

1. Motivation

It is known that time-dependent vacuum expectation value of the background field in the presence of perturbative corrections to non-perturbative production causes the production of particles in

- **non-adiabatic area** (vacuum change, rotation of the basis, rescattering)
- **adiabatic area** (perturbation theory)

2. Fundamentals

Bogoliubov transformation:

$$a_k^{\text{out}} = \alpha_k a_k^{\text{in}} + \beta_k a_k^{\text{in}\dagger}$$

$$a_k^{\text{out}\dagger} = \alpha_k^* a_k^{\text{in}\dagger} + \beta_k^* a_k^{\text{in}}$$

Occupation number of produced particles:

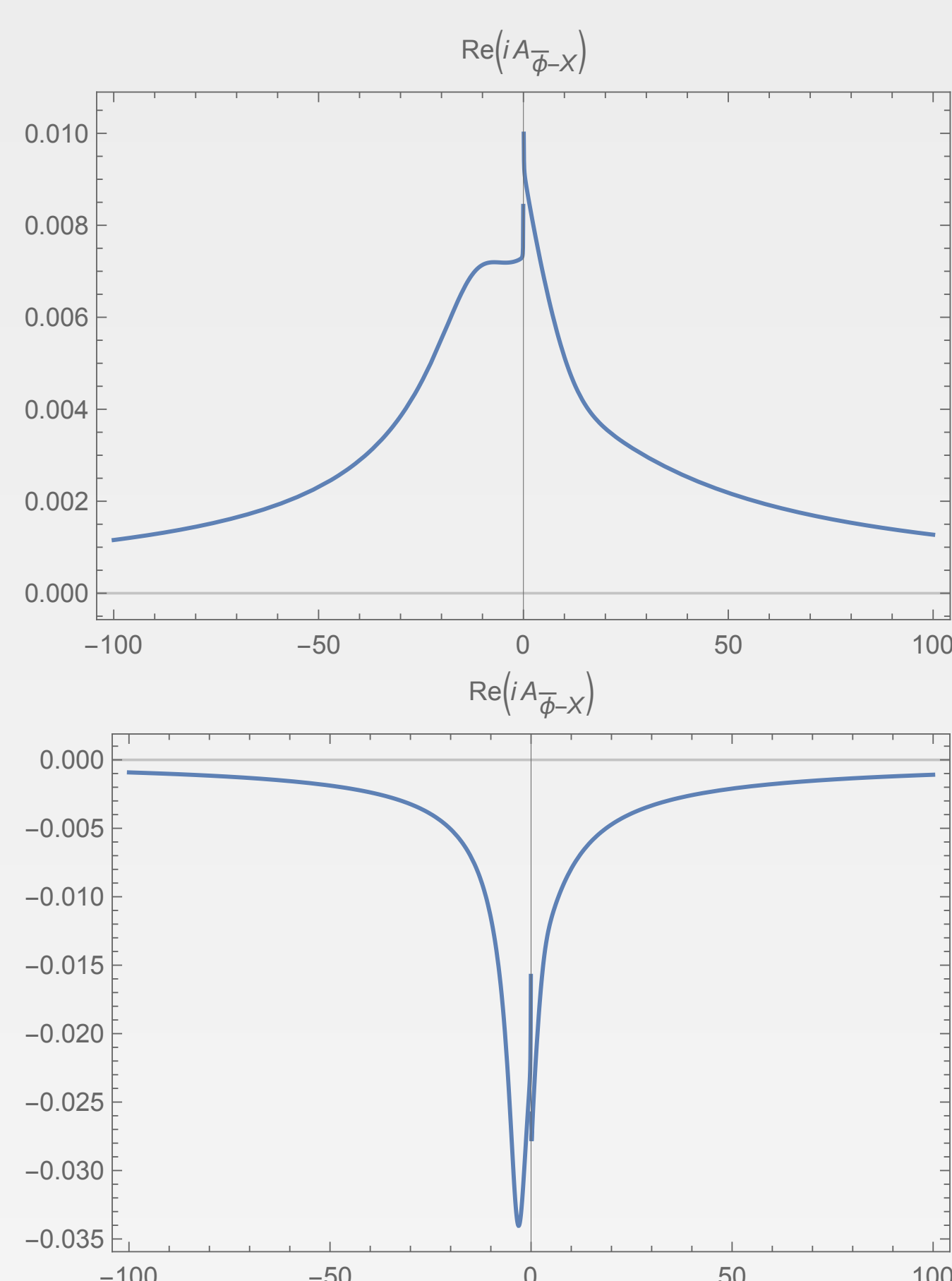
$$n_k \equiv \langle 0^{\text{in}} | N_k | 0^{\text{in}} \rangle = \langle 0^{\text{in}} | a_k^{\text{out}\dagger} a_k^{\text{out}} | 0^{\text{in}} \rangle = V |\beta_k|^2$$

Adiabaticity:

- **adiabatic area:** $\dot{\omega}_k / \omega_k^2 < 1$
 $n_k(t) \approx \frac{\rho_k}{\omega_k} \approx \frac{1}{\omega_k} \left(\sqrt{\omega_k e^{\pm i \int \omega}} \right)^2 \approx \text{const}$
- **non-adiabatic area:** $\dot{\omega}_k / \omega_k^2 > 1$
 $n_k(t) \neq \text{const}$
 \Rightarrow **particle production**

5. Rotation of the basis

When $\langle \chi \rangle \neq 0$ we get additional production due to rotation of the basis for fermionic mass eigenstates.



$$\langle \chi \rangle = v_\chi t + i\mu_\chi$$

above: $v_\chi = 0.5, v_\phi = 0.3, \mu_\chi = 0.01, \mu_\phi = 0.02$
below: $v_\chi = -0.5, v_\phi = 0.3, \mu_\chi = 0.01, \mu_\phi = 0.02$

Overall effect is a few orders of magnitude smaller than the production connected with the change of the vacuum.

7. Literature

- L.Kofman, A.Linde, X.Liu, A.Maloney, L.McAllister, E.Silverstein, arXiv:0403001 [hep-th]
- S.Enomoto, S.Iida, N.Maekawa, T.Matsuda, arXiv:1310.4751 [hep-ph]
- S.Enomoto, O.Fuksińska, Z.Lalak, arXiv:1412.7442 [hep-ph]

3. Change of the vacuum

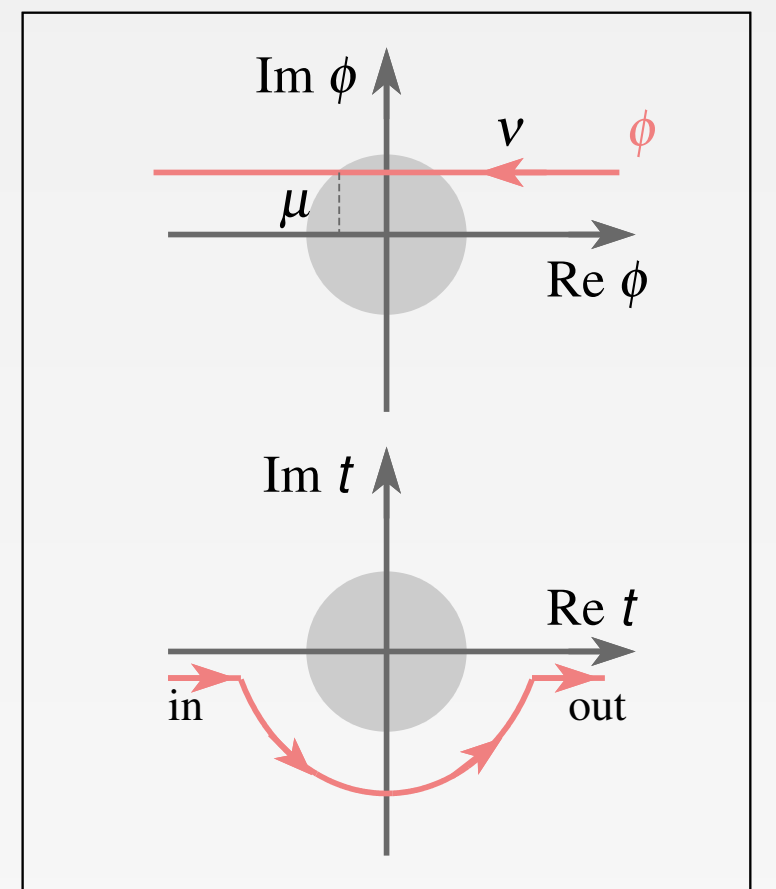
- superpotential: $W = \frac{g}{2} \Phi X^2$
- asymptotically: $\langle \phi \rangle(t) = vt + i\mu, \langle \chi \rangle = \langle \psi_\chi \rangle = \langle \psi_\phi \rangle = 0$
- background field in **non-adiabatic region**: χ and ψ_χ particles are produced
- produced particles induce a new **linear potential** and an attractive force ("oscillations")
- occupation number: $n_k^\chi = V \cdot |e^{-i \int^t dt' \omega_k(t')}|^2 = V \cdot \exp\left(-\pi \frac{k^2 + g^2 \mu^2}{gv}\right)$

- each time the **number density** $\left(n = \int \frac{d^3 k}{(2\pi)^3} \frac{n_k}{V}\right)$ of produced particles is:

$$- n_\phi \approx 0 \text{ and } n_{\psi_\phi} \approx 0$$

$$- n_{\psi_\chi} = \frac{(gv)^{3/2}}{(2\pi)^3} e^{-\pi g \mu^2 / v}$$

$$- n_\chi = \frac{(gv)^{3/2}}{(2\pi)^3} e^{-\frac{\pi(g^2 \mu^2 + m_\chi^2)}{gv}}$$



4. Rescattering

Equation of motion for a scalar field Ψ with the **source term** $J(x)$:

$$\left(\partial^2 + M^2(x)\right)\Psi(x) + J(x) = 0$$

Yang-Feldman equation:

$$\Psi(x) = \sqrt{Z} \Psi^{\text{as}}(x) - iZ \int_{t_{\text{as}}}^{x^0} dy^0 \int d^3 y [\Psi^{\text{as}}(x), \Psi^{\text{as}}(y)] J(y)$$

Generalized Bogoliubov transformation:

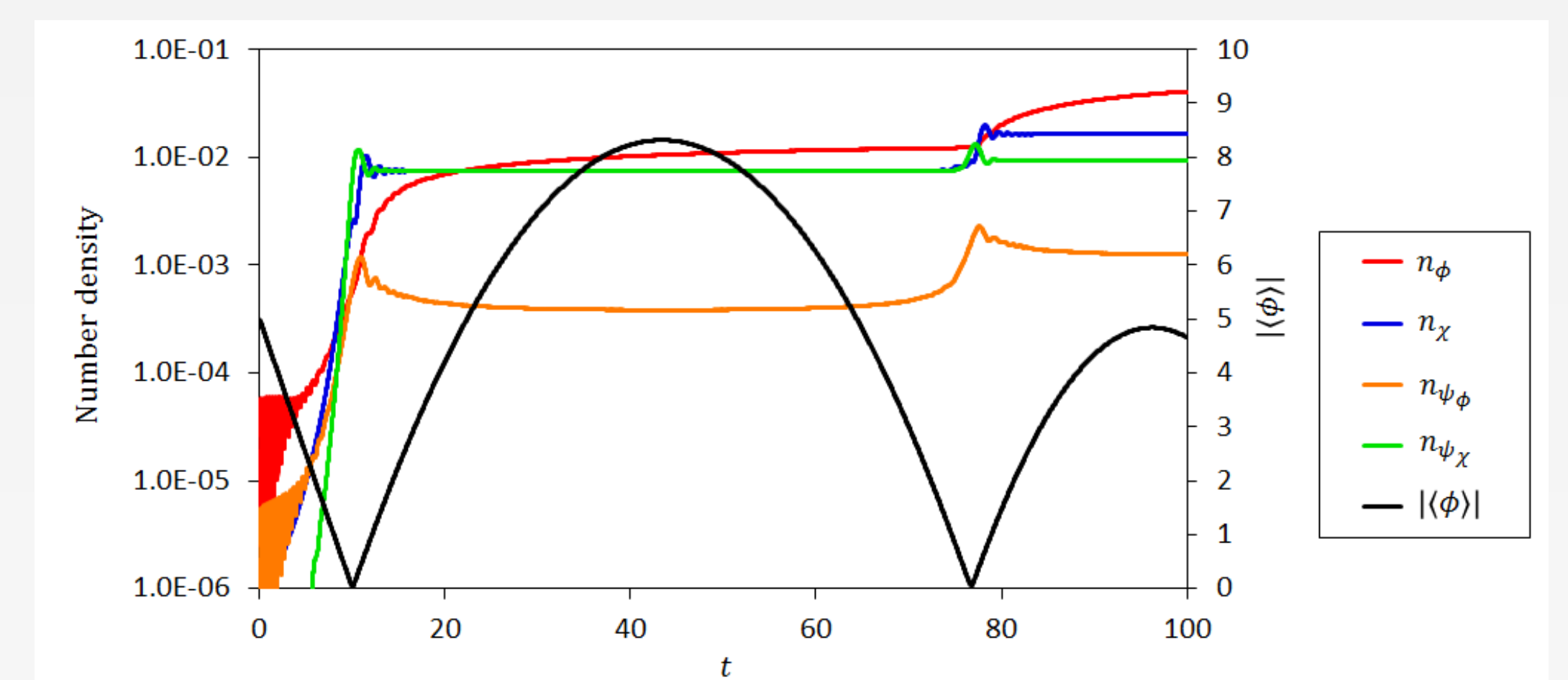
$$a_k^{\text{out}} = \alpha_k a_k^{\text{in}} + \beta_k a_{-k}^{\text{in}\dagger} - i\sqrt{Z} \int d^4 x e^{-i\vec{k}\cdot\vec{x}} \left(-\beta_k \Psi_k^{\text{in}}(x^0) + \alpha_k \Psi_k^{\text{in}*}(x^0) \right) J(x)$$

Occupation number:

$$n_k = \begin{cases} V |\beta_k|^2 + \dots & (\beta_k \neq 0) \\ 0 + Z \left| \int d^4 x e^{-i\vec{k}\cdot\vec{x}} \Psi_k^{\text{in}*} J | 0^{\text{in}} \right|^2 & (\beta_k = 0) \end{cases}$$

g	n_χ	n_{ψ_χ}	n_ϕ	n_{ψ_ϕ}
0.1	45.85	50.66	1.83	1.66
0.5	47.33	47.74	4.26	0.66
0.8	45.26	45.36	8.72	0.66
1.5	36.8	37.04	25.03	1.13
1.6	35.67	35.94	27.16	1.24
1.8	32.85	33.14	32.59	1.41
2	43.45	43.61	12.27	0.67

Number density of produced species as a part of the whole production (in %).



$$\phi(t=0) = 5 + 0.05i, \dot{\phi}(t=0) = -0.5, g = 2$$

6. Perturbative production

Perturbative production in the adiabatic region can be described by **Boltzmann equation**:

$$\dot{n}_k - Hk \frac{\partial n_k}{\partial k} = \frac{1}{E} \int \frac{d^3 p_2 d^3 p_3}{4(2\pi)^5 E_2 E_3} \delta^{(4)}(K - p_2 - p_3) \left(\sum_m (-n_{\vec{k}} + n_2 n_3) |\mathcal{M}_m|^2 \right) +$$

$$+ \frac{1}{E} \int \frac{d^3 p_4 d^3 p_5 d^3 p_6}{8(2\pi)^2 E_4 E_5 E_6} \delta^{(4)}(K + p_4 - p_5 - p_6) \left(\sum_m (-n_{\vec{k}} n_4 + n_5 n_6) |\mathcal{M}_m|^2 \right),$$

where $K = (E, \vec{k})$ and $k = |\vec{k}|$.

Preliminary conclusion:
perturbative production is comparable to non-perturbative one

