# Associated Higgs production at NLO with GoSam

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Alexander von Humboldt Stiftung/Foundation

## Publications

#### Based on the work described in:

- S. Alioli, S. Badger, ..., HvD, ... et al., "Update of the Binoth Les Houches Accord for a standard interface between Monte Carlo tools and one-loop programs" arXiv: 1308.3462
- HvD, G. Luisoni, P. Mastrolia, E. Mirabella, G. Ossola and T. Peraro, arXiv:1308.3462 [hep-ph]. "NLO QCD correction to Higgs boson production in association with a top quark pair and a jet", arXiv:1307.8437 [hep-ph].
- G. Cullen, HvD, N. Greiner, G. Luisoni, P. Mastrolia, E. Mirabella, G. Ossola., T. Peraro, and F. Tramontano, "NLO QCD corrections to Higgs boson production plus three jets in gluon fusion", arXiv:1307.4737 [hep-ph].
- S. Heinemeyer, C. Mariotti, ..., HvD, ... et al., "Handbook of LHC Higgs Cross Sections: 3. Higgs Properties", arXiv:1307.1347 [hep-ph]
- HvD, N. Greiner, G. Luisoni, P. Mastrolia, E. Mirabella, G. Ossola., T. Peraro, J.F. von Soden-Fraunhofen, and F. Tramontano, "NLO QCD corrections to the production of Higgs plus two jets at the LHC," Phys. Lett. B 721, 74 (2013).
- P. Mastrolia, E. Mirabella, G. Ossola., T. Peraro, HvD, "The Integrand Reduction of One- and Two-Loop Scattering Amplitudes" PoS LL2012 028 (2012).

# Outline

- Motivation
- GoSam
- Interfaces with external MCs
- Scattering amplitudes at one-loop
- Determining the parametric form of the numerator
- Extended rank numerator
- Higgs + 2 jets in GF
- Higgs + 3 jets in GF
- Higgs in association with top pair and jet
- Higgs plus 2 and 3 jets in VBF
- Conclusions and Outlook

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► Boson discovered by Atlas and CMS → Higgs?

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Need to determine properties: spin, CP properties, couplings

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Vector Boson Fusion



Gluon Fusion via top loop

Boson discovered by Atlas and CMS  $\rightarrow$  Higgs? Need to determine properties: ► spin, CP properties, couplings SM Higgs production 10 5 LHC  $\sigma[fb]$ Vector Boson Fusion  $gg \rightarrow h$  $10^{4}$ 10 3 → Wh  $z.aa \rightarrow tth$ 10<sup>2</sup> qb · → qth  $qq \rightarrow Zh$ TeV4LHC Higgs working group 300 400 500 100 200 m, [GeV] Gluon Fusion via top loop

# Motivation for NLO



- Reduce theoretical error
- Strong dependence on renormalization and factorization scale

# Motivation for NLO



- Reduce theoretical error
- Strong dependence on renormalization and factorization scale
- Development of more general framework for NLO automation

$$\sigma^{NLO} = \int_{m+1} \left[ d^{(4)} \sigma^R - d^{(4)} \sigma^A \right] + \int_m \left[ d^{(4)} \sigma^B + \int_{loop} d^{(d)} \sigma^V + \int_1 d^{(d)} \sigma^A \right]$$

$$\sigma^{NLO} = \int_{m+1} \left[ d^{(4)} \sigma^R - d^{(4)} \sigma^A \right] + \int_m \left[ d^{(4)} \sigma^B + \int_{loop} d^{(d)} \sigma^V + \int_1 d^{(d)} \sigma^A \right]$$

$$\uparrow$$

#### NLO cross section consists of:

Leading Order: Born diagram

$$\sigma^{NLO} = \int_{m+1} \left[ d^{(4)} \sigma^R - d^{(4)} \sigma^A \right] + \int_m \left[ d^{(4)} \sigma^B + \int_{loop} d^{(d)} \sigma^V + \int_1 d^{(d)} \sigma^A \right]$$

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- Leading Order: Born diagram
- Virtual corrections: loop diagrams

$$\sigma^{NLO} = \int_{m+1} \left[ d^{(4)} \sigma^R - d^{(4)} \sigma^A \right] + \int_m \left[ d^{(4)} \sigma^B + \int_{loop} d^{(d)} \sigma^V + \int_1 d^{(d)} \sigma^A \right]$$

$$\uparrow$$

- Leading Order: Born diagram
- Virtual corrections: loop diagrams
- Real corrections: Radiation

$$\begin{split} \sigma^{NLO} &= \int_{m+1} \left[ d^{(4)} \sigma^R - d^{(4)} \sigma^A \right] + \int_m \left[ d^{(4)} \sigma^B + \int_{loop} d^{(d)} \sigma^V + \int_1 d^{(d)} \sigma^A \right] \\ &\uparrow \qquad \uparrow \end{split}$$

- Leading Order: Born diagram
- Virtual corrections: loop diagrams
- Real corrections: Radiation
- Subtraction terms to regulate infinities

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- Leading Order: Born diagram
- Real corrections: Radiation
- Subtraction terms to regulate infinities

### GoSam

The GoSam collaboration:

G. Cullen, HvD, N. Greiner, G. Heinrich, G. Luisoni, P. Mastrolia, E. Mirabella, G. Oscola, T. Porera, J. Poiebel, J. Schlenk, J.F. von Soden-Fraunhofen,







 $\begin{array}{l} {\rm SAMURAI} \\ d-{\rm dimensional Integrand-Level Reduction} \rightarrow {\rm Current \ default} \\ {\rm Automated \ Model-independent \ Computation \ of \ the \ full \ Rational \ Term} \\ {\rm Mastrolia, \ Ossola, \ Reiter, \ Tramontano} \end{array}$ 

 $\label{eq:GOLEM95} GOLEM95$  Tensorial Reduction  $\rightarrow$  Rescue System Binoth, Guillet, Heinrich, Pilon, Reiter

 $\label{eq:NINJA} \begin{array}{l} \mathsf{NINJA} \to \mathsf{Talk} \text{ of } \mathsf{T}. \text{ Peraro} \\ \\ \mathsf{Integrand}\text{-}\mathsf{Level} \ \mathsf{Reduction} + \mathsf{Laurent} \ \mathsf{Expansion} \to \mathsf{Stable} \ \mathsf{and} \ \mathsf{Fast}!! \\ \\ \\ \mathsf{Mastrolia}, \ \mathsf{Mirabella}, \ \mathsf{Peraro} \end{array}$ 



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Associated Higgs production at NLO with GoSam + c1235678,0

# **Binoth Les Houches Accord**



"Update of the Binoth Les Houches Accord for a standard interface between Monte Carlo tools and one-loop programs", arXiv: 1308.3462

[Alioli, Badger, ..., HvD, ... et al. (2013)]

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# Interfaces with external MC

GoSam + MadGraph+MadDipole+MadEvent → ad-hoc interface [Greiner]

• GoSam + Sherpa  $\rightarrow via BLHA$ 

 GoSam + POWHEG → via BLHA [Luisoni, Nason, Oleari, Tramontano]

# Interfaces with external MC

► GoSam + MadGraph+MadDipole+MadEvent → ad-hoc interface [Greiner]

• GoSam + Sherpa  $\rightarrow via BLHA$ 

 GoSam + POWHEG → via BLHA [Luisoni, Nason, Oleari, Tramontano]

GoSam + HERWIG

→ work in progress [Greiner, Heinrich, von Soden-Fraunhofen]

#### GoSam + AMC@NLO

 $\rightarrow$  work in progress

[HvD, Frederix, Frixione, Hirschi, Luisoni, Mastrolia, Ossola, Peraro]









$$\mathcal{M}_n \equiv \int \mathcal{A}_n(\bar{q}) \, d\bar{q} \equiv \int d^{-2\epsilon} \mu \int d^4 q \frac{N(q,\mu^2)}{\bar{D}_0 \dots \bar{D}_{n-1}}$$



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• Decompose:  $c_{5,0} \xrightarrow{d+2} + c_{4,0} \xrightarrow{d+4} + c_{3,0} \xrightarrow{d+4} + c_{3,7} \xrightarrow{d+4} + c_{2,0} \xrightarrow{d+4} + c_{2,0} \xrightarrow{d+4} + c_{1,0} \xrightarrow{d+4} + \int d\bar{q} \frac{c_{5,0}\mu^2}{D_0 D_1 D_2 D_3 D_4} + \int d\bar{q} \frac{c_{4,0} + c_{4,4}\mu^4}{D_0 D_1 D_2 D_3} + \int d\bar{q} \frac{c_{3,0} + c_{3,7}\mu^2}{D_0 D_1 D_2} + \int d\bar{q} \frac{c_{2,0} + c_{2,9}\mu^2}{D_0 D_1} + \int d\bar{q} \frac{c_{1,0}}{D_0}$ 

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$$\mathcal{M}_n \equiv \int \mathcal{A}_n(\bar{q}) \, d\bar{q} \equiv \int d^{-2\epsilon} \mu \int d^4q \frac{N(q,\mu^2)}{\bar{D}_0 \dots \bar{D}_{n-1}}$$

• Decompose:  $c_{5,0} \xrightarrow{d+2} + c_{4,0} \xrightarrow{d+4} + c_{3,0} \xrightarrow{d+6} + c_{3,7} \xrightarrow{d+6} + c_{2,0} \xrightarrow{d+6} + c_{2,9} \xrightarrow{d+4} + c_{1,0} \xrightarrow{d+6} + c_{1,0}$ 

• computation of  $\mathcal{M}_n \rightarrow$  computation of coefficients

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$$c_{5,0} \xrightarrow{\mathbf{d}+2} + c_{4,0} \xrightarrow{\mathbf{d}+4} + c_{4,4} \xrightarrow{\mathbf{d}+4} + c_{3,0} \xrightarrow{\mathbf{d}+4} + c_{3,7} \xrightarrow{\mathbf{d}+4} + c_{2,0} \xrightarrow{\mathbf{d}+4} + c_{2,0} \xrightarrow{\mathbf{d}+4} + c_{1,0} \xrightarrow{\mathbf{d}+4} + c_{1$$

$$c_{5,0} \xrightarrow{\mathbf{d}+2} + c_{4,0} \xrightarrow{\mathbf{d}+4} + c_{4,4} \xrightarrow{\mathbf{d}+4} + c_{3,0} \xrightarrow{\mathbf{d}+4} + c_{3,7} \xrightarrow{\mathbf{d}+4} + c_{2,0} \xrightarrow{\mathbf{d}+4} + c_{2,0} \xrightarrow{\mathbf{d}+4} + c_{1,0} \xrightarrow{\mathbf{d}+4} + \int d\bar{q} \frac{c_{5,0}\mu^2}{D_0 D_1 D_2 D_3 D_4} + \int d\bar{q} \frac{c_{4,0} + c_{4,4}\mu^4}{D_0 D_1 D_2 D_3} \\ + \int d\bar{q} \frac{c_{3,0} + c_{3,7}\mu^2}{D_0 D_1 D_2} + \int d\bar{q} \frac{c_{2,0} + c_{2,9}\mu^2}{D_0 D_1} + \int d\bar{q} \frac{c_{1,0}}{D_0}$$

• integral  $\rightarrow$  integrand:

$$\int d^{-2\epsilon} \mu^{2} d^{4}q \mathcal{A}_{n}(q) = \int d\bar{q} \frac{c_{5,0}\mu^{2}}{D_{0}D_{1}D_{2}D_{3}D_{4}} + \int d\bar{q} \frac{c_{4,0} + c_{4,4}\mu^{4}}{D_{0}D_{1}D_{2}D_{3}} + \int d\bar{q} \frac{c_{4,0} + c_{4,4}\mu^{4}}{D_{0}D_{1}D_{2}D_{3}} + \int d\bar{q} \frac{c_{3,0} + c_{3,7}\mu^{2}}{D_{0}D_{1}D_{2}} + \int d\bar{q} \frac{c_{2,0} + c_{2,9}\mu^{2}}{D_{0}D_{1}} + \int d\bar{q} \frac{c_{1,0}}{D_{0}}$$

► integral → integrand:  

$$A_n(q) = \frac{c_{5,0}\mu^2 + f_{01234}(q,\mu^2)}{D_0 D_1 D_2 D_3 D_4} + \frac{c_{4,0} + c_{4,4}\mu^4 + f_{0123}(q,\mu^2)}{D_0 D_1 D_2 D_3} + \frac{c_{3,0} + c_{3,7}\mu^2 + f_{012}(q,\mu^2)}{D_0 D_1 D_2} + \frac{c_{2,0} + c_{2,9}\mu^2 + f_{01}(q,\mu^2)}{D_0 D_1} + \frac{c_{1,0} + f_0(q,\mu^2)}{D_0}$$

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$$\int d^{-2\epsilon} \mu^{2} d^{4} q \mathcal{A}_{n}(q) = \int d\bar{q} \frac{c_{5,0}\mu^{2}}{D_{0}D_{1}D_{2}D_{3}D_{4}} + \int d\bar{q} \frac{c_{4,0} + c_{4,4}\mu^{4}}{D_{0}D_{1}D_{2}D_{3}} + \int d\bar{q} \frac{c_{3,0} + c_{3,7}\mu^{2}}{D_{0}D_{1}D_{2}} + \int d\bar{q} \frac{c_{2,0} + c_{2,9}\mu^{2}}{D_{0}D_{1}} + \int d\bar{q} \frac{c_{1,0}}{D_{0}}$$

► integral → integrand:  

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$$\int d^{-2\epsilon} \mu^2 \int d^4 q \quad \frac{f_{ij\dots}(q,\mu^2)}{D_i D_j \dots} = 0$$

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# Determining the parametric form of the numerator

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### Determining the parametric form of the numerator

$$\mathcal{A}_n = \sum_{ijklm} \frac{\Delta_{ijklm}(q,\mu^2)}{D_i D_j D_k D_l D_m} + \sum_{ijkl} \frac{\Delta_{ijkl}(q,\mu^2)}{D_i D_j D_k D_l} + \sum_{ijk} \frac{\Delta_{ijk}(q,\mu^2)}{D_i D_j D_k} + \sum_{ij} \frac{\Delta_{ij}(q,\mu^2)}{D_i D_j} + \sum_i \frac{\Delta_i(q,\mu^2)}{D_i} + \sum_i \frac{\Delta_i$$

- Form residues process independent
- Values of coefficients process dependent

### Determining the parametric form of the numerator

$$\mathcal{A}_n = \sum_{ijklm} \frac{\Delta_{ijklm}(q,\mu^2)}{D_i D_j D_k D_l D_m} + \sum_{ijkl} \frac{\Delta_{ijkl}(q,\mu^2)}{D_i D_j D_k D_l} + \sum_{ijk} \frac{\Delta_{ijk}(q,\mu^2)}{D_i D_j D_k} + \sum_{ij} \frac{\Delta_{ij}(q,\mu^2)}{D_i D_j} + \sum_i \frac{\Delta_i(q,\mu^2)}{D_i} + \sum_i \frac{\Delta_i$$

- Form residues process independent
- Values of coefficients process dependent
- Implemented in Samurai [Ossola, Reiter, Tramontano, Mastrolia, 2010]

# Rankcounting, normal rank



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 Only q propagators and 3-gluon-vertices contribute one power of q to numerator

# Rankcounting, normal rank



 Only q propagators and 3-gluon-vertices contribute one power of q to numerator


# Rankcounting, normal rank



 Only q propagators and 3-gluon-vertices contribute one power of q to numerator



 $\blacktriangleright$   $r_N \leq \#D$ 

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$$\Delta_{ijk\ell m}(\bar{q}) = \operatorname{Res}_{ijk\ell m} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\}$$
 1 coefficient

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$$\Delta_{ijk\ell m}(\bar{q}) = \text{Res}_{ijk\ell m} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} \right\}$$

1 coefficient



$$\Delta_{ijk\ell m}(\vec{q}) = \operatorname{Res}_{ijk\ell m} \left\{ \frac{N(\vec{q})}{D_0 \cdots D_{n-1}} \right\}$$

$$1 \text{ coefficient}$$

$$\Delta_{ijk\ell}(\vec{q}) = \operatorname{Res}_{ijk\ell} \left\{ \frac{N(\vec{q})}{D_0 \cdots D_{n-1}} - \sum_{i < m}^{n-1} \frac{\Delta_{ijk\ell m}(\vec{q})}{D_i D_i D_j D_k D_\ell D_m} \right\}$$

$$5 \text{ coefficients}$$

$$\Delta_{ijk}(\bar{q}) = \operatorname{Res}_{ijk} \left\{ \frac{N(\bar{q})}{\bar{D}_0 \cdots \bar{D}_{n-1}} - \sum_{i < < m}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell \bar{D}_m} - \sum_{i < < \ell}^{n-1} \frac{\Delta_{ijk\ell}(\bar{q})}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} \right\} \quad 10 \text{ coefficients}$$

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$$\Delta_{ijk\ell m}(\bar{q}) = \operatorname{Res}_{ijk\ell m} \left\{ \frac{N(\bar{q})}{D_0 \cdots D_{n-1}} \right\}$$
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5 coefficients
$$\Delta_{ijk}(\bar{q}) = \operatorname{Res}_{ijk} \left\{ \frac{N(\bar{q})}{D_0 \cdots D_{n-1}} - \sum_{k < m}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_j D_k D_k D_m} - \sum_{k < m}^{n-1} \frac{\Delta_{ijk\ell}(\bar{q})}{D_i D_j D_k D_k D_m} \right\}$$
10 coefficients

$$- \underbrace{ \Delta_{ij}(\bar{q}) = \operatorname{Res}_{ij} \left\{ \frac{N(\bar{q})}{D_0 \cdots D_{n-1}} - \sum_{i < < m}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_j D_k D_\ell D_m} - \sum_{i < < k}^{n-1} \frac{\Delta_{ijk\ell}(\bar{q})}{D_i D_j D_k D_\ell} - \sum_{i < k}^{n-1} \frac{\Delta_{ijk}(\bar{q})}{D_i D_j D_k} \right\} \quad 10 \text{ coefficients}$$

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$$\Delta_{ijk\ell m}(\vec{q}) = \operatorname{Res}_{ijk\ell m}\left\{\frac{N(\vec{q})}{D_0 \cdots D_{n-1}}\right\}$$
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$$\Delta_{ijk}(\vec{q}) = \operatorname{Res}_{ijk}\left\{\frac{N(\vec{q})}{D_0 \cdots D_{n-1}} - \sum_{i < m}^{n-1} \frac{\Delta_{ijk\ell m}(\vec{q})}{D_i D_k D_k D_k D_m} - \sum_{i < m}^{n-1} \frac{\Delta_{ijk\ell}(\vec{q})}{D_i D_j D_k D_\ell D_m}\right\}$$
10 coefficients
$$\Delta_{ijk}(\vec{q}) = \operatorname{Res}_{ijk}\left\{\frac{N(\vec{q})}{D_0 \cdots D_{n-1}} - \sum_{i < m}^{n-1} \frac{\Delta_{ijk\ell m}(\vec{q})}{D_i D_j D_k D_\ell D_m} - \sum_{i < m}^{n-1} \frac{\Delta_{ijk\ell}(\vec{q})}{D_i D_j D_k D_\ell D_m}\right\}$$
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$$\underbrace{ \Delta_i(\vec{q}) = \operatorname{Res}_i \left\{ \frac{N(\vec{q})}{D_0 \cdots D_{n-1}} - \sum_{i < < m}^{n-1} \frac{\Delta_{ijk\ell m}(\vec{q})}{D_i D_j D_k D_\ell D_m} - \sum_{i < < \ell}^{n-1} \frac{\Delta_{ijk\ell}(\vec{q})}{D_i D_j D_k D_\ell} + \sum_{i < < \ell}^{n-1} \frac{\Delta_{ijk\ell}(\vec{q})}{D_i D_j D_k} - \sum_{i < < \ell}^{n-1} \frac{\Delta_{ij\ell}(\vec{q})}{D_i D_j} \right\}$$
 5 coefficients

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$$\Delta_{ijk\ell m}(\bar{q}) = \operatorname{Res}_{ijk\ell m}\left\{\frac{N(\bar{q})}{D_{0}\cdots D_{n-1}}\right\}$$

$$1 \text{ coefficient}$$

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$$5 \text{ coefficients}$$

$$-\sum_{k$$$$$$$$$$$$

Hexagon: 
$$\binom{6}{5} \cdot 1 + \binom{6}{4} \cdot 5 + \binom{6}{3} \cdot 10 + \binom{6}{2} \cdot 10 + \binom{6}{1} \cdot 5 = 386$$
 coefficients

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#### Rankcounting, higher rank



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#### Rankcounting, higher rank



• One effective vertex: 
$$r_N \leq \#D + 1$$

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#### Extended rank Integrand decomposition algorithm

$$\Delta_{ijk\ell m}(\bar{q}) = \operatorname{Res}_{ijk\ell m} \left\{ \frac{N(\bar{q})}{D_0 \cdots D_{n-1}} \right\} \qquad 1 \rightarrow 1 \text{ coefficients}$$

$$\Delta_{ijk\ell}(\bar{q}) = \operatorname{Res}_{ijk\ell} \left\{ \frac{N(\bar{q})}{D_0 \cdots D_{n-1}} - \sum_{i

$$\Delta_{ijk\ell}(\bar{q}) = \operatorname{Res}_{ijk\ell} \left\{ \frac{N(\bar{q})}{D_0 \cdots D_{n-1}} - \sum_{i

$$\Delta_{ij}(\bar{q}) = \operatorname{Res}_{ij} \left\{ \frac{N(\bar{q})}{D_0 \cdots D_{n-1}} - \sum_{i

$$- \sum_{i<\ell} \frac{\Delta_{ij}(\bar{q})}{D_i D_i D_m} - \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_j D_k D_\ell D_m} - \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_j D_k D_\ell D_m} + 5 \rightarrow 15 \text{ coefficients}$$

$$- \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk}(\bar{q})}{D_i D_j D_k} - \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_j D_k D_\ell} - \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_j D_k D_\ell D_m} + 5 \rightarrow 15 \text{ coefficients}$$

$$- \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk}(\bar{q})}{D_i D_j D_k} - \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_j D_k D_\ell} + \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_i D_k D_\ell} + \sum_{i<\ell}^{n-1} \frac{\Delta_{ijk\ell m}(\bar{q})}{D_k D_\ell} + \sum_{i<\ell$$$$$$$$

Samurai  $\rightarrow$  XSamurai [HvD et al.]

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#### Associated Higgs production at NLO with GoSam

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 $\blacktriangleright \ \Delta(q,\mu^2)$  multivariate polynomial in q and  $\mu^2$ 

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- Systematic sampling: DFT

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$$P(x) = c_0 + c_1 x + c_2 x^2 + \dots + c_n x^n$$

- $\blacktriangleright \ \Delta(q,\mu^2)$  multivariate polynomial in q and  $\mu^2$
- Systematic sampling: DFT

$$P(x) = c_0 + c_1 x + c_2 x^2 + \dots + c_n x^n$$

$$x_k = \rho \exp\left[-2\pi i \frac{k}{n+1}\right]$$



- $\Delta(q,\mu^2)$  multivariate polynomial in q and  $\mu^2$
- Systematic sampling: DFT

$$P(x) = c_0 + c_1 x + c_2 x^2 + \dots + c_n x^n$$
$$x_k = \rho \exp\left[-2\pi i \frac{k}{n+1}\right]$$
$$P_k = P(x_k) = \sum_{l=0}^n c_l \rho^l \exp\left[-2\pi i \frac{k}{(n+1)}l\right]$$

- $\Delta(q,\mu^2)$  multivariate polynomial in q and  $\mu^2$
- Systematic sampling: DFT

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$$c_l = \frac{\rho^{-l}}{n+1} \sum_{k=0}^n P_k \exp\left[2\pi i \frac{k}{n+1}l\right]$$

[Mastrolia, Ossola, Papadopoulos, Pittau (2008)]

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$$q = \sum_{i=1}^{4} x_i e_i \Rightarrow \mu^2, x_1, x_2, x_3, x_4$$
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- At triple cut:  $\Delta(\mu^2, x_3, x_4)$  Condition:  $x_3x_4 = C(x_1, x_2) = C \Rightarrow \Delta(\mu^2, x_3, C/x_3)$

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  - ▶ Use DFT twice,  $\Delta(\mu^2, x3, C/x3)$  and  $\Delta(\mu^2 C/x_4, x_4)$  solutions  $\propto \frac{1}{1-C}$ , problem if C = 1

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- ▶ Branching: if(C=0): Use  $\Delta(\mu^2, x_3, C/x_3)$  and  $\Delta(\mu^2 C/x_4, x_4)$ else: Use  $\Delta(\mu^2, x_3, C/x_3)$

At double cut:  $\Delta(\mu^2, x_1, x_3, x_4)$  with  $x_3x_4 = F(x_1) = Ax_1^2 + Bx_1 + C$  lot of branchings:

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- At single cut: Δ(μ<sup>2</sup>, x<sub>1</sub>, x<sub>2</sub>, x<sub>3</sub>, x<sub>4</sub>) with x<sub>3</sub>x<sub>4</sub> x<sub>1</sub>x<sub>2</sub> = G similar to the triple cut

## Towards Higgs jets in GF @ NLO

H+0j	1 NLO
$gg \rightarrow H$	1 NLO
H+1j	62 NLO
$qq \rightarrow Hqq$	14 NLO
$qg \rightarrow Hqg$	48 NLO
H+2j	926 NLO
$qq' \rightarrow Hqq'$	32 NLO
qq  ightarrow Hqq	64 NLO
$qg \rightarrow Hqg$	179 NLO
$gg \rightarrow Hgg$	651 NLO
H+3j	13179 NLO
$qq' \rightarrow Hqq'g$	467 NLO
$qq \rightarrow Hqqg$	868 NLO
$qg \rightarrow Hqgg$	2519 NLO
$gg \rightarrow Hggg$	9325 NLO

Computational Challenges:

- Over 10,000 diagrams
- Higher-Rank terms
- 60 Rank-7 hexagons



#### $\text{Complex calculations} \rightarrow \text{GoSam enhanced}$

grouping, optimalization through Form4.0, numerical polarization vectors, parallelization

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#### Higgs + 2 jets in GF @ NLO

- Results obtained with GoSam+Sherpa
- Agreement with MCFM (v6.4) [Campbell, Ellis, Williams]



HvD, Greiner, Luisoni, Mastrolia, Mirabella, Ossola, Peraro, von Soden-Fraunhofen, Tramontano (2013) (also appeared in *Handbook of LHC Higgs Cross Sections: 3. Higgs Properties*)

### Higgs + 3 jets in GF: virtual part

Cullen, HvD, Greiner, Luisoni, Mastrolia, Mirabella, Ossola, Peraro, Tramontano, **arXiv:1307.4737** 

# Virtual parts computed with GoSam

• • • • • • • • • • • • • • • • • • •		00000000000000000000000000000000000000	$\frac{2 \operatorname{Ste}\left\{\mathcal{M}^{\operatorname{Hee-level}},\mathcal{M}^{\operatorname{Hee-level}}\right\}}{\left(\alpha_s/2\pi\right)\left \mathcal{M}^{\operatorname{tree-level}}\right ^2} \equiv \frac{a_{-2}}{\epsilon^2} + \frac{a_{-1}}{\epsilon} + a_0$
	<u>3</u>		$\begin{array}{cccccccccccccccccccccccccccccccccccc$
SUBPROCESS	DIAGRAMS	TIME/PS-POINT [sec]	*
$q\bar{q} \rightarrow Hq'\bar{q}'g$	467	0.29	1 XV/ XV/1
$q\bar{q} \rightarrow Hq\bar{q}g$	868	0.60	50
$gg \rightarrow Hq\bar{q}g$	2519	3.9	<sup>-ov</sup> [ Y Y]
$gg \rightarrow Hggg$	9325	20	$\frac{1}{\pi/2}$ $\frac{1}{\pi}$ $\frac{1}{3\pi/2}$ $\frac{2\pi}{2\pi}$
			Angle $\theta$ around $y$ -axis

#### Number of Feynman diagrams and time per

PS-point point for each subprocess

#### Tests: gauge invariance and IR poles

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#### Higgs + 3 jets GF @ NLO: cross-section

Cross sections are obtained with a hybrid setup:

- GoSam + Sherpa for Born and of the virtual contributions
- MadGraph+MadDipole+MadEvent for reals/subtraction/integrated dipoles



$$\mu_F = \mu_R = \frac{\hat{H}_T}{2} = \mu_0$$

$$\hat{H}_T = \sqrt{m_H^2 + p_{T,H}^2} + \sum_i |p_{T,i}|$$

Tests performed on the cross section:

- NLO H+2 jets: Agreement between hybrid scheme and GoSam+Sherpa
- LO H+3 jets: Agreement between MADGRAPH and Sherpa
- ► NLO H+3 jets: Independence from α-parameter (subtraction+int. dipoles)

Cullen, HvD, Greiner, Luisoni, Mastrolia, Mirabella, Ossola, Peraro, Tramontano arXiv:1307.4737

### Higgs + 3 jets GF @ NLO: distributions



 $pp \rightarrow Hjjj$  generated by GoSam can be paired with available MC programs for further phenomenological analyses.

Cullen, HvD, Greiner, Luisoni, Mastrolia, Mirabella, Ossola, Peraro, Tramontano arXiv:1307.4737

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#### $pp \rightarrow Ht\bar{t}$ + 1 jet @ NLO

First Application of GOSAM/NINJA + SHERPA  $\rightarrow$  Talk of T. Peraro

$t\bar{t}H + 1j$	1895 NLO
$qq \rightarrow H t \bar{t} g$	320 NLO
$gg \to H t \bar{t} g$	1575 NLO

- Two different mass scales: Higgs and Top
- 51 hexagons in the gluon-gluon channel



#### Htītj results





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## Vector Boson Fusion

H+2j	240 NLO
$us \rightarrow Hdc$	24 NLO
$uc \rightarrow Huc$	24 NLO
$us \rightarrow Hus$	24 NLO
$ds \rightarrow Hds$	24 NLO
$ud \rightarrow Hud$	48 NLO
$uu \rightarrow Huu$	48 NLO
$dd \rightarrow H dd$	48 NLO
H+3j	2160 NLO
$\begin{array}{c} \textbf{H+3j} \\ us \rightarrow Hdcg \end{array}$	2160 NLO 216 NLO
$\begin{array}{c} \textbf{H+3j} \\ us \rightarrow Hdcg \\ uc \rightarrow Hucg \end{array}$	2160 NLO 216 NLO 216 NLO
$H+3j$ $us \rightarrow Hdcg$ $uc \rightarrow Hucg$ $us \rightarrow Husg$	2160 NLO 216 NLO 216 NLO 216 NLO
$H+3j$ $us \rightarrow Hdcg$ $uc \rightarrow Hucg$ $us \rightarrow Husg$ $ds \rightarrow Hdsg$	2160 NLO 216 NLO 216 NLO 216 NLO 216 NLO
$H+3j$ $us \rightarrow Hdcg$ $uc \rightarrow Hucg$ $us \rightarrow Husg$ $ds \rightarrow Hdsg$ $ud \rightarrow Hudg$	2160 NLO 216 NLO 216 NLO 216 NLO 216 NLO 432 NLO
$\begin{array}{c} \textbf{H+3j}\\ \hline us \rightarrow Hdcg\\ uc \rightarrow Hucg\\ us \rightarrow Husg\\ ds \rightarrow Hdsg\\ ud \rightarrow Hudg\\ uu \rightarrow Huug \end{array}$	2160 NLO 216 NLO 216 NLO 216 NLO 216 NLO 432 NLO 432 NLO


## Conclusions

- GoSam is a powerful framework for the automatic computation of one loop virtual amplitudes
- Samurai has been extended to deal with higher rank numerators: Xsamurai
- GoSam is interfaced to a lot of Monte Carlo Programs
- GoSam has been used for a lot of different phenomenology studies, among which
  - Higgs plus two jet in GF
  - Higgs plus three jet in GF
  - Higgs  $t\bar{t}$  jet
  - Higgs plus two and three jet in VBF in progress

## Outlook

- Even more processes on the way
- Interaction with MC and experimental collaborations
- Additional code improvements towards GoSam 2.0
- Multi-loop integrand reduction in the making (talk T.Peraro on Thursday)



## **BACKUP SLIDES**

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Associated Higgs production at NLO with GoSam