sqrt{r^2 - R^2}} - R^2 \int_R^\infty dr \beta(r) \frac{\ell(r)\sigma_r^2}{\sqrt{r^2 - R^2}} \right]$$

Our works: data - Ultra-diffuse galaxies

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- mass modelling: stars + DM
- data: velocity dispersions \Rightarrow Jeans eq. (in DHOST)

$$\sigma_{los}^2(R) = \frac{2G_N}{I(R)} \int_R^\infty dr K\left(\frac{r}{R}\right) \ell(r) \frac{M(r)}{r} \quad (GR)$$

$$\sigma_{los}^2(R) = \frac{2G_N}{I(R)} \int_R^\infty dr K\left(\frac{r}{R}\right) \ell(r) \frac{M_{eff}(r)}{r} \quad (DHOST)$$

$$M_{eff}(r) = M(r) + r^2 \Xi_1 M''(r)$$

Our works: data - Ultra-diffuse galaxies

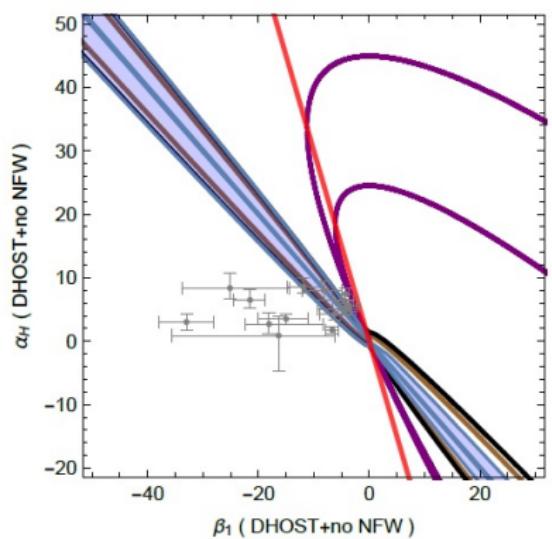
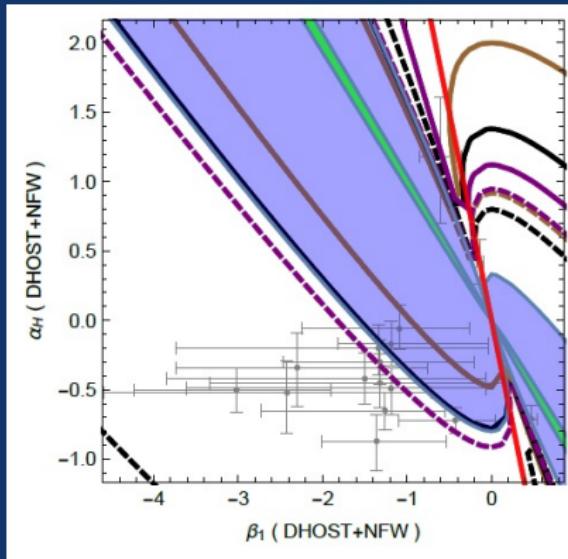
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- mass modelling: stars + DM
- data: velocity dispersions \Rightarrow Jeans eq. (in DHOST)
- Anisotropy parameter $\beta(r)$: two models (1501.03167)

$$\beta \equiv \beta_c ; \quad \beta(r) = \beta_0 + (\beta_\infty - \beta_0) \frac{r}{r + r_a}$$

- Relation between stars and DM
 - Stellar-to-halo mass relation (SHMR) (1703.04542)
 - Stellar and DM component decoupled

Results: data - Ultra-diffuse galaxies

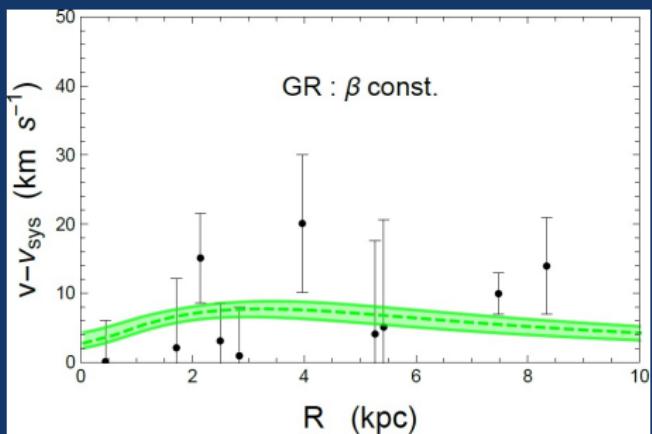
- EFT constraints: DF2 (black), DF4 (brown) and DF44 (purple)



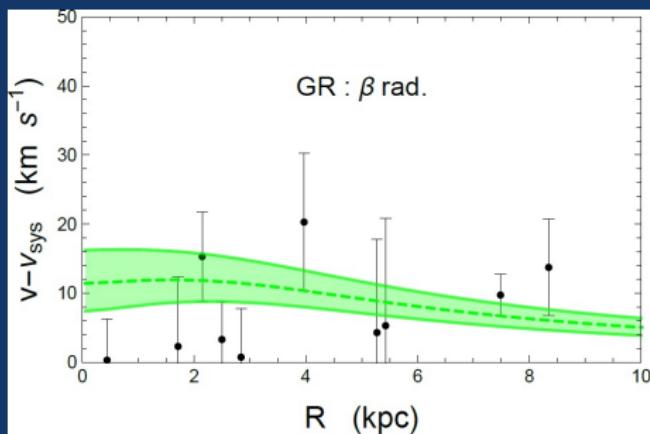
Results: data - Ultra-diffuse galaxies

- DF2: Baryonic case - GR

Reference case



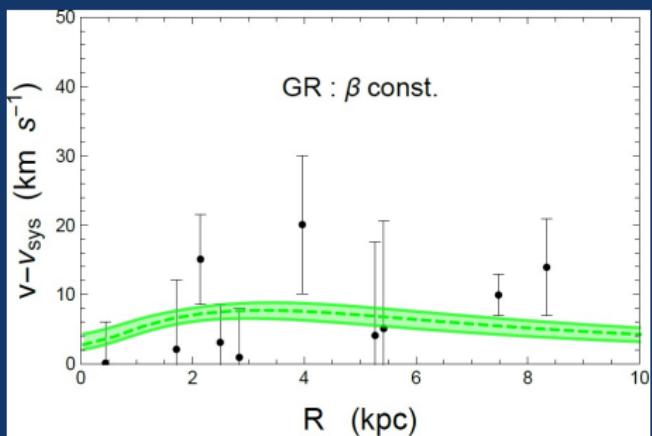
Statistically disfavoured



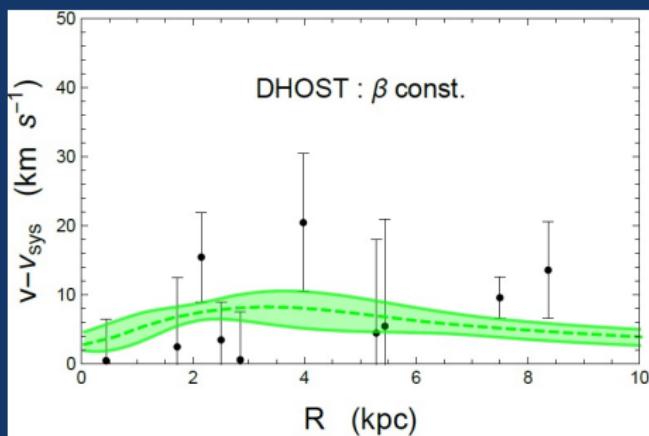
Results: data - Ultra-diffuse galaxies

- DF2: Baryonic case - GR vs DHOST (no DM!)

Reference case

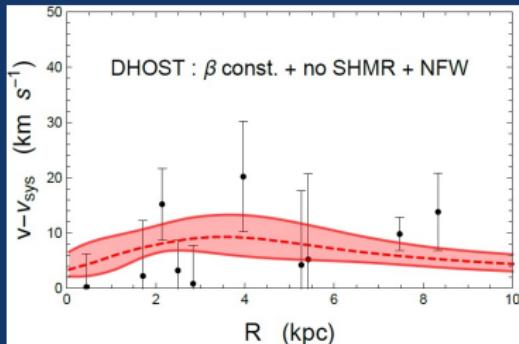
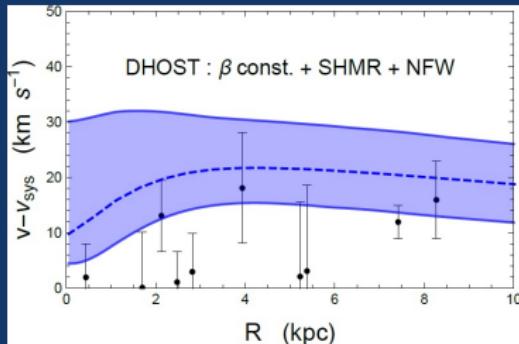
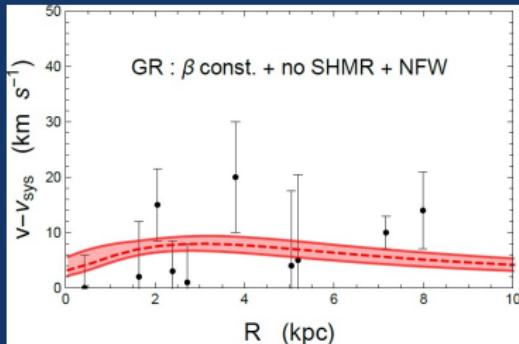
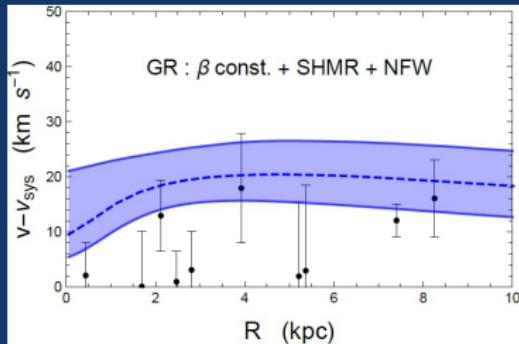


Statistically equivalent!



Results: data - Ultra-diffuse galaxies

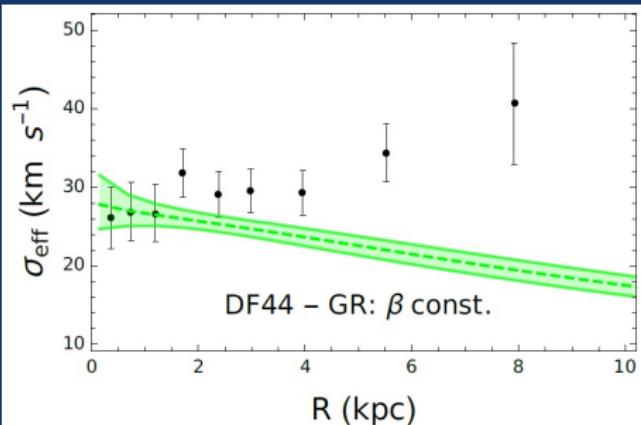
- DF2: DM case - GR vs DHOST



Results: data - Ultra-diffuse galaxies

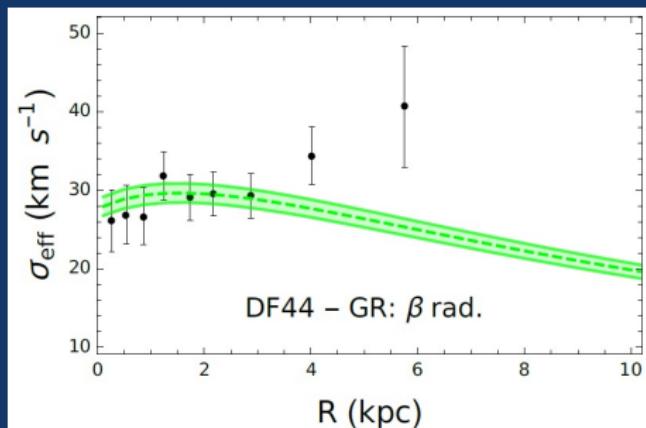
- Dragonfly44: Baryonic case - GR

Unable to fit data!



DF44 – GR: β const.

Unable to fit data!

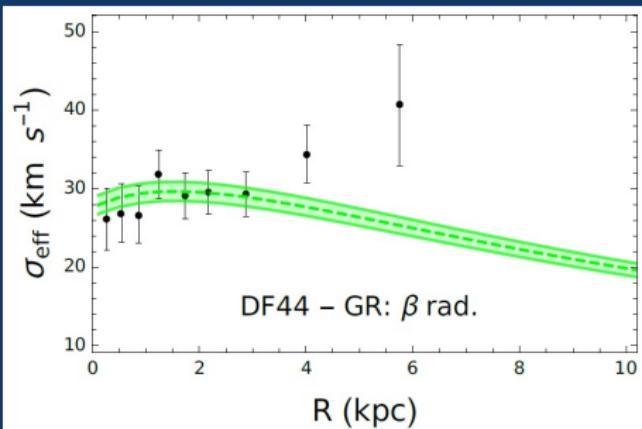


DF44 – GR: β rad.

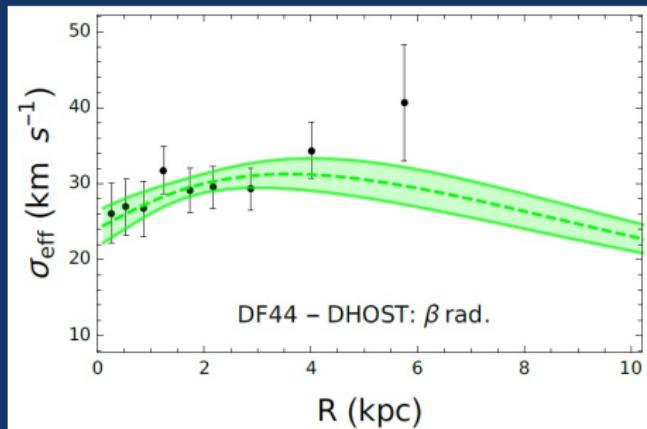
Results: data - Ultra-diffuse galaxies

- Dragonfly44: Baryonic case - GR vs DHOST (no DM!)

Unable to fit data!

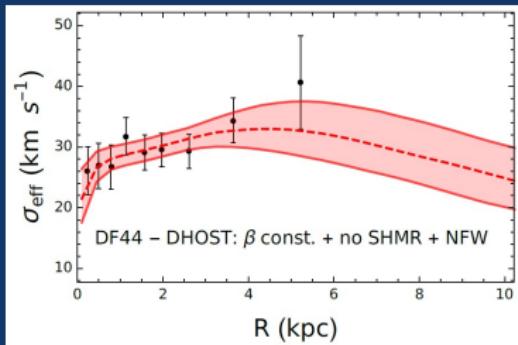
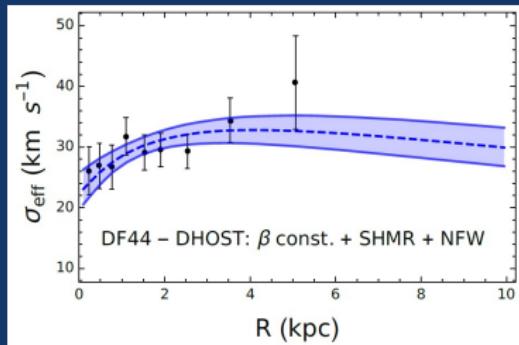
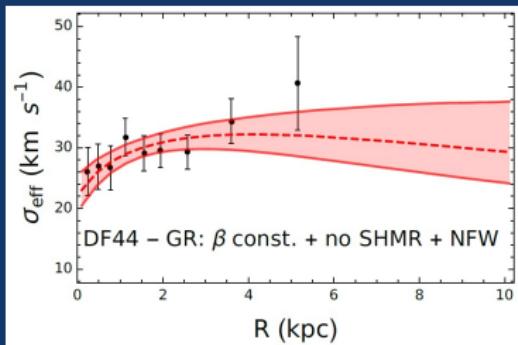
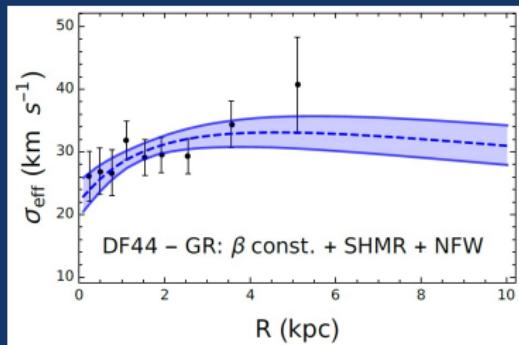


Better, but unable to fit data!



Results: data - Ultra-diffuse galaxies

- DF44: DM case - GR vs DHOST



Conclusions

From clusters of galaxies:

- Below 5 Mpc the screening mechanism is partially broken \Rightarrow DHOST can mimic the role of DE in galaxy clusters
- DHOST as DE can solve the natural tension between X-rays/lensing
- DHOST as DM is statistically disfavored compared to GR

From Ultra-diffuse galaxies:

- Lacking-DM galaxies (DF2 and DF4) are compatible with a pure baryonic mass modeling
- DHOST as DE does not produce significant changes w.r.t. GR \Rightarrow DE does not play a key role at galactic scales
- DHOST as DM
 - As successful as GR for lacking DM galaxies (DF2, DF4)
 - Unable to reproduce DM effects \Rightarrow DM is needed (Dragonfly44)

DZIĘKUJĘ BARDZO!